

The Design of an Innovative Automotive Suspension  
for Formula SAE Racing Applications

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### Abstract

This thesis details an analytical approach to an innovative suspension system design for implementation to the Formula SAE collegiate competition. It focuses specifically on design relating to geometry, mathematical modeling, energy element relationships, and computer analysis and simulation to visualize system behavior. The bond graph approach is utilized for a quarter car model to facilitate understanding of the analytical process, then applied to a comparative analysis between two transverse half car models. The second half car model contains an additional transverse linkage with a third damper, and is compared against the baseline of the first half car model without the additional linkage. The transverse third damper is an innovative design said to improve straight-line tire contact during single-sided disturbance, help mitigate the adverse effects of squat and dive, while not inhibiting the function of the anti-roll bar in cornering capability. Additional work is done investigating an optimization of suspension geometry through mathematical modeling in MATLAB of a four-bar linkage system. This code helps visualize the complex motion of the upright and calculates the wheel camber rate and variation to compare against tire data analysis to match maximum tire performance characteristics with camber angle.

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## Innovative Formula SAE Suspension Design

### **Chapter 1: Introduction/Background**

Undoubtedly, suspension systems are critical to the dynamic performance of a race car. From the word's roots, it means literally to "suspend," or keep something raised from the ground. This is exactly what vehicle suspension systems do; the body and chassis of a vehicle is kept suspended from the ground through a series of linkages and connections to energy storage and dissipative elements, which control the contact between the tires and the ground. In other words, springs (energy storage) and dampers (energy dissipation) are mounted between the vehicle's frame and control arms to translate the weight of the vehicle through these elements and into the wheels, which push down against the ground and by Newton's 3<sup>rd</sup> law push back up against the wheels to keep the car suspended. However, there is more to a suspension system than simply keeping something raised above the ground – if this were the case then we could simply use fixed non-compliant connections between the wheels and the frame to keep the frame from dragging on the ground. The disadvantage of such a system comes when the wheels hit a bump or pothole, and the entirety of that force is translated directly into the frame, likely causing fracture and failure. To avoid this, springs were placed in the connection between wheels and chassis so that when hitting a bump, some of that impact energy goes into compressing and extending the springs which minimizes the energy absorption required by the frame. Still, this caused another issue of riding discomfort as the springs freely oscillate up and down until enough energy is naturally dissipated to return back to a steady-state equilibrium position (which can take quite a while). Thus, another element

was added to quickly dissipate the energy that causes prolonged oscillation – the damper. However, not just any configuration of these elements will suffice; there are specific ratios and combinations of mass, spring stiffness, and damping coefficients that yield a comfortable riding experience and maintain the optimal contact patch between each of the tires and the road.

In researching how to design a racing vehicle's suspension system, nearly every source has something different to say about how and where to begin. A vehicle's suspension system is the most important subsystem of a dynamic performance racing vehicle. To narrow the scope of this investigation, this research is focused on ways in which the suspension system of a FSAE racing vehicle can be optimized. Thus, the results of this investigation may be directly applied to the design of Liberty University's Formula SAE team. The suspension system is where the rubber meets the road, so the design is focused on best controlling that physical interface across a range of loading conditions. In dynamic racing applications, as the sprung mass of the vehicle chassis and driver are being accelerated they produce a reactive force that must be balanced and utilized to both maintain and increase traction and contact with the road, to maximize the ability of the wheels to keep a secure grip and maintain as high a speed as possible without sliding off the track.

As such, a suspension system is designed with two primary purposes in mind; first to accelerate from a stop to top speed and decelerate to a stop as quickly as possible in a straight line, and second to maintain the highest speed possible through a turn. In the first of these, the goal is to allow the vehicle's inertia to cause a controlled "squat" on the

power producing wheels as the effective center of mass is shifted towards the rear wheels producing a greater normal force over the rear wheels to increase the possible force of friction before slipping. For the second of these, the goal is similar, in that the inertia and momentum of the sprung mass must be controlled from excessive roll when changing yaw as the vehicle travels through a turn (See Figure 1).

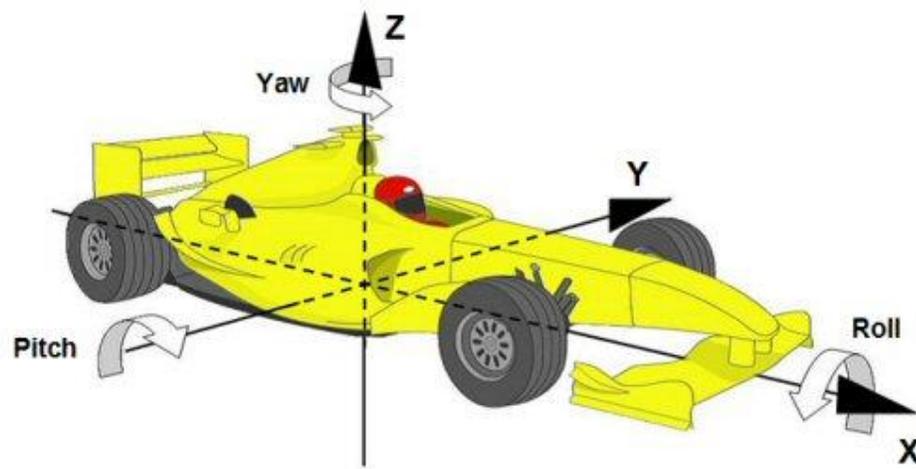


Figure 1: Axis definition and nomenclature of rotational movements about each axis.

When turning, the center of mass is effectively shifted towards the outside of the turn, putting more normal force on those outside wheels (called load transfer). If excessive this can be dangerous, and potentially cause the vehicle to either slip from the track or roll completely if the center of mass is displaced beyond the limits of the outer wheels from too much side loading. The degree of body roll is controlled by the stiffness of the springs and anti-roll bar (essentially a torsional spring) in the suspension, while the dampers are designed to control a rapid and hard-hitting force like a bump or pothole in the road, to dissipate the energy of the oscillatory result of a wheel displacement due to the nature of springs. Both of these energy components (the springs and dampers) in a

vehicle's suspension system are of critical importance to control the vehicle's dynamic behavior. Frequency response analysis may be done with a variety of input parameters to find coefficients that are best optimized for the vehicle's performance. Such analysis is demonstrated in the latter sections of this investigation. However, another important component of suspension design is the system geometry, providing the lengths and connection angles of the linkages between wheels and chassis. The conventional suspension system referenced here is the common double wishbone suspension, pictured in Figure 2, as found in most performance vehicles. It is acknowledged that there is a plethora of various types of suspension such as MacPherson Strut, leaf spring, and trailing arm to name a few, but within the field of racing the double wishbone is the most widely implemented design.

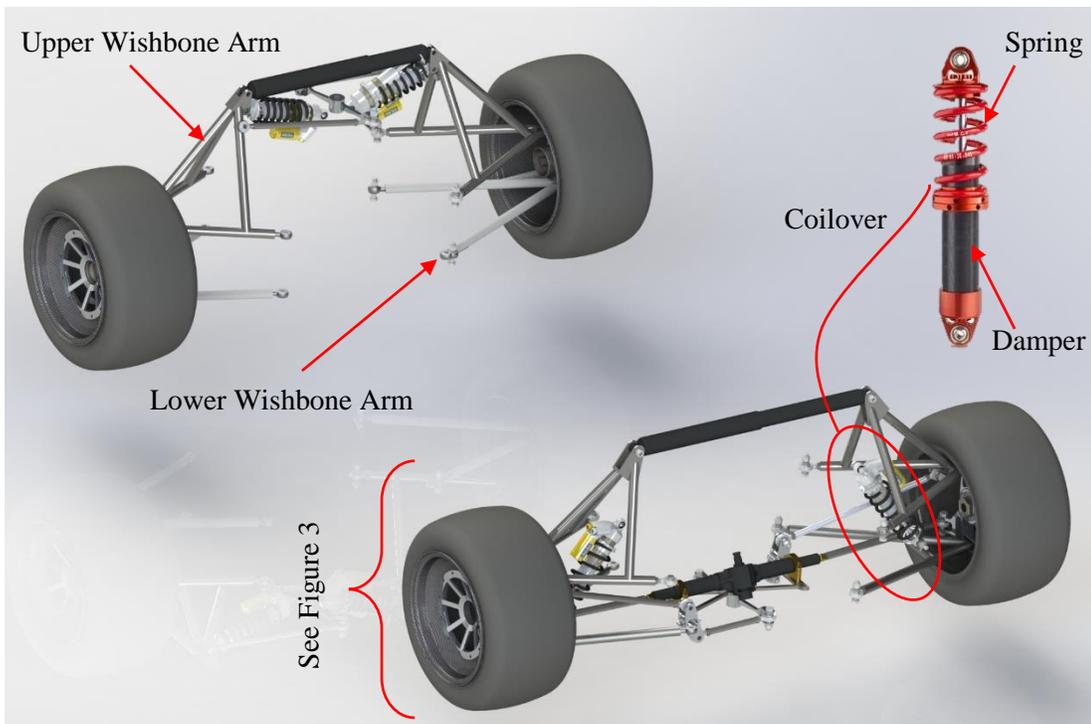


Figure 2: Double wishbone suspension system, partial car model.

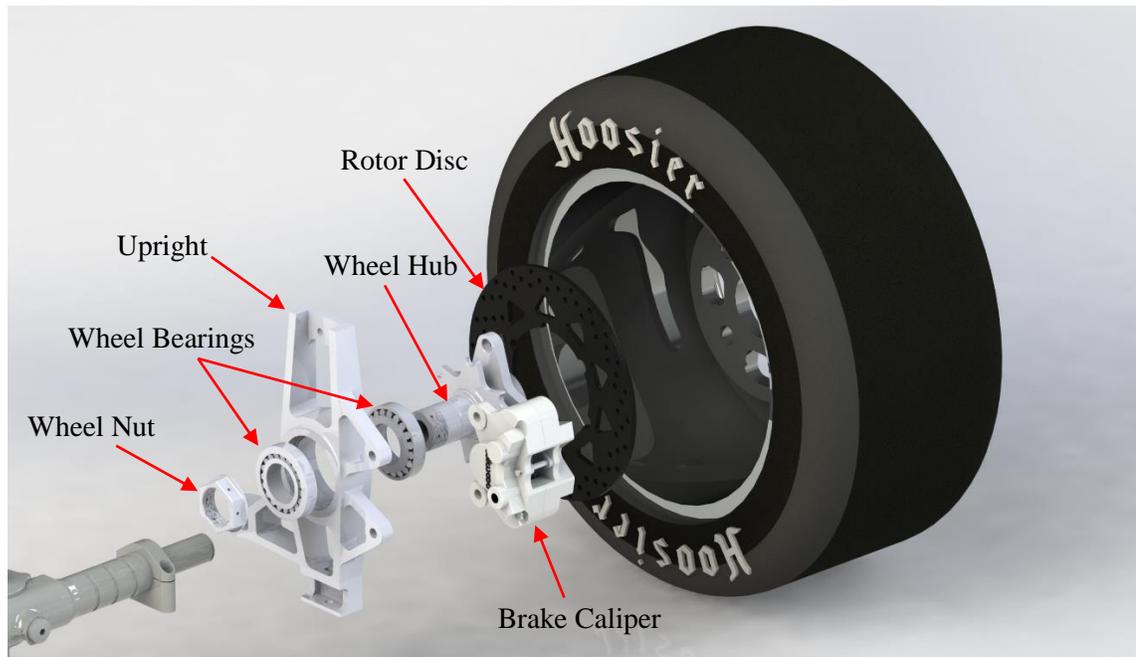


Figure 3: Exploded view of wheel assembly for component nomenclature.

The movement of a double wishbone suspension system can be mathematically reduced to an analysis of a double-rocker four-bar linkage. Consider the chassis to be the inertial frame of reference from which other movement is analyzed, despite normally thinking of things moving relative to the ground. The stationary link is the straight line between the wishbone connection points on the chassis. In other words, two points on the chassis should never move relative to one another because it is a rigid structure. (Any compliance or deformations of the chassis members shall be negligible.) The two wishbone arms both function as rockers, simply pivoting from their pinned (single axis rotation) connections to the chassis as their ends move up and down with the movement of the wheel. The component of interest is the wheel assembly. The upright is in fixed connection to the wheel hub, which is connected to the wheel by the lug bolts. Therefore, these will all be effectively considered the same assembly, called the “wheel.” The movement of the wheel under load determines how the driver inputs of power and turning

connect to the road to provide feedback and give direction to the driver's movements. Unless a double wishbone suspension geometry is designed to be a perfect parallelogram, the wheel must change its angle as it translates up and down, as demonstrated by the four-bar linkage analysis. Therefore, the wheel is in complex movement, both translating and rotating relative to the chassis. This has a direct effect on camber angle (See Figure 4), which is a key parameter within the calculations of a wheel's contact patch. The contact patch ultimately determines the maximum friction force capable of being transmitted between the tire and road before slipping.

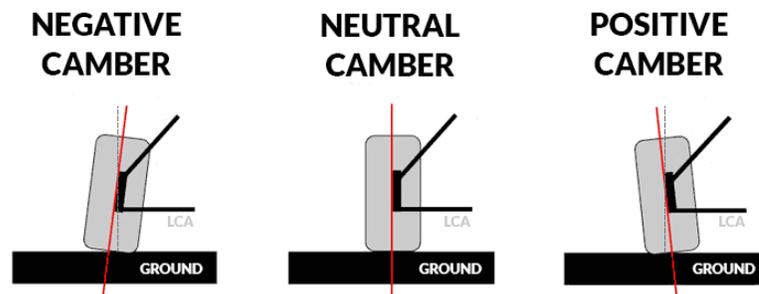


Figure 4: Visual depiction of camber angle.

Generally, it is most desirable for a vehicle to have neutral camber in wheel alignment for even tread wear and maximum contact patch area during normal driving. However, given that racing is not normal driving, a perfectly neutral camber is not always desirable because it does not yield a maximum contact patch area when the driver takes very hard turns and the wheels are displaced according to their suspension geometry path. In fact, it is better to create negative camber on the wheels on the outside of a turn, as the wheel is forced upward by the increased effective load it is supporting through its springs causing them to compress. The precise desired angle of negative camber is more difficult

to calculate, but it is important to visualize the path of the wheel and camber angle throughout projected wheel displacement (See Figure 5).

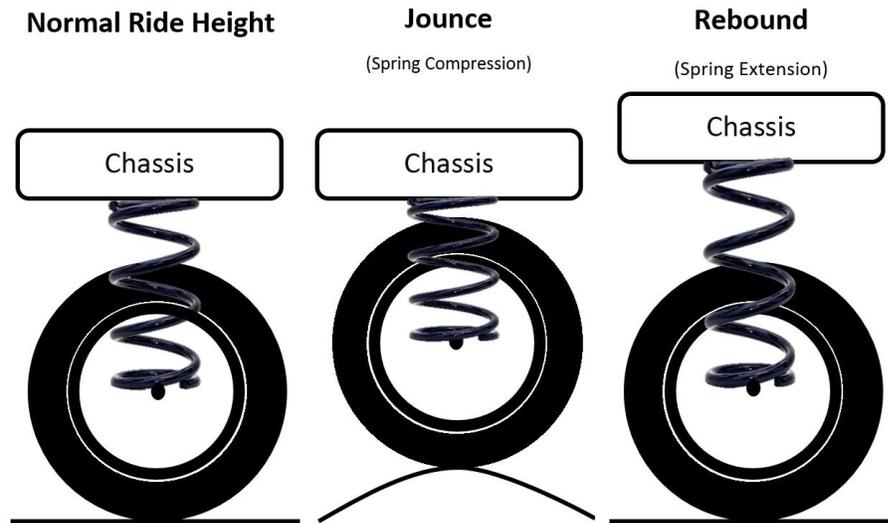


Figure 5: Depiction of jounce and rebound relative to spring compression and extension.

Therefore, I created a program in MATLAB (see Four-bar Linkage Motion Analysis Section) that could plot the movement of a four-bar linkage through a range of movement, inputting custom lengths of all the links. In this way, with the design of the FSAE vehicle in mind, dimensions directly relating to desired variables of the vehicle could be tested and manipulated. These included track width, ride height, and chassis frame members to design where the suspension components could be mounted, and what their corresponding lengths should be to have a desired complex path demonstrated by the wheel through jounce primarily, but also in rebound of the wheel.

With all the interacting factors that need to be considered during suspension design, it makes suspension optimization a challenging and rewarding field when done right. Further, given the many interdependent variables that may be altered, it leaves

room for seemingly infinite design possibilities, always leaving room for improvement and development.

### **Literature Review**

**Transverse 3<sup>rd</sup> damper integration.** Certainly the most ambitious innovation to racing vehicle design discussed here is the implementation of a third damper mounted transversely between both the front and rear pairs of wheels. This design was pioneered by Koenigsegg; a company whose mission is to take areas of compromise in vehicle design and find ways to minimize or eliminate such compromise. In suspension design, nearly everything is a tradeoff. Koenigsegg saw one of these areas, in which superior cornering capability came at the cost of driver comfort and straight-line traction – and decided to do something about it.

There are many types of vehicle racing. As a result, the racing vehicle is designed very differently to optimize the conditions it is subjected to. For example, a drag racer and a Formula 1 vehicle are each very successful in straight-line and dynamic circuit tracks respectively, but put one in the other's race, and it will fail miserably against the competition. An engineer's job is to utilize physics and mathematics, with a scientific approach, to understand as much as possible about the materials and forces at play to make the greatest use of the strengths and minimize the weaknesses in a design. The incorporation of a third damper is most effective in the design of a suspension intended for a more dynamic circuit – like those raced by Formula 1 vehicles. This racing setting demands that a vehicle be able to corner exceptionally well, but it cannot come at the cost of general traction and driver comfort, which is critical in endurance events.

A traditional approach to this trade-off dilemma is to tune the stiffness of the anti-roll bar, finding a sweet spot based on mathematical modeling and on-track testing so that it is as stiff as allowable without too much compromise in traction. This traction compromise comes from the fact that the anti-roll bar is the only direct connection between two otherwise independently suspended wheels, serving to force both wheels to move in the same direction to reduce the body roll generated by one wheel being forcibly compressed on the outside of a turn when cornering from lateral load transfer. As this outside wheel is compressed under greater load, a torsion is generated in the anti-roll bar, which is then applied to the inside wheel, compressing it as well and therefore allowing the chassis to come back closer to the road on the inside and balance out the roll.

However, as useful as it is to balance out roll in a corner as one wheel imparts a force on the other through the torsion of the anti-roll bar, this same tendency is detrimental to traction in a straight line when bumps on the track surface cause wheel displacements on one side of the vehicle that are then also imparted to the other side. When this happens, a wheel is being unnecessarily “lifted” from the road when the opposite wheel hits a bump. (More realistically the wheel is not fully lifted from the pavement, there is only a reduction in the vertical load applied to the wheel.) Such a reduction in vertical load means that the maximum possible friction force between the tire and road is significantly reduced, because  $f_f = \mu N$  and  $N$  is directly proportional to  $f_f$  (where  $f_f$  is the friction force,  $\mu$  is the friction coefficient, and  $N$  is the normal force).

A third damper, however, mounted transversely across the width of the vehicle similarly to the anti-roll bar establishes a second connection between the two otherwise

independently suspended wheels. This time, there is an energy dissipating effect from the damper, as a force imparted on either wheel is translated by a pushrod into a displacement of the piston within the damper cylinder, effectively converting the kinetic energy from the motion of the wheel into heat as the viscous damping fluid is forced through the capillary holes of the piston head creating viscous friction within the fluid. The damping coefficient of energy dissipation is controllable by the viscosity of the damping fluid, either by filling the chamber with different fluids, or utilizing a fluid whose viscosity can be altered. Magnetorheological fluids are capable of this by the introduction of a magnetic field within the fluid to change viscosity, and will be discussed in the following section.

With an energy dissipating element, the disturbances felt from one wheel to another are greatly reduced, while cornering capability is not impeded. The damper responds effectively to rapid displacements as the damping force is proportional to the velocity of the internal piston, and permits the anti-roll bar to accomplish its purpose in more sustained displacements such as those of cornering. Therefore, a stiff anti-roll bar may be implemented to improve cornering capability, while the negative effects of opposite wheel displacements and loss of traction are mitigated by the dissipation of those forces through the transversely mounted third damper, and driver comfort and road-feel are also improved. Further, with the implementation of magnetorheological fluid and sensor feedback for automated and calculated control of the damping coefficient, the control of the entire suspension system is taken to a new level, and vehicle handling is more capable than ever before.

**Magnetorheological damper integration.** Conventional dampers are constructed with two basic components, the piston and cylinder. (See Figure 6) Though there are many variations of form, including monotube, twin tube, internal bypass and more, it serves the same purpose. The damper's function is to dissipate energy, as a displacement force either in tension or compression from the kinetic energy of wheel movement displaces the piston, and is converted into heat within the hydraulic fluid by way of viscous friction. That heat then slowly moves out into the environment by conduction to connected parts or is carried away into the air as forced convection removes heat energy.

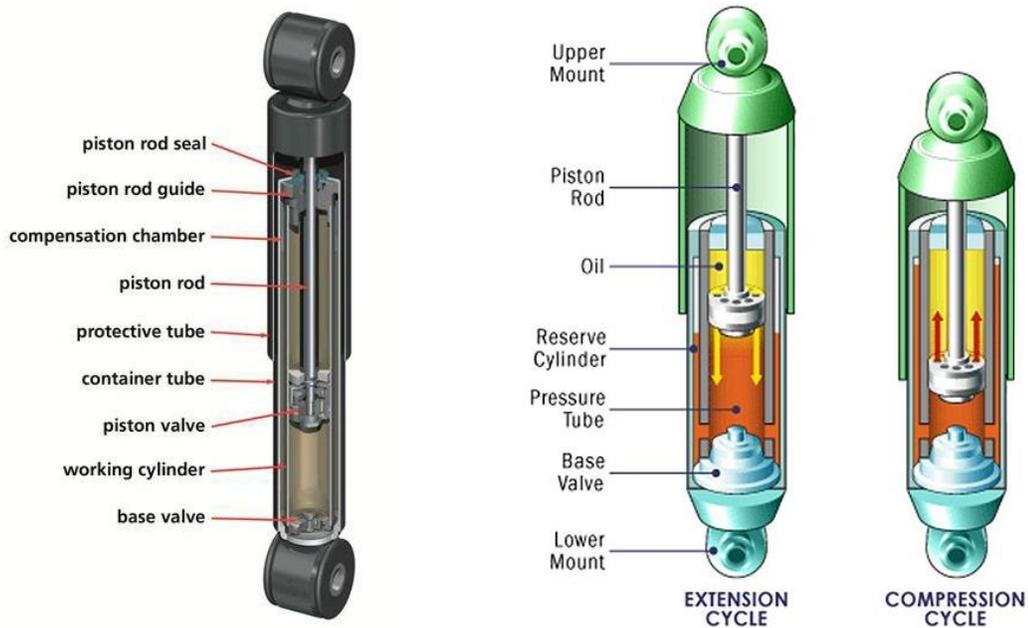


Figure 6: Cutaway view of conventional damper, twin-tube model.

As the piston is displaced within the hydraulic fluid, it increases pressure within the hydraulic fluid on the compression side. This pressure is allowed to reach equilibrium with the side opposite the piston by designing tight passages for the hydraulic fluid to flow either around the sides of the piston at the walls of the cylinder, or through small

holes in the piston head. It takes significant energy to force a viscous fluid through these tight passages, and the viscous shearing happening within the fluid heats it up.

Dampers exist to control the repeated oscillations that would otherwise be present from suspending the chassis on springs. Engineers tune the damping coefficient to a specific value to establish an underdamped, critically damped, or overdamped system. A simplified model of the system behavior can be easily interpreted from the plots of a second order differential equation of the following form:

$$F(x) = m\ddot{x} + b\dot{x} + kx \quad (1)$$

In which  $m$  represents the sprung mass,  $b$  is the damping coefficient, and  $k$  is the spring constant. A derivation of this mass-spring-damper model is given in Appendix A for further exploration, but the plot shown by Figure 7 reveals the behaviors possible by tuning these three coefficients. The parameter zeta taken from that derivation is the damping ratio, defined such that the system is critically damped when zeta equals one.

$$\zeta = \frac{b}{b_c} = \frac{b}{2m\omega_n} = \frac{b}{2m\sqrt{\frac{k}{m}}} \quad (2)$$

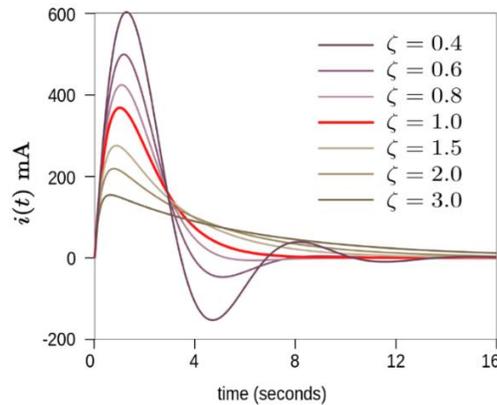


Figure 7: Oscillatory behavior control of mass-spring-damper system by tuning of zeta parameter.

Realistically however, this tuning process can be very laborious and time consuming to make even small changes. Damping coefficients may be changed either by altering the viscosity of the hydraulic fluid or altering the passages through which it must travel when under pressure from the piston. This means disassembling the damper to either drain and replace the hydraulic fluid for one with a different viscosity, or exchanging washers on the piston head to change fluid flow through the passages. With sufficient data from field testing we are able to very precisely predict and measure these damping coefficients. However, there exists a much more versatile means of changing the behavior of a damper.

Magnetorheological fluid is a fluid with iron particles in it that responds to a magnetic stimulus, changing shape and/or viscosity. Engineers have implemented the use of magnetorheological fluid within dampers by adding iron particles to hydraulic fluid, such that when a magnetic field is induced within the fluid, it causes these iron particles to align themselves. This particle alignment increases the viscosity of the fluid by making it more resistant to deformation by shear stress. The magnetic field is generated by a coil of wire running perpendicular to the length of the cylinder that has electric current sent through it. Ampere's Law dictates that the magnitude of the magnetic field generated around a current carrying wire is proportional to the current in the wire. In this way, as magnetorheological fluid viscosity is proportional to the strength of the magnetic field, so it is also proportional to the magnitude of the current sent through the wires. This kind of damper is pictured in Figure 8.

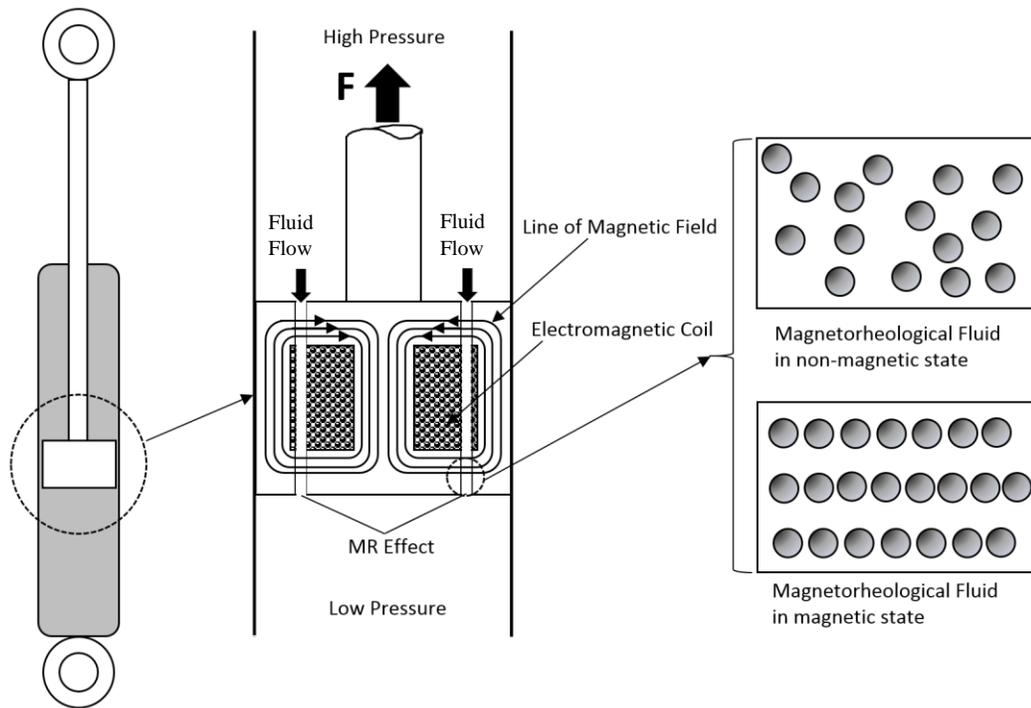


Figure 8: Cutaway view of magnetorheological damper functionality.

Each year more sensors are being placed on various components within a vehicle to have their data fed through the vehicle’s ECU for automatic adjustments based on programmed code. The capacity for rapid tuning and real-time adjustment has revolutionized the racing industry. With magnetorheological dampers, the effective damping coefficient can be changed a fraction of a second after hitting a bump, or as the brakes are being pressed to work against vehicle dive. In fact, Koenigsegg has their dampers networked so as to be able to make remote adjustments to any of their cars at any time wherever they are in the world, quickly enough to aid the vehicle as it races along a track. Hydraulic fluid viscosity adjustments as rapid as current may be sent through a wire is changing the way the racetrack is approached to optimize every second of tire contact with the road through accelerations, corners, banks, and more.

Further, to demonstrate the superior adjustability of magnetorheological dampers, magnetorheological dampers have been designed with current carrying wires wound in opposite directions and timed such that the residual coil inductance from the magnetic field generated by the first coil is negated by the second. For example, the first coil carrying electric current induces a magnetic field, aligning iron particles in the fluid. When the viscosity is desired to be lower however, power is decreased or shut off to the first coil, but it takes a fraction of a second for the magnetic field to dissipate due to residual eddy currents caused by inductance, meaning that those aligned iron particles are not yet released from their positions in the fluid so the viscosity takes time to drop down to a lesser value. These eddy currents are counteracted by adding a second coil to the damper so that as the first is being shut off, the second set activates and produces a magnetic field in the opposite direction working to immediately disrupt the field that existed and scatter the aligned ferrous particles, only activating long enough to disrupt the effects of inductance from the first coil and then shutting off before realigning the particles within a new magnetic field orientation.

## **Chapter 2: Geometry Considerations**

### **Tire Data Analysis**

A racing vehicle's suspension system was stated as the subsystem where the rubber meets the road, and that is exactly the location where the definition of suspension geometry must begin. Professional racing teams spend substantial money on tire testing with carefully calibrated machines. The data from tire testing can be analyzed to determine the optimum orientation of the tire for given loading scenarios, which then

determine how the wheel path should be constrained by the placement of suspension linkages. For example, in a cornering scenario if the outside front wheel is considered, it experiences a greater vertical load (due to load transfer from centripetal acceleration), as well as a significant lateral load which directly correlates to the friction force that keeps the wheel from sliding on the pavement. Many racing tires can increase their lateral load capacity in a corner if there is some negative camber gain to angle the tire against the pavement. However, too much camber can be even worse than none at all. A number of years ago the Tire Test Consortium (TTC) was founded to provide a database for Formula SAE students to access and retrieve this data to analyze for suspension design. This database details exactly the test procedures that govern each set of data for proper interpretation, and includes more than a dozen variables. It is the student's job to select a tire compound and dimensions, and get to work analyzing hundreds of thousands of lines of data to make useful plots that will tell how the tire can best perform in the expected racing conditions. The first of these plots is Slip Angle vs. Lateral Load. Slip angle is understood according to Figure 9 as the angle difference between the steered intended angle, and the actual traveled angle due to tire deformation at the contact patch.

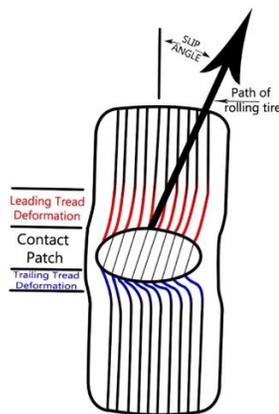


Figure 9: Definition of slip angle; a result of applied lateral load deforming the tire at the contact patch.

By holding all other variables constant while only manipulating one variable at a time, the Slip Angle vs. Lateral Load was plotted across the range of tested inclination angles (which simulate camber) to visualize how the lateral load capacity of the tire changes with camber. This was done for a number of vertical loads because the normal force acting on the tire changes due to load transfer when cornering, so the outside wheels will experience a greater normal force than the inside wheels.

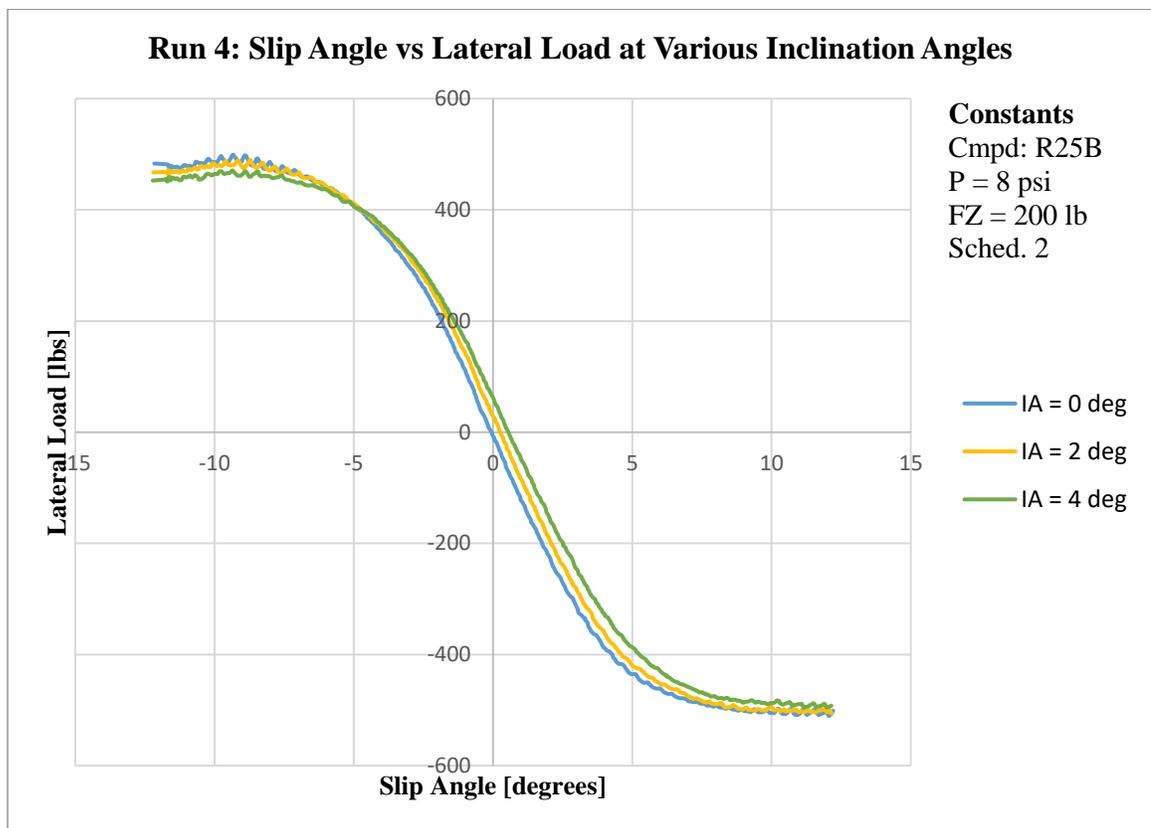


Figure 10: Slip Angle vs. Lateral Load for a Hoosier 16x7.5-10 R25B compound tire at 8psi, with a normal force of 200lb, across three camber angles (IA) of 0, 2, and 4 degrees.

The curve shown in Figure 10 demonstrates the effectiveness of the tire at resisting a lateral load with the friction contact between the tire and road surface over a range of slip angles. This is important because when cornering, Newton’s 1<sup>st</sup> law says

that the vehicle will tend to continue going straight, unless accelerated towards the center of curvature dictated by turning the steering wheel. The centripetal force generated that pushes the car towards the center of curvature is the friction force acting between the tires and road surface. The greater the friction capacity at this contact patch, the more centripetal force can be generated to take turns at higher speeds. The negatively sloped linear portion in the middle means that the tire is effectively resisting lateral loads; for small slip angles the tire linearly increases in lateral load capacity. However, when the curve begins to plateau at either end, this shows that the lateral load capacity of the tire is tapering to its limit at high slip angles, so when the curve goes flat the tire has lost grip on the road and will begin to slide across the pavement because the frictional force is not enough to resist the lateral load from centripetal acceleration. In Figure 10 one can clearly see that with different camber angles, the curve changes slightly and lateral load capacity changes. In fact, it appears in this plot that the camber angle of zero (refer to Figure 4) has the greatest peak lateral load capacity, but in the linear range the higher camber angles have a greater lateral load. However, the answer is far from being reached; these tires must be tested for various compounds, pressures, and normal forces at the very least before being able to determine the desired camber change with wheel travel for the suspension geometry. Figure 11 shows the comparison between two different tire compounds, R25B and LCO, of the same tire dimensions.

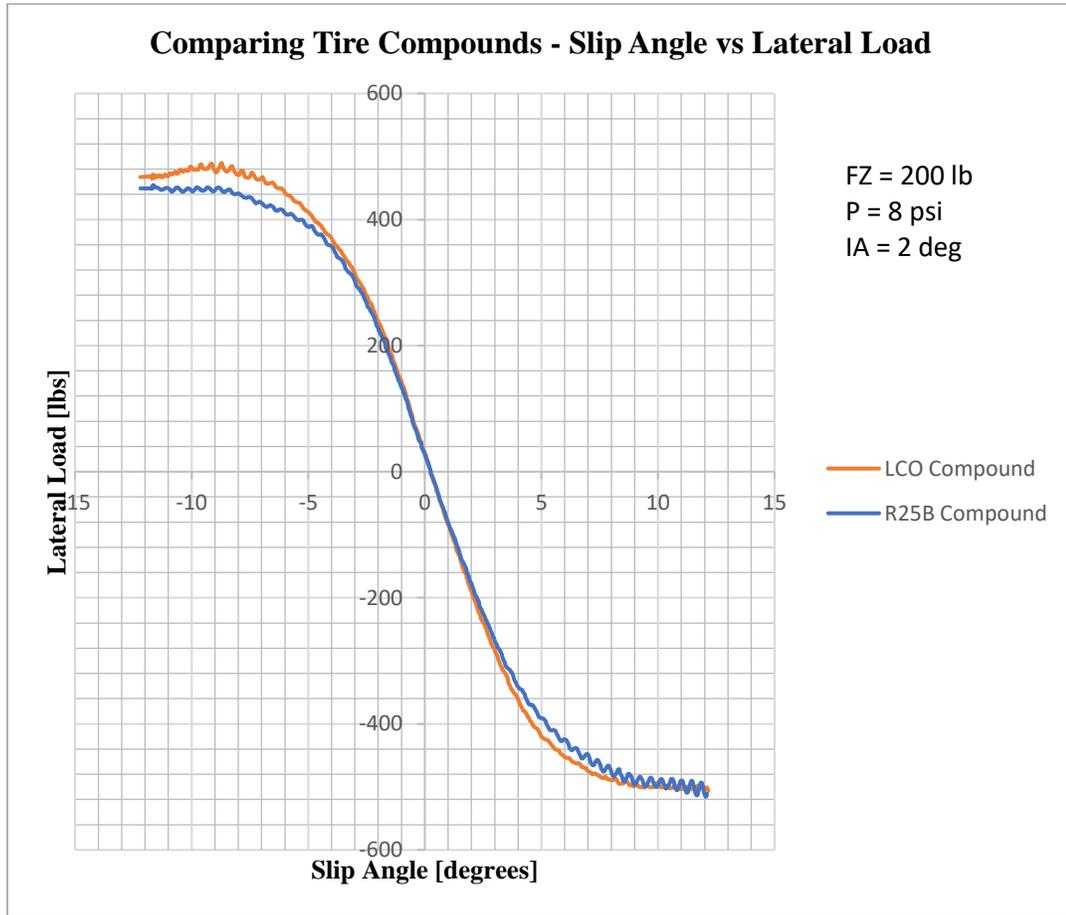


Figure 11: Comparison of R25B and LCO tire compounds at 8psi, camber of 2 degrees, and 200lb normal force.

According to Figure 11, it would appear that the LCO compound is the clear winner with the higher peak lateral load capacity in the same conditions. However, Figure 12 plots the same two tires with the same pressure and normal force, simply at a camber of zero degrees, and the R25B compound takes the lead with a higher peak lateral load capacity and a longer stretch of linear behavior in the curve.

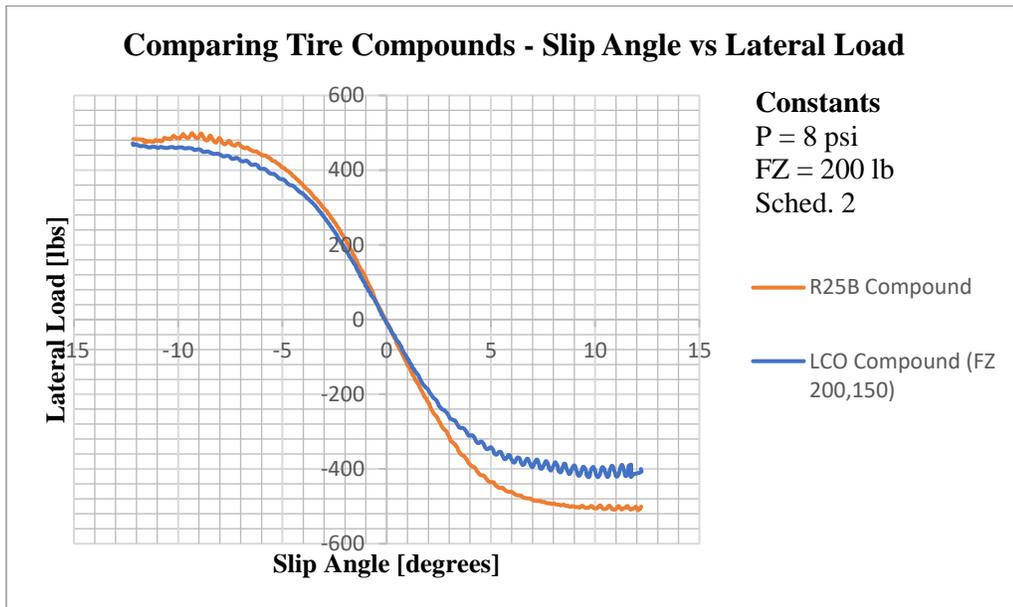


Figure 12: R25B vs LCO tire compounds compared at 8psi, 200lb normal force, and camber of 0 degrees.

Further still, changes in tire pressure also change this curve behavior significantly.

Figure 13 shows the same tire and conditions other than the tire pressure changing between 8psi and 12psi, and it is clear that the higher pressure has a higher peak lateral load capacity and longer stretch of linear behavior between slip angle and lateral load.

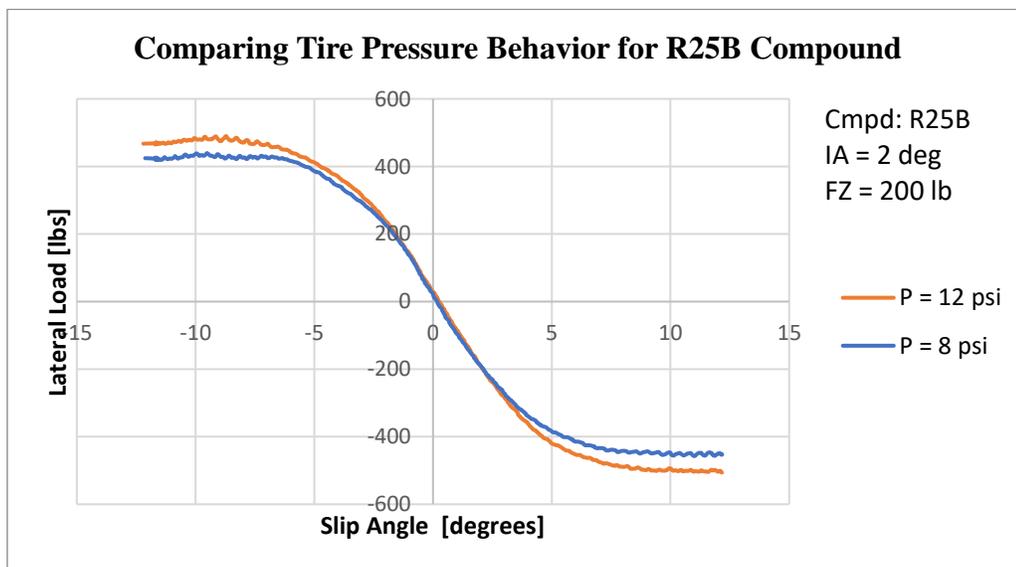


Figure 13: Slip Angle vs Lateral Load plot comparing two different tire pressures for R25B tire compound at 2 degrees of camber and a normal force of 200lb.

Therefore, without detailing the entirety of the tire analysis process, the previous few plots should be sufficient to demonstrate the complexity in objectively defining the desired camber rate in relation to vertical and lateral loads. Once sufficient analysis has been done to make an argument for a specific camber rate limitation based on expected wheel travel, the control arms in a double-wishbone system must then be designed accordingly to yield such a camber rate with the expected wheel travel occurring during lateral load transfer of cornering. The resulting camber at the maximum jounce position during a cornering scenario must be a summation of both the ideal camber from slip angle vs lateral load curves to maximize lateral load capacity, as well as the angle of body roll resulting from the difference in spring compression on opposite sides of the car. In this way, if the tire data shows peak lateral load capacity at 1.5 degrees of camber, and the body rolls 1 degree in cornering from load transfer, the suspension geometry should provide a total camber of 2.5 degrees at the wheel deflection expected in the same cornering scenario.

**Roll center.** The roll center is the location about which the sprung mass will rotate when experiencing body roll when cornering, and is one of the first things that must be defined when designing suspension geometry. The ratio of the height of the roll center relative to the height of the vehicle's center of gravity determines the percentage of anti-roll. Anti-roll is a measure of how much of the roll forces are reacted by the control arms rather than the springs and dampers. Figure 14 shows a SolidWorks sketch of a front view cross-section of the vehicle to place dimensions on things that would not change, with extended lines to find the location of the instantaneous center of rotation of

a wheel (IC). Then a line from the center of the tire contact with the road to the IC location is drawn. The point where this line crosses the vertical line passing through the center of gravity (assuming a laterally symmetric vehicle mass distribution) yields the location of the roll center (RC). The distance between the center of gravity and the roll center is the leverage of the roll moment. The closer they are together, the less body roll is experienced with the same force, thus increasing anti-roll geometry by putting a greater proportion of force into the control arms. However, while less body roll is generally a good thing, it also leaves less room for control of the movement of the sprung mass through the springs and dampers. A general range within the racing industry is about 15 to 30% anti-roll.

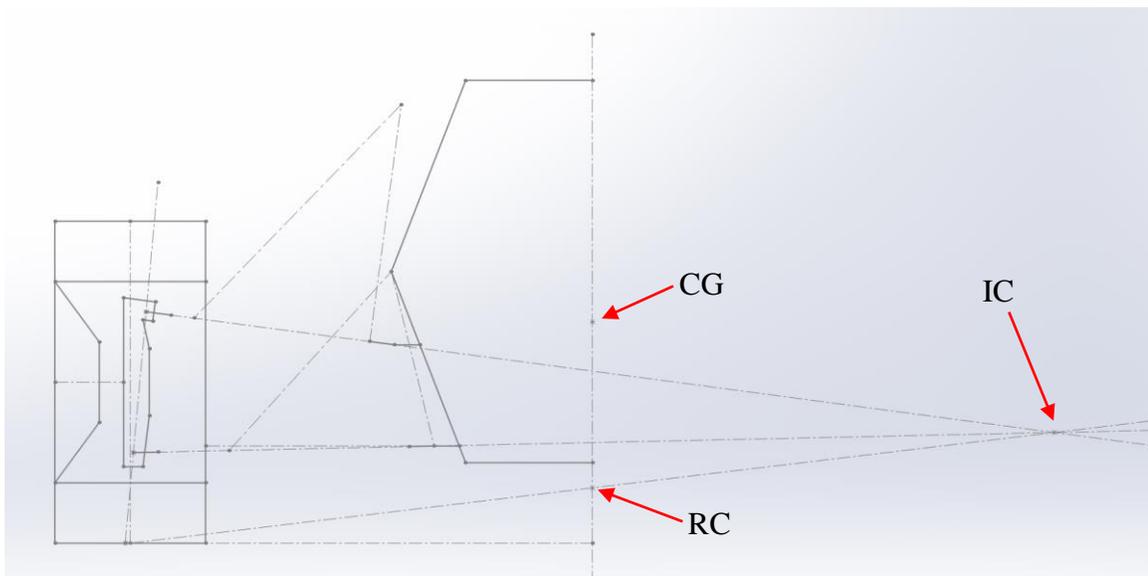


Figure 14: SolidWorks sketch to plan out roll center (RC) calculations. IC = instantaneous center, CG = center of gravity.

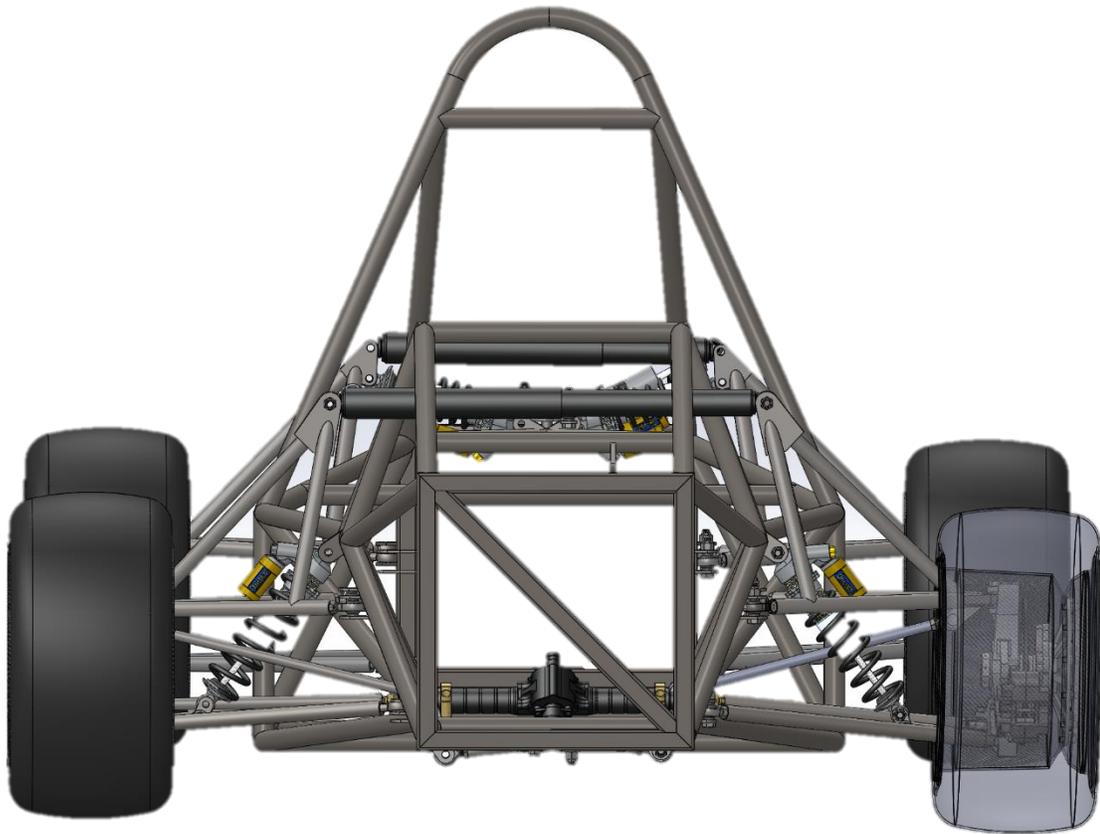


Figure 15: SolidWorks model front view of suspension system shown to match the geometry of the sketch in Figure 14.

The horizontal location of the instantaneous center along the line generated by the center of the tire contact and the roll center is still variable at this point. Moving the IC will adjust the mounting locations of the control arms to the frame, and therefore change the relative lengths of the control arms. Both these adjustments have direct effects on the path that the wheel travels in jounce and rebound. Not only are the arcs that each rocker linkage makes changed, but the equilibrium position of the system relative to the angle from the horizontal of the control arms is also changed with a new IC location. Therefore, a means of visualizing the direct effects of these changes on camber rate was necessary, and led to the creation of a code in MATLAB to solve the four-bar linkage equations and make helpful plots.

**Four-bar linkage motion analysis.** Concurrently with the RC calculations, when mounting positions for the control arms were found that gave the desired percentage of anti-roll, the suspension geometry then had to pass the camber variation test. A MATLAB program was generated to yield the changes in camber angle over a length of wheel travel of a double-wishbone suspension system, analyzing the system as a four-bar linkage. Figure 16 represents the vector analysis notation to locate the points in the linkage system as they move from a changing  $\theta_2$  input. The following Figures Figure 16 through Figure 20 demonstrate the final decided geometry for the rear suspension, having determined an IC location that yields a desired camber change through wheel travel.

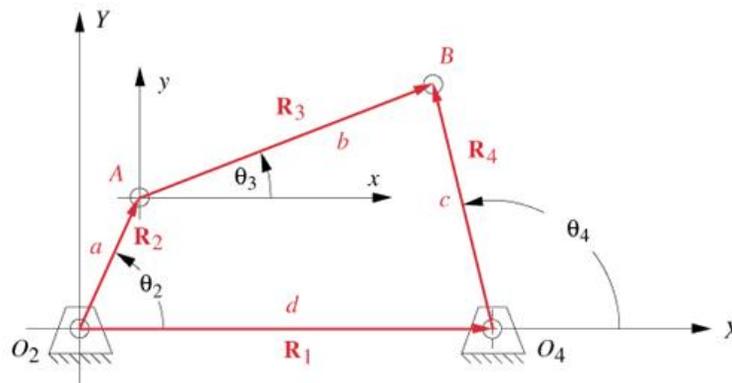


Figure 16: Four-bar linkage system defined by vectors and variables from which MATLAB variables were generated for suspension camber analysis.

The mathematical analysis required to solve this system is not highly relevant to suspension design, but can be found within the code for the MATLAB program in Appendix B. The first plot shown by Figure 17 shows a shape very similar to the one seen in Figure 16, but also shows the instantaneous placement of all the linkages over a  $\theta_2$  variation of 10 degrees. This 10-degree change simulates the maximum desired wheel travel (and minimum required travel) starting from one inch in rebound, passing the equilibrium position, and extending to 1 inch in jounce. Table 1 gives important

dimensions relevant to this system. The plots are rotated 90 degrees counter-clockwise from their natural orientation on the ground because that made writing the code and solving for the locations of each linkage easiest, so keep in mind that for these plots the ground is actually the right side of the plot.

Table 1: *Dimensions of rear suspension system*

Linkage	Location in Figure 9	Length (inches)
Upper control arm	Left side	12.029
Lower control arm	Right side	14.739
Chassis	Bottom	5.540
Upright	Top	7.026

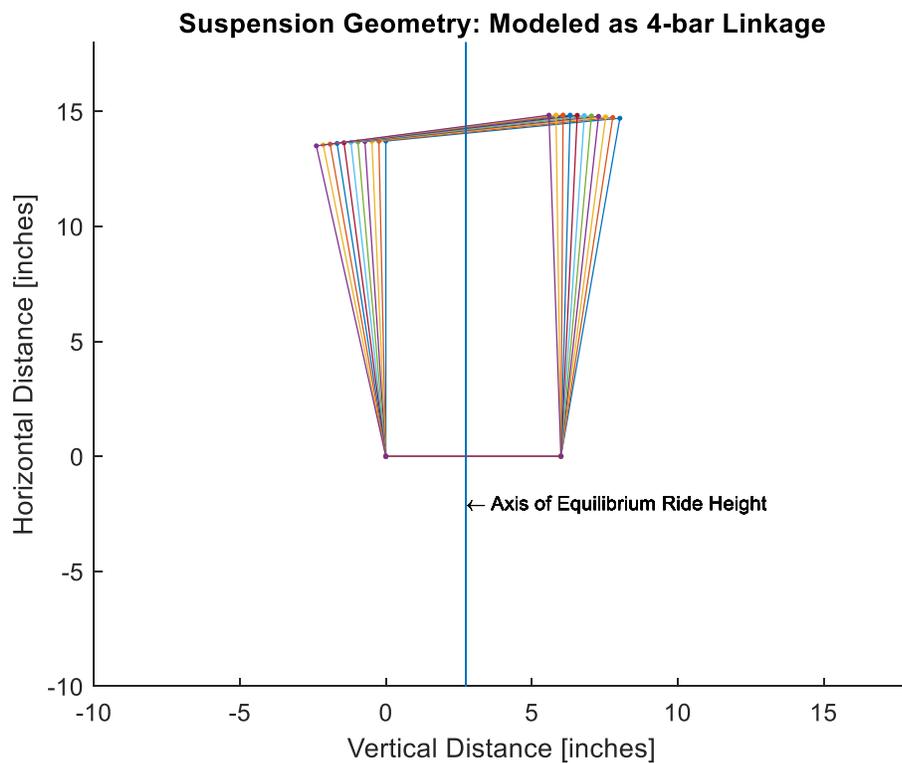


Figure 17: Overall plot of suspension geometry modeled as a four-bar linkage through a 10 degree variation in  $\theta_2$ .

Note that in Figure 17 the upright appears to be more steeply angled than it should be, because this linkage merely connects the control arm mounting points, which must

account for the Kingpin inclination angle. The resultant angle of the wheel is 8.3 degrees different than that of the inclination angle, such that at the equilibrium ride height position the wheel is perpendicular to the ground. When this system is pushed farthest to the right, that simulates the maximum wheel travel in rebound, while pushing it farthest to the left simulates maximum wheel travel in jounce. One can see slight variations in the rotation of the upright linkage as it translates, so an additional plot in Figure 18 was made to view this up close.

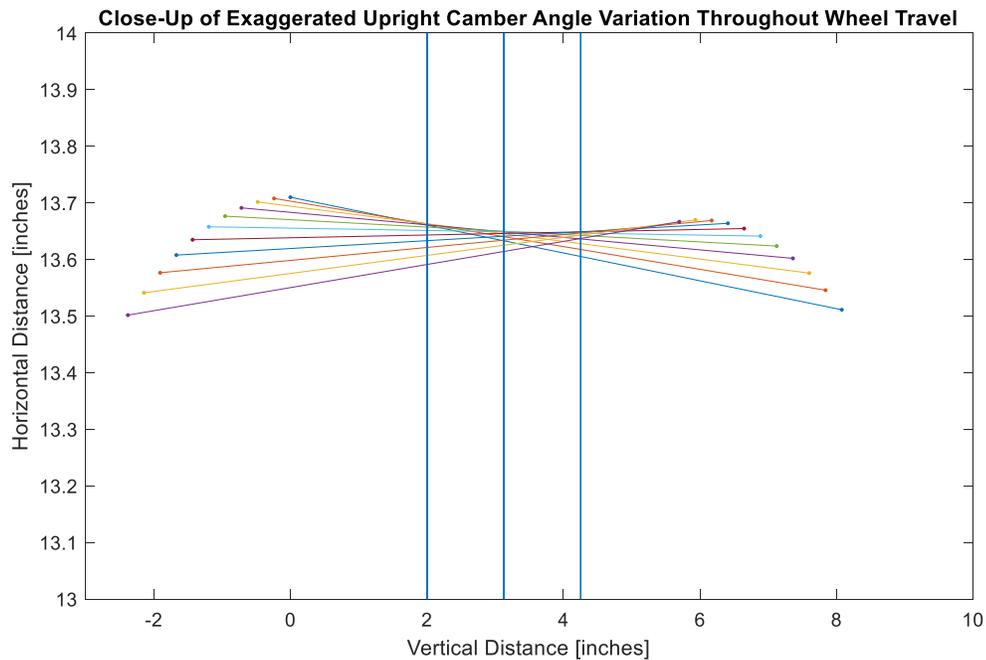


Figure 18: Close-up of the camber variation of the upright through full span of wheel travel jounce and rebound.

This plot isolates the upright linkage motion with unequal scaling in the x and y axes to exaggerate the effects of camber variation causing rotation of the upright. Additionally, this plot has corrected for the difference in inclination angle versus the actual angle of the wheel relative to the ground, so these angles demonstrate the switch between positive and negative camber in wheel travel. One can see clearly the fact that in

When the wheel is being forced towards increasingly negative camber. When cornering, the lateral load applied to the tires causes them to deform as friction between the road and tire pushes the tire towards the center of curvature. However, this friction force causes deformation of the tire and changes the shape of the contact patch. A negative camber on the outside wheel when cornering helps to counteract the loss of contact patch area, such that tire deformations create a fuller contact patch. The full effects of this concept are investigated through empirical testing of tires and plotting the collected data to visualize such relationships. Though Figure 18 helps to show that there exists camber variation with wheel travel, and that it makes a switch by becoming more negatively cambered with jounce travel, we need to know just how much camber there is. With some more calculations in MATLAB, Figure 19 shows exactly that.

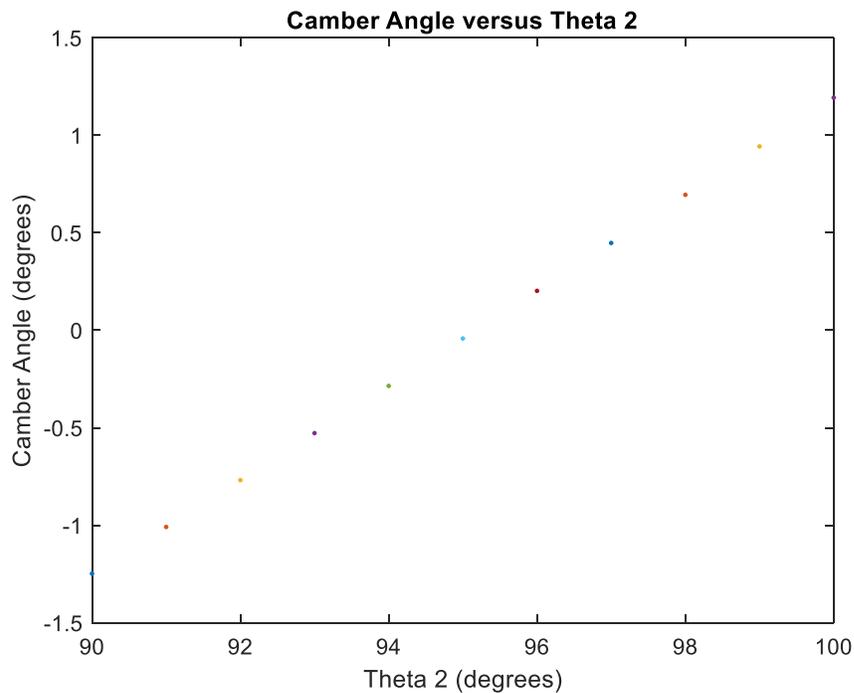


Figure 19: Calculated camber angle of the wheel versus the theta 2 input to the four-bar linkage system, causing wheel travel.

It is evident from this graph that despite the exaggerated visual in Figure 18, there is in fact very little camber variation with wheel travel. This is desired, because more than just a few degrees of camber change can drastically change the way that the tire interacts with the road and responds to various loads. At the equilibrium ride height position, which the wheels will be at when driving in a straight line (without intense launch and braking forces) the best contact patch is made when the tires are perpendicular to the road surface. This is evident here by the camber angle of zero degrees at the middle of wheel travel (when  $\theta_2$  is about 95 degrees). During cornering or launch when the wheels are pushed up into jounce, the camber increases to about 1.25 degrees of negative camber, and similarly in rebound.

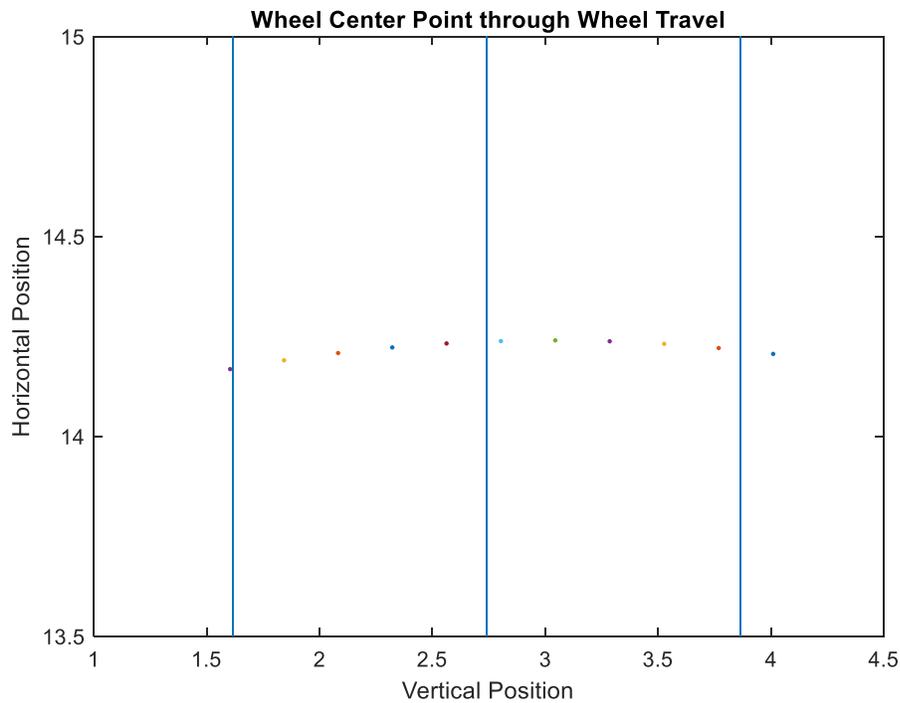


Figure 20: Horizontal displacement of wheel center through vertical wheel travel.

One final consideration is the amount of scrub the tire experiences through wheel travel. Plotted in Figure 20 is the wheel center as it moves horizontally while translating vertically during wheel travel. Scrub is a measure of the lateral distance that the center of the tire translates laterally due to its wheel travel path. More scrub causes more drag and resistance on the tires which takes energy away from the vehicle that could be forward velocity. As shown in Figure 20 there is very little scrub, with maximum displacement less than 0.10 inches, further proving the success of this geometry.

### **Chapter 3: Energy and Behavioral Considerations**

The following three sections utilize the bond graph approach (Karnopp, 2012) to investigate the behavior of a dynamic system. More specifically, the behavior of interest is the motion of certain masses within the system in relation to an input velocity over a wide range of frequencies. This allows the engineer to calculate the range of frequencies which have significant impacts on the desired functionality of the system due to the effects of resonance and damping. A system should be designed to dampen out possible resonance within the frequency range of expected operation. The first section details an introductory analysis to the familiar quarter car model, a system whose behavior is well known within the automotive industry. The second section goes a step further to develop an analysis of a transverse half car model responding to a single-sided velocity input, something not well investigated before. The third section goes even beyond that to investigate the behaviors of the transverse half car model but with an additional third damper mounted transversely as seen in the Swedish Koenigsegg hypercars as a means of improved suspension control and vehicle handling. Such a three-damper half-car system is touted to have better traction and control along straights and resist excessive squat on launch and dive in braking, all while not hindering the functionality of the anti-roll bar in cornering. The developed plots from this mathematical analysis will test this hypothesis for potential applications to an FSAE competition vehicle.

## Conventional Quarter Car Model Analysis

### Schematic.

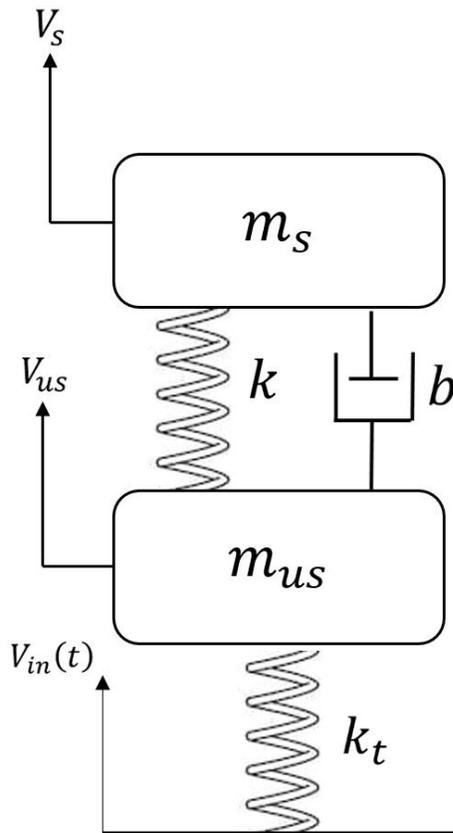


Figure 21: Schematic of quarter car model showing idealized model connected by familiar energy components.

### List of Variables and Their Meaning:

$V_s$  = sprung mass velocity

$m_s$  = sprung mass, representing the portion of mass of the body/chassis of the car distributed to this wheel

$k$  = spring stiffness of the physical helical springs present in the suspension system

$b$  = damping coefficient of the physical damper present in the suspension system

$m_{us}$  = un-sprung mass, includes mass of the wheel, tire, wheel hub, brake rotors and calipers, and control arms.

$V_{us}$  = un-sprung mass velocity

$k_t$  = stiffness coefficient of the tire; the pressurized rubber tire has compliant properties

**Bond Graph.**

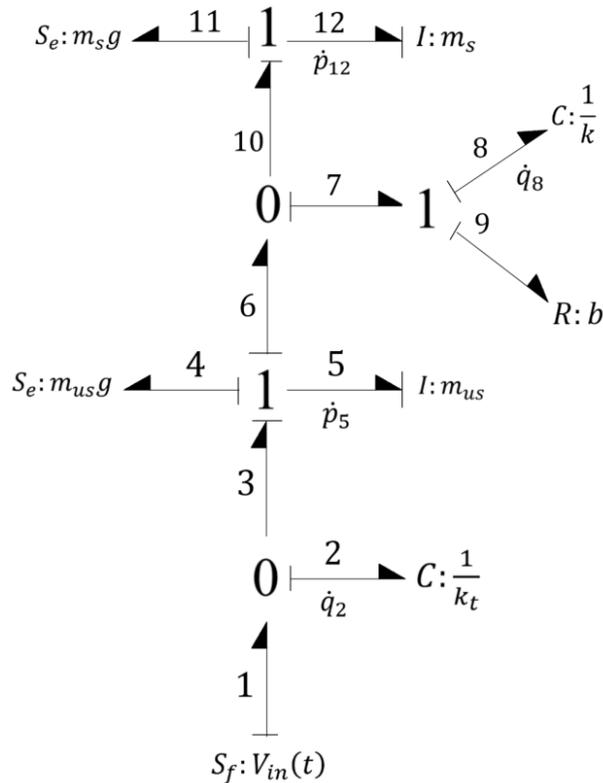


Figure 22: Fully augmented and numbered bond graph of quarter car model.

**Derivation of state space equations and matrix.** Given that the applied causality to the bond graph of the simple quarter car model is all in integral form (the causal stroke indicating that effort is flowing into the inertial elements and out of the capacitive elements) the derivation of the system’s state equations is very straight forward. In this system there exist 4 state variables,  $(q_2, p_5, q_8, p_{12})$  and 3 input variables,  $(V_{in}, m_{us}g, m_s g)$ . It shall be noted that the primary input to this system is  $V_{in}$  which represents the change in velocity caused by the movement of the wheel against a variable road surface which may have bumps, holes, drops, etc. The other input variables come from an analysis including the effect of gravity on both the sprung and un-sprung masses and are in fact not particularly necessary if the assumption is made that the springs in the

system are already at their equilibrium position. However, for ease of intuition in this first quarter car model, they will be kept in the analysis to avoid confusion as these effort sources will be included in the analyses of the following half car models.

$$\dot{q}_2 = V_{in}(t) - \frac{p_5}{m_{us}} \quad (3)$$

$$\dot{p}_5 = q_2 k_t - m_{us} g - q_8 k - b \left( \frac{p_5}{m_{us}} - \frac{p_{12}}{m_s} \right) \quad (4)$$

$$\dot{q}_8 = \frac{p_5}{m_{us}} - \frac{p_{12}}{m_s} \quad (5)$$

$$\dot{p}_{12} = q_8 k + b \left( \frac{p_5}{m_{us}} - \frac{p_{12}}{m_s} \right) - m_s g \quad (6)$$

These state variable equations can then be put into state space matrix form,  $\dot{Y} = [A]\bar{x} + [B]\bar{u}$ , shown as follows:

$$\begin{bmatrix} \dot{q}_2 \\ \dot{q}_8 \\ \dot{p}_5 \\ \dot{p}_{12} \end{bmatrix} = \begin{bmatrix} 0 & 0 & -\frac{1}{m_{us}} & 0 \\ 0 & 0 & \frac{1}{m_{us}} & -\frac{1}{m_s} \\ k & -k & -\frac{b}{m_{us}} & \frac{b}{m_s} \\ 0 & k & \frac{b}{m_{us}} & -\frac{b}{m_s} \end{bmatrix} \begin{bmatrix} q_2 \\ q_8 \\ p_5 \\ p_{12} \end{bmatrix} + \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} V_{in} \\ m_{us} g \\ m_s g \end{bmatrix} \quad (7)$$

With the state space equations in matrix form, manipulation can be done to put the system into the Laplace domain, to effectively solve the system of first order differential equations by algebraic manipulation. This is done by operating on the A matrix, subtracting each element from a 4x4 matrix made by multiplying the Laplacian operator S by the identity matrix of equal size to the A matrix. Essentially, [A] becomes [SI - A], and the rest of the system is put into the Laplace domain, shown by equation 8.

$$\begin{bmatrix} SQ_2(S) \\ SQ_8(S) \\ SP_5(S) \\ SP_{12}(S) \end{bmatrix} = \begin{bmatrix} S & 0 & \frac{1}{m_{us}} & 0 \\ 0 & S & -\frac{1}{m_{us}} & \frac{1}{m_s} \\ -k & k & S + \frac{b}{m_{us}} & -\frac{b}{m_s} \\ 0 & -k & -\frac{b}{m_{us}} & S + \frac{b}{m_s} \end{bmatrix} \begin{bmatrix} q_2(S) \\ q_8(S) \\ p_5(S) \\ p_{12}(S) \end{bmatrix} + \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} V_{in}(S) \\ m_{us}g(S) \\ m_s g(S) \end{bmatrix} \quad (8)$$

From this point, various transfer functions may be found. A transfer function relates a ratio of an output to an input of the system, to see how the behavior is affected. The most insightful transfer function of the quarter car system is the relationship between the velocity of the sprung mass of the vehicle and the input velocity as the wheel receives bumps from the road. This relationship relates directly to driver comfort, making it a desirable system behavior to know. The transfer function is solved by Cramer's rule, taking the quotient of two determinants. The numerator replaces the output variable vector with the input variable vector before solving the determinants. In this case the output variable is the velocity of the sprung mass, which may be solved with the state variable of  $p_{12}$ , the momentum of the sprung mass. It is known  $p = mV$ , so solving by simple algebra the velocity may be found by  $V = \frac{p}{m}$ .

$$\frac{V_{m_s}(S)}{V_{in}(S)} = \frac{1}{m_s} * \frac{\begin{vmatrix} S & 0 & \frac{1}{m_{us}} & 1 \\ 0 & S & -\frac{1}{m_{us}} & 0 \\ -k & k & S + \frac{b}{m_{us}} & 0 \\ 0 & -k & -\frac{b}{m_{us}} & 0 \end{vmatrix}}{\begin{vmatrix} S & 0 & \frac{1}{m_{us}} & 0 \\ 0 & S & -\frac{1}{m_{us}} & \frac{1}{m_s} \\ -k & k & S + \frac{b}{m_{us}} & -\frac{b}{m_s} \\ 0 & -k & -\frac{b}{m_{us}} & S + \frac{b}{m_s} \end{vmatrix}} \quad (9)$$

$$\frac{V_{m_s}}{V_{in}}(S) = \frac{k^2 + S b k}{(m_{us}S^4 + b S^3 + 2 k S^2)m_s + b m_{us}S^3 + m_{us}S^2k + b S k + k^2} \quad (10)$$

With a transfer function between the input velocity and the velocity of the sprung mass, various analyses may be done on the system to see how it behaves with specific given parameters. Results of such an analyses are demonstrated by Figure 23 with plots for the amplitude ratio of the sprung mass to the input velocity and the phase angle plotted along an increasing frequency. However, it is necessary first to list the values assigned to the parameters in the equation to give context and meaning to the plots. These are found in Table 2.

Table 2: Values assigned to parameters in quarter car equations used to create plots of frequency response behavior.

Parameter	Value	Units
$m_{us}$	8.0	kg
$k$	4000	Nm/deg
$k_t$	123464	N/m
$m_s$	182/4 *for 400lb vehicle	kg
$b$	$0.6(2)m_s \sqrt{\frac{k}{m_s}}$	Ns/m

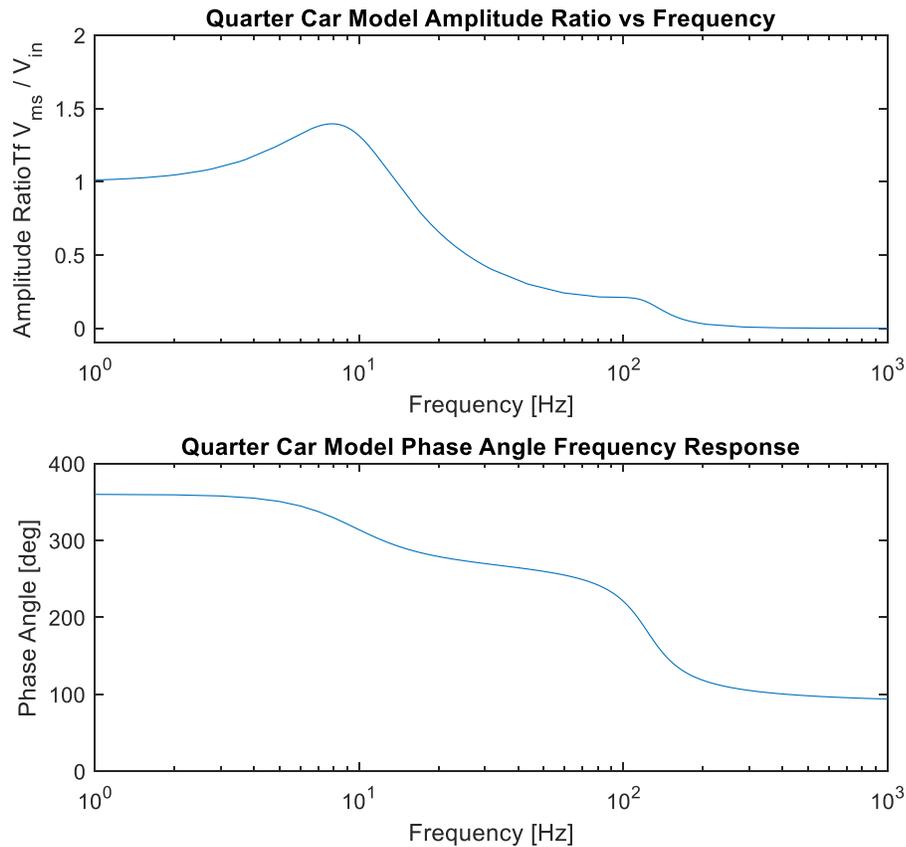


Figure 23: Quarter car model frequency response plots of amplitude ratio and phase angle.

The goal of the quarter car model is to essentially isolate the input velocity from the sprung mass velocity with a damped system, such that the sprung mass returns to a condition of steady state with a comfortable number of oscillations, as energy is dissipated by the dampers. It was originally thought that the system should be critically damped, as there is a very specific relationship between the mass, spring constant, and damping coefficient that will permit the system to be critically damped of  $b_c = 2m\sqrt{\frac{k}{m}}$ . However, upon further research, it is determined that for FSAE applications, and even in most racing applications the goal is not to control oscillations such that only one complete oscillation is made before the sprung mass returns to steady state. In fact, it is

more desired that the damping ratio  $\frac{b}{b_c}$  be somewhere around 0.5 to 0.7 for the best road feel and vehicle control. With a damping ratio in this range, there is not such drastic lateral load transfer in cornering, so the inside wheel on the turn is allowed to keep more vertical load on it which helps to increase grip and overall cornering capacity. This was taken into consideration when applying values to the MATLAB calculations for the plots in this report. It shall be acknowledged that most performance dampers are designed to have different damping coefficients in compression and rebound. However, for the purposes of analysis one can simply consider multiple plots of various damping ratios and evaluate the curves correspondingly.

Despite all this, the quarter car model is very well known as automotive companies seek to outfit their vehicles with the best possible comfort to increase sales and satisfaction. What is less well known, and has not been analyzed much before with the bond graph approach is the relationship between the two lateral halves of the vehicle's suspension, meaning the right and left side of the car, as opposed to front to back half. Longitudinal models have been studied to investigate the behaviors of dive and squat. Lateral suspension behavior was investigated here in the conventional and Koenigsegg half car models in the following sections.

Though the primary focus of this investigation is on lateral behavior in a half car model, it is necessary to also introduce the effects of dive and squat, as a third damper system serves to counteract the negative effects of these behaviors. Dive pulls the nose of the car down when braking as load transfer due to negative acceleration (slowing down) causes spring compression in the front end, requiring the front wheels to handle the

majority of the braking power. This is never a beneficial quality; maximum braking power is obtained when all four wheels are loaded equally. Squat pulls the rear of the car down when accelerating, as load transfer due to positive acceleration (launch, or increasing speed) causes additional spring compression in the rear. This is sometimes a desired behavior as rear-wheel drive vehicles can transmit more power during launch with a greater normal force.

### Anti-Squat Control Arm Geometry

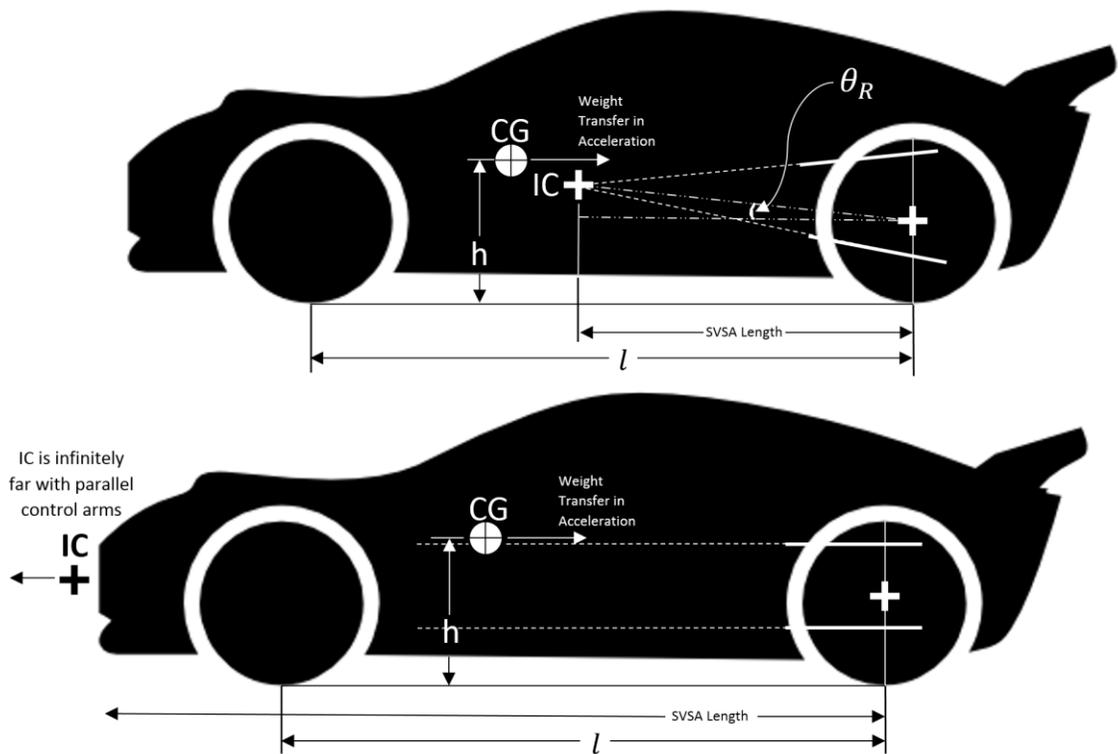


Figure 24: Anti-Squat control arm geometry (top) versus no anti-squat geometry with parallel control arms (bottom).  $l$  is the length of the wheelbase, and  $h$  is the height of the center of gravity (CG).

Both dive and squat are not particularly desired in racing applications, as they tend to have negative effects on the contact patch between the tire and the road due to changes in camber angle in straight-line driving, as well as changes in aerodynamic

influence and ride height. Changes in the orientation of the control arms can alter the percent of anti-dive or anti-squat geometry the vehicle has, by putting longitudinal load transfer more into tensile and compressive forces acting on the control arms and linkages rather than entirely into the springs and dampers. Figure 24 helps to visualize the suspension orientation changes made to counteract squat. The percentage of anti-squat made by the angle of the control arms can be calculated as follows with reference to Figure 24.

$$\%Anti - Squat = \frac{\tan(\theta_R)}{\frac{h}{l}} * 100 \quad (11)$$

The top vehicle depicts anti-squat geometry of about 52%, assuming  $\theta_R = 5deg$ ,  $h = 11.0in$ , and  $l = 65.0in$ . The bottom vehicle in the Figure demonstrates 0% anti-squat with both control arms made parallel to the ground. This effectively puts the instantaneous center (IC) infinitely far away so  $\theta_R$  goes to zero, and the numerator of equation 11 is then zero.

However, though it is beneficial to distribute some of the load out of the springs to reduce spring compression and therefore total squat/dive, the distribution of these loads demand that the control arms be made stronger to withstand greater stresses. Assuming the same material is used for the control arms, greater strength requires more material, resulting in more un-sprung mass.

Everything not directly contained by the chassis is considered “sprung mass,” (because it is suspended by the springs). Everything else, including the wheels, tires, wheel hubs, brake rotors, calipers, and control arms are considered “un-sprung mass.” It

is imperative for the benefits of a vehicle's frequency response and overall dynamics to have the least possible un-sprung mass; this means there is less inertia and momentum related to wheel movements with irregularities of the road surface, resulting in less influence of residual motion into the sprung mass. When the control arms are forced to be more massive to handle increased stresses by directing squat and dive forces into them, it increases the un-sprung mass and makes the system more difficult to control.

A third damper system serves to minimize the magnitude of squat and dive experienced during longitudinal load transfer by dissipating some of this energy into the hydraulic fluid of the damper. When two wheels both move up at the same time very rapidly (as in the case of hard braking from high speed) and a damper is connected transversely between them, it is compressed from both sides. This movement forces the piston to move rapidly within its hydraulic fluid, transferring energy into heat that would otherwise go into compression or extension of the springs.

### **Conventional Transverse Half Car Model Analysis**

The transverse half car model investigates the concepts of what happens to body roll when a single wheel experiences a velocity change input. Traditionally, a component known as the "sway bar" or "anti-roll bar" is installed that connects two wheels transversely by way of a U-shaped bar that elastically deforms in torsion to act as a spring. This is done such that when the body experiences roll, and the outer wheels travel up as their springs and dampers compress, it will cause torsion on the sway bar that will also pull up on the other wheel, effectively working to help balance out the effects of roll in a turn to maintain the best contact patch possible between the tires and the road.

**Schematic.**

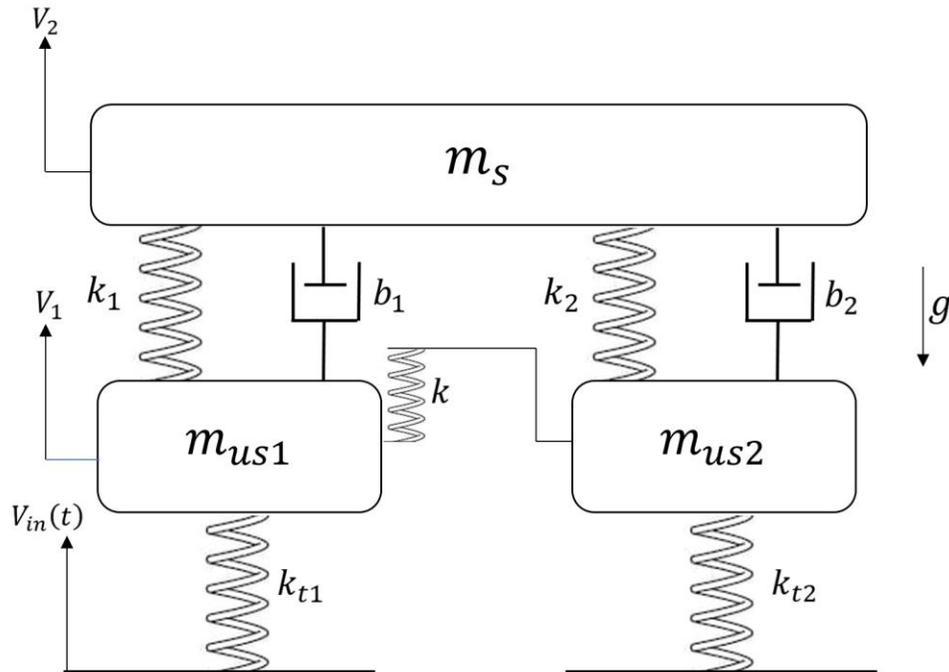


Figure 25: Schematic of conventional half car model showing idealized model connected by familiar energy components.

**Bond Graph.**

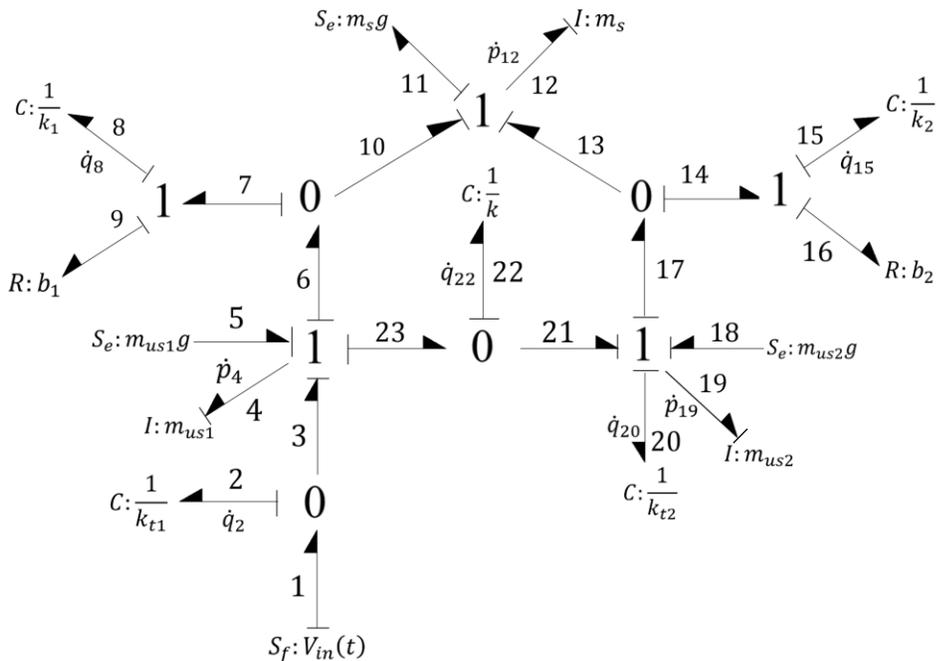


Figure 26: Bond graph of conventional half car model fully augmented and numbered.

**Derivation of state space equations and matrix.** The following 8 equations list the final equations derived from following the analysis of the bond graph given the assigned augmentation with compression considered positive. The hand calculations for these equations may be found in Appendix D.

$$\dot{q}_2 = V_{in}(t) - \frac{p_4}{m_{us1}} \quad (12)$$

$$\dot{q}_8 = \frac{p_4}{m_{us1}} - \frac{p_{12}}{m_s} \quad (13)$$

$$\dot{q}_{15} = \frac{p_{19}}{m_{us2}} - \frac{p_{12}}{m_s} \quad (14)$$

$$\dot{q}_{20} = \frac{p_{19}}{m_{us2}} \quad (15)$$

$$\dot{q}_{22} = \frac{p_4}{m_{us1}} - \frac{p_{19}}{m_{us2}} \quad (16)$$

$$\dot{p}_4 = q_2 k_{t1} + m_{us1} g - q_8 k_1 - b_1 \left( \frac{p_4}{m_{us1}} - \frac{p_{12}}{m_s} \right) - q_{22} k \quad (17)$$

$$\dot{p}_{12} = -m_s g + q_8 k_1 + b_1 \left( \frac{p_4}{m_{us1}} - \frac{p_{12}}{m_s} \right) + q_{15} k_2 + b_2 \left( \frac{p_{19}}{m_{us2}} - \frac{p_{12}}{m_s} \right) \quad (18)$$

$$\dot{p}_{19} = q_{22} k + m_{us2} g - q_{15} k_2 - b_2 \left( \frac{p_{19}}{m_{us2}} - \frac{p_{12}}{m_s} \right) - q_{20} k_{t2} \quad (19)$$

To work with these equations requires that they be put into state-space matrix form. They are listed in the order that they are to be arranged in the matrix shown by equation 20.

$$\begin{bmatrix} \dot{q}_2 \\ \dot{q}_8 \\ \dot{q}_{15} \\ \dot{q}_{20} \\ \dot{q}_{22} \\ \dot{p}_4 \\ \dot{p}_{12} \\ \dot{p}_{19} \end{bmatrix} = \begin{bmatrix} 0 & 0 & 0 & 0 & 0 & -\frac{1}{m_{us1}} & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & \frac{1}{m_{us1}} & -\frac{1}{m_s} & -\frac{1}{m_s} \\ 0 & 0 & 0 & 0 & 0 & 0 & -\frac{1}{m_s} & \frac{1}{m_{us2}} \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & \frac{1}{m_{us2}} \\ 0 & 0 & 0 & 0 & 0 & \frac{1}{m_{us1}} & 0 & -\frac{1}{m_{us2}} \\ k_{t1} & -k_1 & 0 & 0 & -k & -\frac{b_1}{m_{us1}} & \frac{b_1}{m_s} & 0 \\ 0 & k_1 & k_2 & 0 & 0 & \frac{b_1}{m_{us1}} & \left( -\frac{b_1 - b_2}{m_s} \right) & \frac{b_2}{m_{us2}} \\ 0 & 0 & -k_2 & -k_{t2} & k & 0 & \frac{b_2}{m_s} & -\frac{b_2}{m_{us2}} \end{bmatrix} \begin{bmatrix} q_2 \\ q_8 \\ q_{15} \\ q_{20} \\ q_{22} \\ p_4 \\ p_{12} \\ p_{19} \end{bmatrix} + \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} V_{in} \\ m_{us1} g \\ m_{us2} g \\ m_s g \end{bmatrix} \quad (20)$$

Though this matrix in equation 20 is substantially larger than the example demonstrated by the quarter car model, the method of analysis is the same. We then converted to the Laplacian domain, so  $[A]$  becomes  $[SI - A]$ .

$$\begin{bmatrix} SQ_2(S) \\ SQ_8(S) \\ SQ_{15}(S) \\ SQ_{20}(S) \\ SQ_{22}(S) \\ SP_4(S) \\ SP_{12}(S) \\ SP_{19}(S) \end{bmatrix} = \begin{bmatrix} S & 0 & 0 & 0 & 0 & \frac{1}{m_{us1}} & 0 & 0 \\ 0 & S & 0 & 0 & 0 & -\frac{1}{m_{us1}} & \frac{1}{m_s} & \frac{1}{m_s} \\ 0 & 0 & S & 0 & 0 & 0 & \frac{1}{m_s} & -\frac{1}{m_{us2}} \\ 0 & 0 & 0 & S & 0 & 0 & 0 & -\frac{1}{m_{us2}} \\ 0 & 0 & 0 & 0 & S & S - \frac{1}{m_{us1}} & 0 & \frac{1}{m_{us2}} \\ -k_{t1} & k_1 & 0 & 0 & k & \frac{b_1}{m_{us1}} & S - \frac{b_1}{m_s} & 0 \\ 0 & -k_1 & -k_2 & 0 & 0 & -\frac{b_1}{m_{us1}} & S + \left(\frac{b_1 - b_2}{m_s}\right) & -\frac{b_2}{m_{us2}} \\ 0 & 0 & k_2 & k_{t2} & -k & 0 & -\frac{b_2}{m_s} & S + \frac{b_2}{m_{us2}} \end{bmatrix} \begin{bmatrix} Q_2(S) \\ Q_8(S) \\ Q_{15}(S) \\ Q_{20}(S) \\ Q_{22}(S) \\ P_4(S) \\ P_{12}(S) \\ P_{19}(S) \end{bmatrix} + \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} V_{in}(S) \\ m_{us1}g(S) \\ m_{us2}g(S) \\ m_s g(S) \end{bmatrix} \quad (21)$$

From this point, the transfer functions of interest may be derived. The interest here was in determining the relationships between the input velocity and each wheel in the system, as one wheel is directly affected by the input velocity to its tire, while the other feels the input velocity only as it relates through the sprung mass and anti-roll bar. This way, we can viably compare the behavior of each wheel responding to the same input, with the only difference being the addition of the third damper. Therefore, I will derive the transfer functions of  $\frac{V_{m_{us1}}}{V_{in}}(S)$  and  $\frac{V_{m_{us2}}}{V_{in}}(S)$ , utilizing the Cramer's rule replacement and knowing fundamentally momentum is equal to mass times volume ( $p = mV$ ) as demonstrated by the quarter car model example.

The first of these transfer functions was set up as shown in equation 22, however for conservation of space the writing out of this step for subsequent transfer functions was omitted, skipping directly to the result of the quotient of the determinants stated by

Cramer's rule. However, it shall be noticed in this instance that the 6<sup>th</sup> column of the matrix in the numerator corresponding to the momentum of  $m_{us1}$  is replaced by the vector corresponding to the input variable  $V_{in}$  as dictated by the method of Cramer's rule. Then the quotient was divided by  $\frac{1}{m_{us1}}$  to effectively convert the momentum of the first un-sprung mass to its velocity to match the need of the transfer function.

$$TF_1 = \frac{V_{m_{us1}}}{V_{in}}(S) = \frac{1}{m_{us1}} * \begin{array}{c} \left| \begin{array}{ccccccccc} S & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & S & 0 & 0 & 0 & 0 & \frac{1}{m_s} & \frac{1}{m_s} & 0 \\ 0 & 0 & S & 0 & 0 & 0 & \frac{1}{m_s} & -\frac{1}{m_{us2}} & 0 \\ 0 & 0 & 0 & S & 0 & 0 & 0 & -\frac{1}{m_{us2}} & 0 \\ 0 & 0 & 0 & 0 & S & 0 & 0 & \frac{1}{m_{us2}} & 0 \\ -k_{t1} & k_1 & 0 & 0 & k & 0 & S - \frac{b_1}{m_s} & 0 & 0 \\ 0 & -k_1 & -k_2 & 0 & 0 & 0 & S + \left(\frac{b_1 - b_2}{m_s}\right) & -\frac{b_2}{m_{us2}} & 0 \\ 0 & 0 & k_2 & k_{t2} & -k & 0 & -\frac{b_2}{m_s} & S + \frac{b_2}{m_{us2}} & 0 \end{array} \right| \\ \hline \left| \begin{array}{ccccccccc} S & 0 & 0 & 0 & 0 & \frac{1}{m_{us1}} & 0 & 0 & 0 \\ 0 & S & 0 & 0 & 0 & -\frac{1}{m_{us1}} & \frac{1}{m_s} & \frac{1}{m_s} & 0 \\ 0 & 0 & S & 0 & 0 & 0 & \frac{1}{m_s} & -\frac{1}{m_{us2}} & 0 \\ 0 & 0 & 0 & S & 0 & 0 & 0 & -\frac{1}{m_{us2}} & 0 \\ 0 & 0 & 0 & 0 & S & S - \frac{1}{m_{us1}} & 0 & \frac{1}{m_{us2}} & 0 \\ -k_{t1} & k_1 & 0 & 0 & k & \frac{b_1}{m_{us1}} & S - \frac{b_1}{m_s} & 0 & 0 \\ 0 & -k_1 & -k_2 & 0 & 0 & -\frac{b_1}{m_{us1}} & S + \left(\frac{b_1 - b_2}{m_s}\right) & -\frac{b_2}{m_{us2}} & 0 \\ 0 & 0 & k_2 & k_{t2} & -k & 0 & -\frac{b_2}{m_s} & S + \frac{b_2}{m_{us2}} & 0 \end{array} \right| \end{array} \quad (22)$$

This equation was solved in MATLAB as it would be extremely tedious and impractical to solve by hand. The code for this and the following determinant calculations may be found in Appendix C. The result even simplified is very lengthy, and will be separated into numerator and denominator portions for clarity.

$$\begin{aligned}
 TF_{1NUM} = & \left( (S^3 b_2 k_{t1} + S^2 k k_{t1} + S^2 k_2 k_{t1} + S^2 k_{t1} k_{t2} + S^4 k_{t1} m_{us2}) m_s^2 \right. \\
 & + (k k_1 k_{t1} - 2 S^2 b_2^2 k_{t1} + k k_2 k_{t1} + k_1 k_2 k_{t1} + k_1 k_{t1} k_{t2} + k_2 k_{t1} k_{t2} \\
 & + S^2 b_1 b_2 k_{t1} + S^3 b_1 k_{t1} m_{us2} - S^3 b_2 k_{t1} m_{us2} + S^2 k_1 k_{t1} m_{us2} + S^2 k_2 k_{t1} m_{us2} \\
 & + S b_1 k k_{t1} - S b_2 k k_{t1} + S b_1 k_2 k_{t1} + S b_2 k_1 k_{t1} - 2 S b_2 k_2 k_{t1} + S b_1 k_{t1} k_{t2} \\
 & \left. - S b_2 k_{t1} k_{t2}) m_s + k_1 k_2 k_{t1} m_{us2} + S b_2 k_1 k_{t1} m_{us2} \right)
 \end{aligned}$$

$$\begin{aligned}
 TF_{1DENOM} = & \left( (S^4 b_1 b_2 + S^3 b_1 k + S^3 b_2 k + S^3 b_1 k_2 + S^3 b_2 k_1 + S^3 b_1 k_{t2} + S^3 b_2 k_{t1} \right. \\
 & + S^5 b_1 m_{us2} + S^5 b_2 m_{us1} + S^2 k k_1 + S^2 k k_2 + S^2 k k_{t1} + S^2 k k_{t2} + S^2 k_1 k_2 \\
 & + S^2 k_1 k_{t2} + S^2 k_2 k_{t1} + S^2 k_{t1} k_{t2} + S^4 k m_{us1} + S^4 k m_{us2} + S^4 k_1 m_{us2} \\
 & + S^4 k_2 m_{us1} + S^4 k_{t1} m_{us2} + S^4 k_{t2} m_{us1} + S^6 m_{us1} m_{us2}) m_s^2 \\
 & + (k k_1 k_{t1} - 2 S^2 b_2^2 k - 2 S^2 b_2^2 k_1 - 2 S^2 b_2^2 k_{t1} - 2 S^4 b_2^2 m_{us1} - 2 S^3 b_1 b_2^2 \\
 & + k k_1 k_{t2} + k k_2 k_{t1} + k k_2 k_{t2} + k_1 k_2 k_{t1} + k_1 k_2 k_{t2} + k_1 k_{t1} k_{t2} + k_2 k_{t1} k_{t2} \\
 & - 2 S^2 b_1 b_2 k - 2 S^2 b_1 b_2 k_2 + S^2 b_1 b_2 k_{t1} - S^2 b_1 b_2 k_{t2} + S^4 b_1 b_2 m_{us1} \\
 & - S^4 b_1 b_2 m_{us2} + S^3 b_1 k m_{us1} + S^3 b_1 k m_{us2} - S^3 b_2 k m_{us1} - S^3 b_2 k m_{us2} \\
 & + S^3 b_1 k_2 m_{us1} + S^3 b_2 k_1 m_{us1} + S^3 b_1 k_2 m_{us2} - S^3 b_2 k_1 m_{us2} - 2 S^3 b_2 k_2 m_{us1} \\
 & + S^3 b_1 k_{t1} m_{us2} + S^3 b_1 k_{t2} m_{us1} - S^3 b_2 k_{t1} m_{us2} - S^3 b_2 k_{t2} m_{us1} \\
 & + S^5 b_1 m_{us1} m_{us2} - S^5 b_2 m_{us1} m_{us2} + S^2 k k_1 m_{us1} + S^2 k k_2 m_{us1} \\
 & + S^2 k k_2 m_{us2} + S^2 k_1 k_2 m_{us1} + S^2 k_1 k_2 m_{us2} + S^2 k_1 k_{t1} m_{us2} + S^2 k_1 k_{t2} m_{us1} \\
 & + S^2 k_2 k_{t1} m_{us2} + S^2 k_2 k_{t2} m_{us1} + S^4 k_1 m_{us1} m_{us2} + S^4 k_2 m_{us1} m_{us2} \\
 & - 2 S b_2 k k_1 - 2 S b_2 k k_2 + S b_1 k k_{t1} + S b_1 k k_{t2} - S b_2 k k_{t1} - S b_2 k k_{t2} \\
 & - 2 S b_2 k_1 k_2 + S b_1 k_2 k_{t1} + S b_2 k_1 k_{t1} + S b_1 k_2 k_{t2} - S b_2 k_1 k_{t2} - 2 S b_2 k_2 k_{t1} \\
 & + S b_1 k_{t1} k_{t2} - S b_2 k_{t1} k_{t2}) m_s + k_1 k_2 k_{t1} m_{us2} + 2 S b_2 k k_1 m_{us2} \\
 & \left. + S b_2 k_1 k_{t1} m_{us2} + S^3 b_2 k_1 m_{us1} m_{us2} + S^2 k_1 k_2 m_{us1} m_{us2} \right)
 \end{aligned}$$

It is evident that this system is of great complexity with a matrix of 8 state variables, so the following mathematical results will not be displayed within the body of this report, but rather will be handled internally within MATLAB, only plotting the insightful behavioral results from a frequency response analysis on this system. The next step in the analysis is to replace the Laplacian operator  $S$  with  $j\omega$  to simulate a harmonic

input, ultimately permitting an analysis of the frequency response and amplitude ratio of the system. There exists a convenient function within MATLAB that can take the coefficients of the S values of a transfer function system to plot both the frequency response and phase angle, but given the length and complexity of this system, it is easier to make these calculations “by hand” within the MATLAB code as reflected in Appendix C, rather than manually searching through to find the coefficients of each degree of S. Once these calculations are made, plots are created according to Figure 27. For relevance of these plots, they were made with parameter values according to Table 3.

Table 3: Values assigned to parameters in half car equations used to create plots of frequency response behavior.

Parameter	Value	Units
$m_s$	200	kg
$m_{us1}$	8.0	kg
$m_{us2}$	8.0	kg
$k_1$	4000	N/m
$k_2$	4000	N/m
$k$	50000	Nm/deg
$k_{t1}$	123464	N/m
$k_{t2}$	123464	N/m
$m_s$	182/4 *for 400lb vehicle	kg
$b_1$	$\frac{0.6(2)m_s}{4} \sqrt{\frac{k_1}{\left(\frac{m_s}{4}\right)}}$	Ns/m
$b_2$	$\frac{0.6(2)m_s}{4} \sqrt{\frac{k_1}{\left(\frac{m_s}{4}\right)}}$	Ns/m

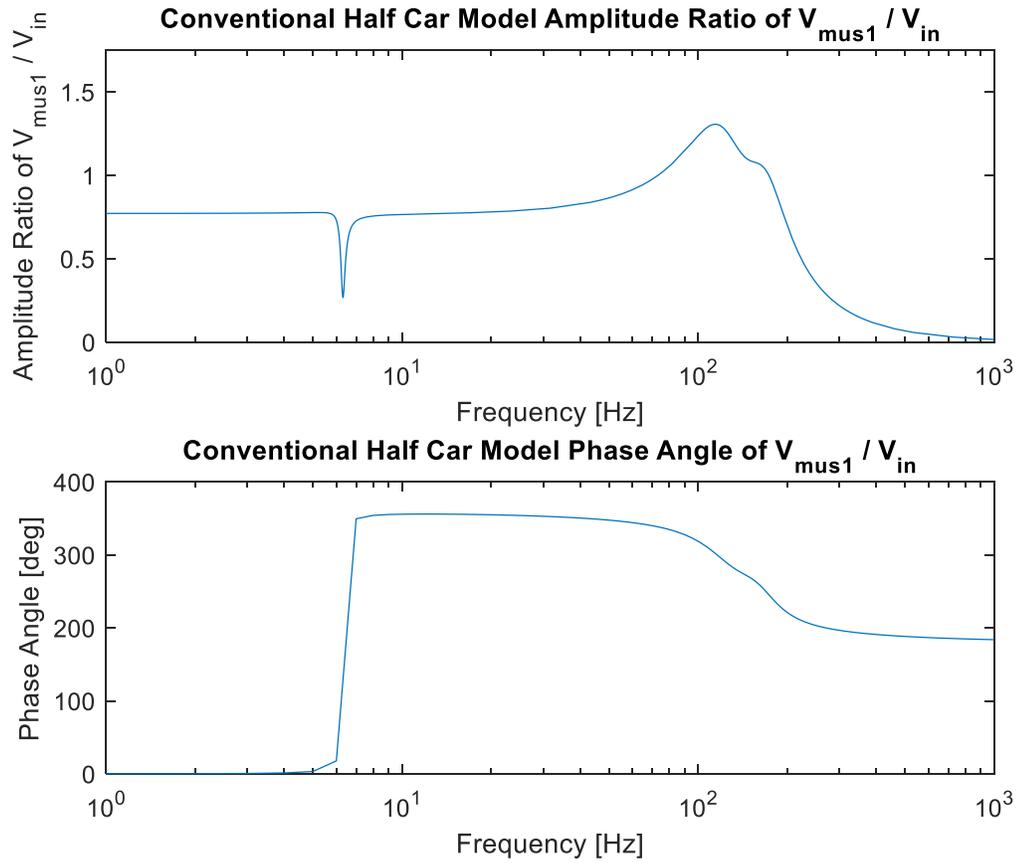


Figure 27: Frequency response plots of the left un-sprung mass velocity to the input velocity in the conventional half-car model.

By a similar process to equation 22, the transfer functions for  $\frac{V_{ms}}{V_{in}}(S)$  and  $\frac{V_{mus2}}{V_{in}}(S)$  were calculated, converted to a harmonic input for frequency response analysis, and the amplitude ratio and phase angles were plotted and calculated. The resulting plots of these calculations are given below.

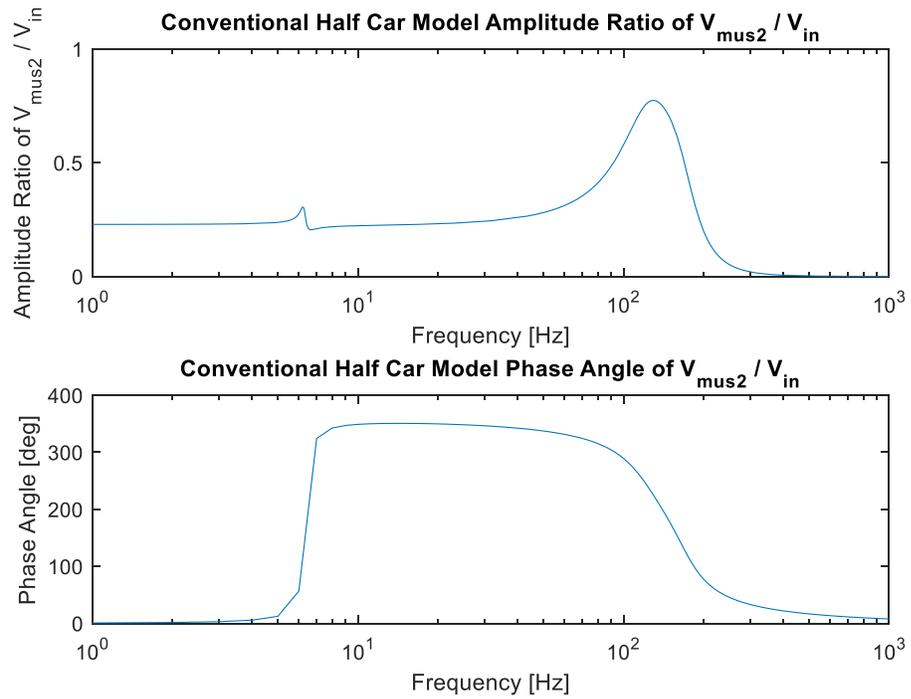


Figure 28: Frequency response plots of the right un-sprung mass velocity to the input velocity in the conventional half-car model.

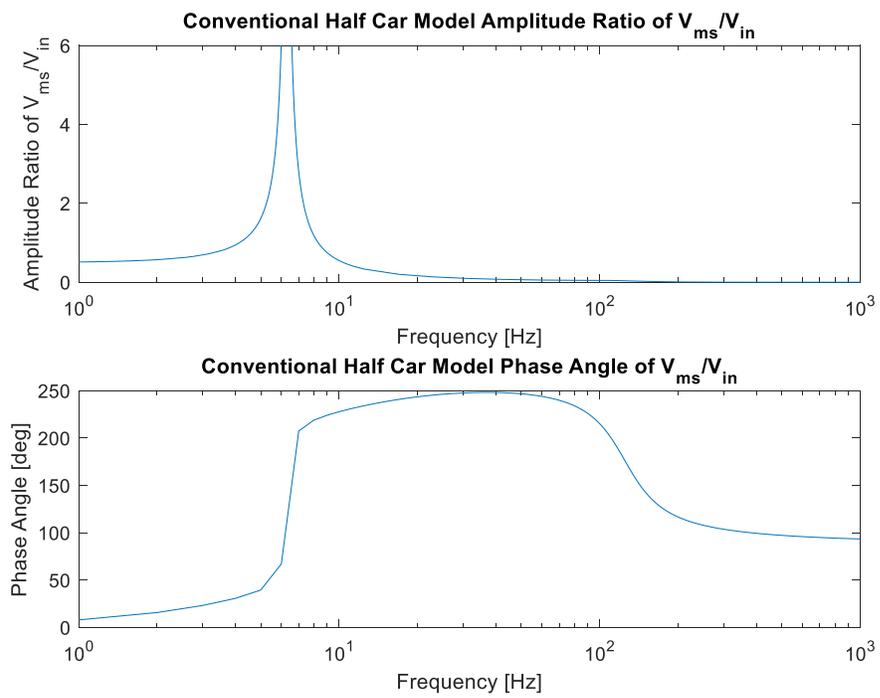


Figure 29: Frequency response plots of the sprung mass velocity to the input velocity in the conventional half-car model.

Now for the most important plot of this conventional model, is to overlay the frequency response amplitude ratio between the left and right un-sprung masses. This will clearly show the difference between the behavior of the two wheels on either side of the car when only one wheel experiences a bump or change in velocity while the other continues on flat ground, but is forced to move because of the sway bar connection between the two wheels.

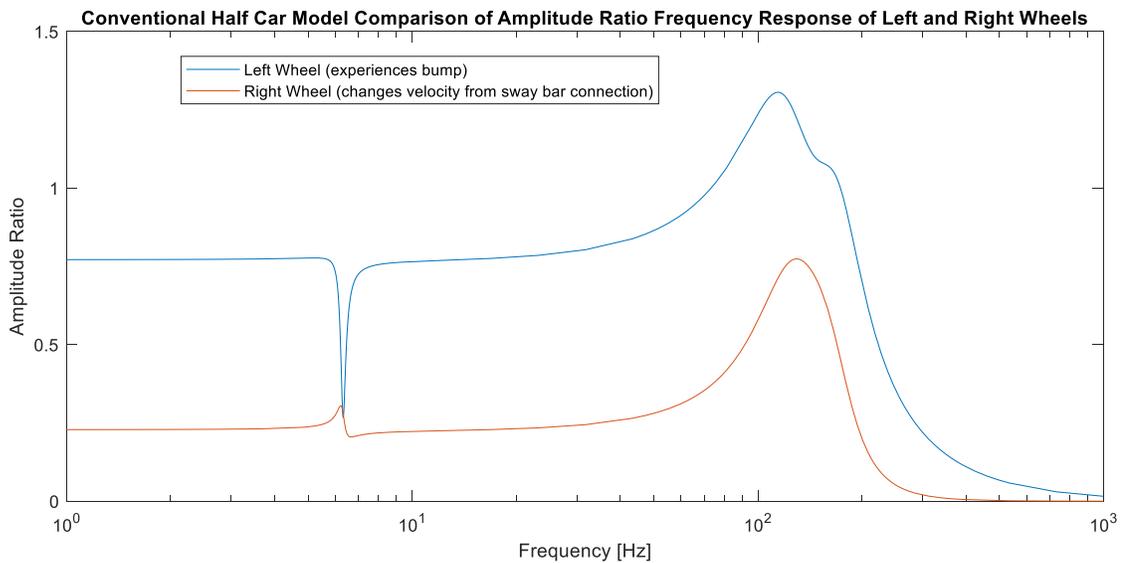


Figure 30: Visual comparison of behavior between left and right wheels only connected by the sway bar.

### Introduction to Koenigsegg Design

Christian von Koenigsegg, a Swedish hypercar manufacturer sought to innovate at every opportunity possible. One of these brilliant innovations is his famous Triplex suspension, which utilizes the addition of a third damper mounted transversely between the two upper wishbone control arms of wheels on either side of the car. With this addition and many others, he has successfully made some of the uncontested best racing performance vehicles ever manufactured, winning all sorts of records and firsts.

**Koenigsegg Design Transverse Half-Car Model Analysis**

**Schematic.**

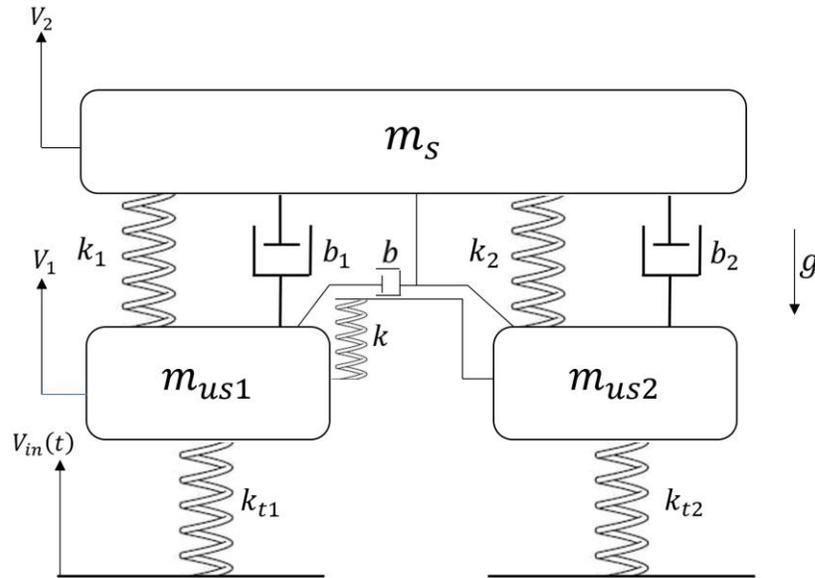


Figure 31: Schematic of innovative half car model showing idealized model connected by familiar energy components.

**Bond Graph.**

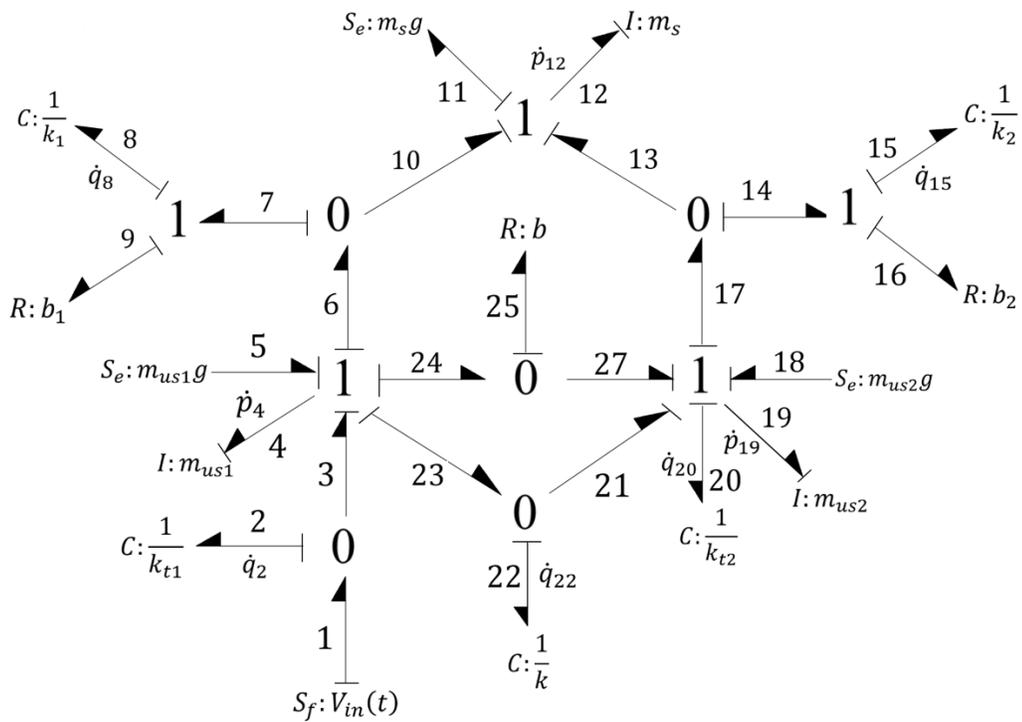


Figure 32: Bond graph of innovative half car model fully augmented and numbered.

**Derivation of state space equations and matrix.** The derivation of the state space equations was quite similar to those derived for the conventional transverse half car model being that the bond graphs are so similar and maintain integral causality. Very conveniently, the derivations of all of the  $\dot{q}$  equations were exactly the same, so I will refer you back to the previous section's derivations for those expressed equations. However, all three of the  $\dot{p}$  equations are slightly different with the addition of the third damper to the system. This is intuitively expected, given that additional damping naturally changes the momenta of masses within a system when acted upon by an input velocity. These new equations are derived as follows:

$$\dot{p}_4 = q_2(k_{t1}) + q_8(-k_1) + q_{22}(-k) + p_4 \left( \frac{-b_1 - b}{m_{us1}} \right) + p_{12} \left( \frac{b_1}{m_s} \right) + p_{19} \left( -\frac{b}{m_{us2}} \right) + m_{us1}g \quad (24)$$

$$\dot{p}_{12} = q_8(k_1) + q_{15}(k_2) + p_4 \left( \frac{b_1}{m_{us1}} \right) + p_{12} \left( \frac{-b_1 - b_2}{m_s} \right) + p_{19} \left( \frac{b_2}{m_{us2}} \right) - m_s g \quad (25)$$

$$\dot{p}_{19} = q_{15}(-k_2) + q_{20}(-k_{t2}) + q_{22}(k) + p_4 \left( -\frac{b}{m_{us1}} \right) + p_{12} \left( \frac{b_2}{m_s} \right) + p_{19} \left( \frac{-b - b_2}{m_{us2}} \right) + m_{us2}g \quad (26)$$

Since these state variable equations are new, the  $[A]$  matrix of the state space matrix equation was re-created, as shown with the vector replacements from equations 24 - 26.

$$\begin{bmatrix} \dot{q}_2 \\ \dot{q}_8 \\ \dot{q}_{15} \\ \dot{q}_{20} \\ \dot{q}_{22} \\ \dot{p}_4 \\ \dot{p}_{12} \\ \dot{p}_{19} \end{bmatrix} = \begin{bmatrix} 0 & 0 & 0 & 0 & 0 & -\frac{1}{m_{us1}} & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & \frac{1}{m_{us1}} & -\frac{1}{m_s} & -\frac{1}{m_s} \\ 0 & 0 & 0 & 0 & 0 & 0 & -\frac{1}{m_s} & \frac{1}{m_{us2}} \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & \frac{1}{m_{us2}} \\ 0 & 0 & 0 & 0 & 0 & \frac{1}{m_{us1}} & 0 & -\frac{1}{m_{us2}} \\ k_{t1} & -k_1 & 0 & 0 & -k & \frac{-b_1 - b}{m_{us1}} & \frac{b_1}{m_s} & -\frac{b}{m_{us2}} \\ 0 & k_1 & k_2 & 0 & 0 & \frac{b_1}{m_{us1}} & \left( \frac{-b_1 - b_2}{m_s} \right) & \frac{b_2}{m_{us2}} \\ 0 & 0 & -k_2 & -k_{t2} & k & -\frac{b}{m_{us1}} & \frac{b_2}{m_s} & -\frac{-b - b_2}{m_{us2}} \end{bmatrix} \begin{bmatrix} q_2 \\ q_8 \\ q_{15} \\ q_{20} \\ q_{22} \\ p_4 \\ p_{12} \\ p_{19} \end{bmatrix} + \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} V_{in} \\ m_{us1}g \\ m_{us2}g \\ m_s g \end{bmatrix} \quad (27)$$

If one were to compare the two A matrices, it is apparent that the only differences appear in the 9 entries of the lower right corner of the matrix, corresponding to the momentum variables. From this point, the same mathematical process as was done in the conventional half car analysis was carried out. The plots of the behavior of the system based on the same three transfer functions as derived previously, are shown in Figure 33 through Figure 36. For plot relevance, the values assigned to the equation parameters are given in Table 4.

Table 4: Values assigned to parameters in Koenigsegg equations used to create plots of frequency response behavior.

Parameter	Value	Units
$m_s$	200	kg
$m_{us1}$	8.0	kg
$m_{us2}$	8.0	kg
$k_1$	4000	N/m
$k_2$	4000	N/m
$k$	50000	Nm/deg
$k_{t1}$	123464	N/m
$k_{t2}$	123464	N/m
$m_s$	182/4 *for 400lb vehicle	kg
$b_1$	$\frac{0.6(2)m_s}{4} \sqrt{\frac{k_1}{\left(\frac{m_s}{4}\right)}}$	Ns/m
$b_2$	$\frac{0.6(2)m_s}{4} \sqrt{\frac{k_1}{\left(\frac{m_s}{4}\right)}}$	Ns/m
$b$	(variable coeff.) * $b_1$	Ns/m

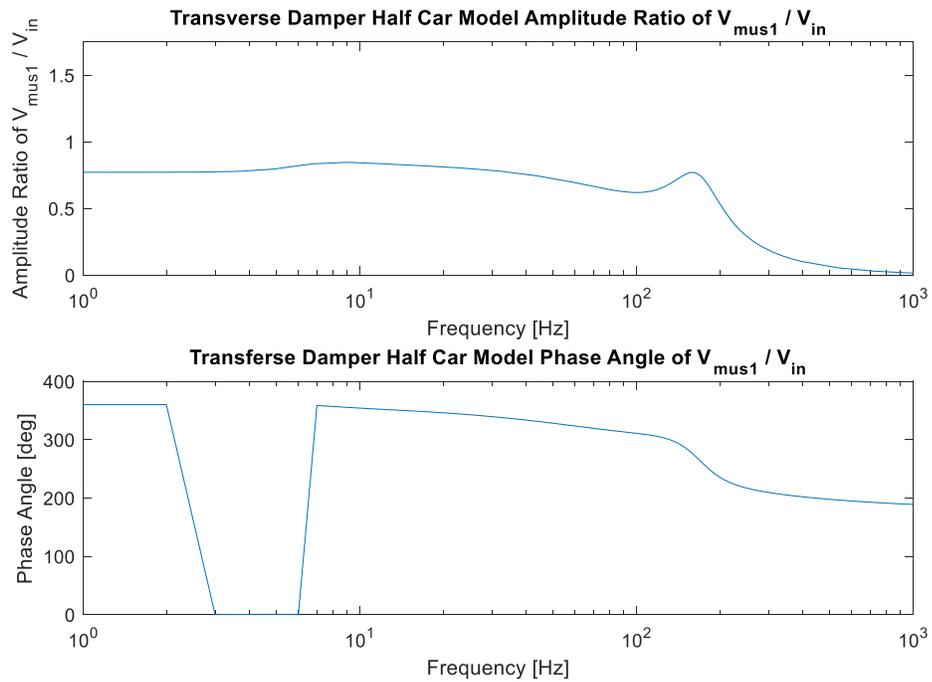


Figure 33: Frequency response plots of the left un-sprung mass velocity to the input velocity in the transverse damper half-car model.

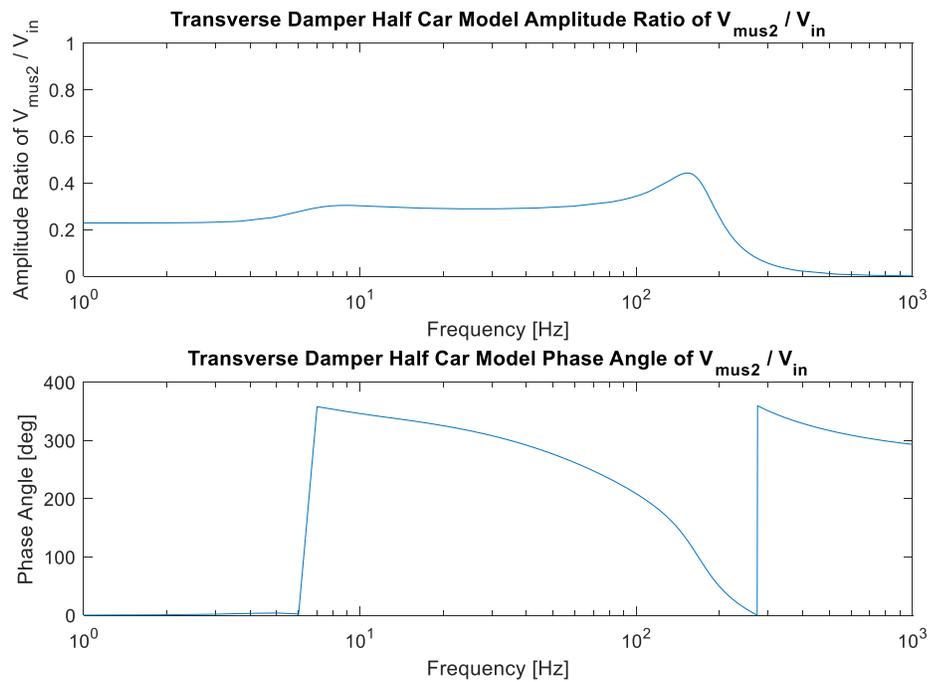


Figure 34: Frequency response plots of the right un-sprung mass velocity to the input velocity in the transverse damper half-car model.

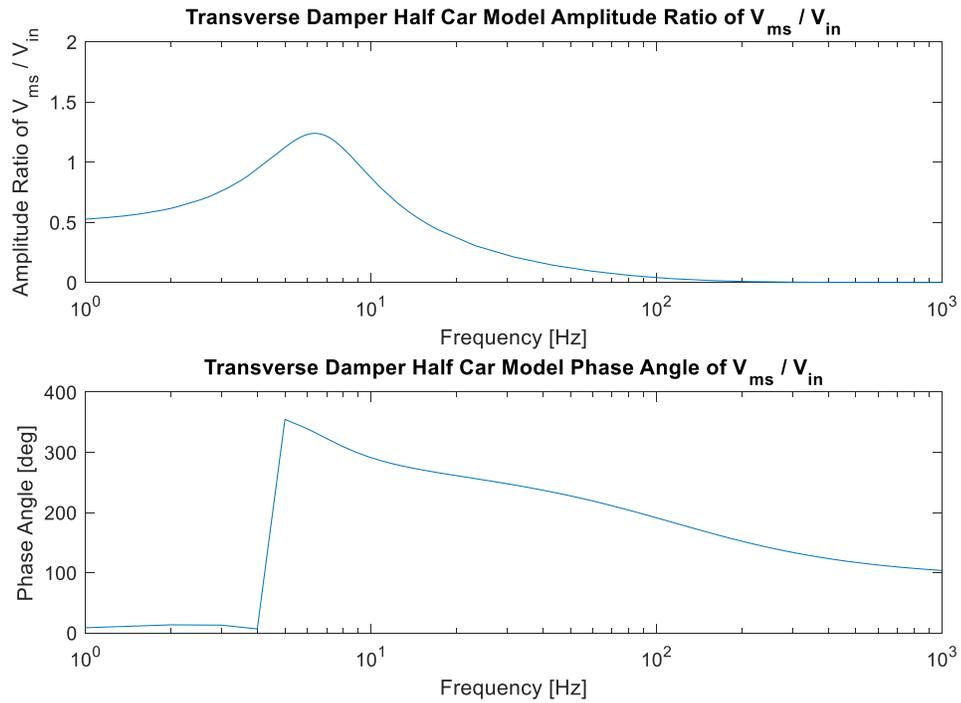


Figure 35: Frequency response plots of the sprung mass velocity to the input velocity in the transverse damper half-car model.

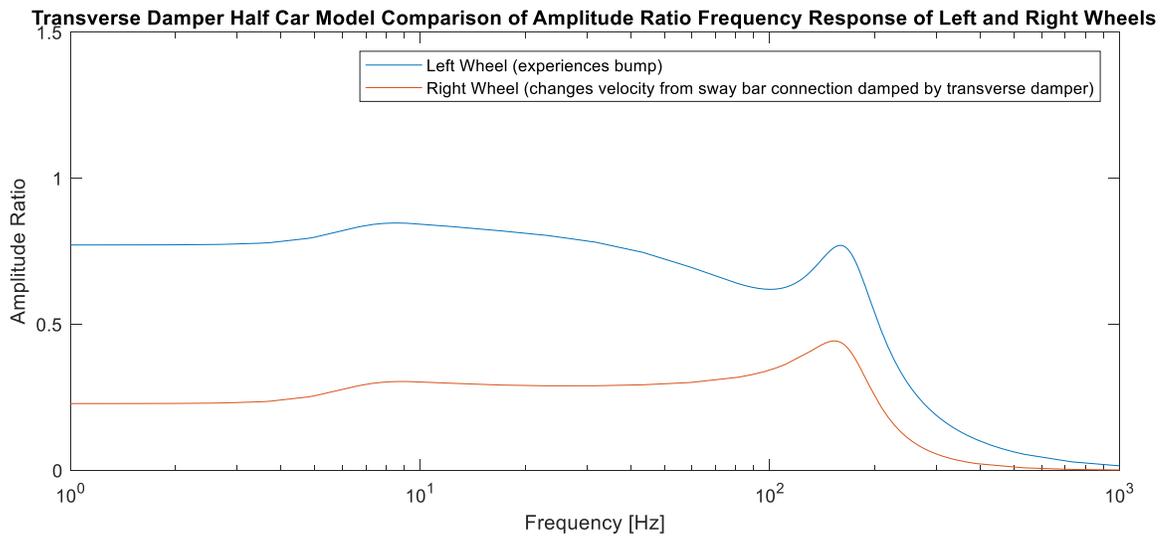


Figure 36: Visual comparison of behavior between left and right wheels only connected by the sway bar and transverse damper.

### Discussion

At long last, the final plot that is meant to truly tell the difference between the effectiveness of a system with a transversely mounted third damper compared to one without.

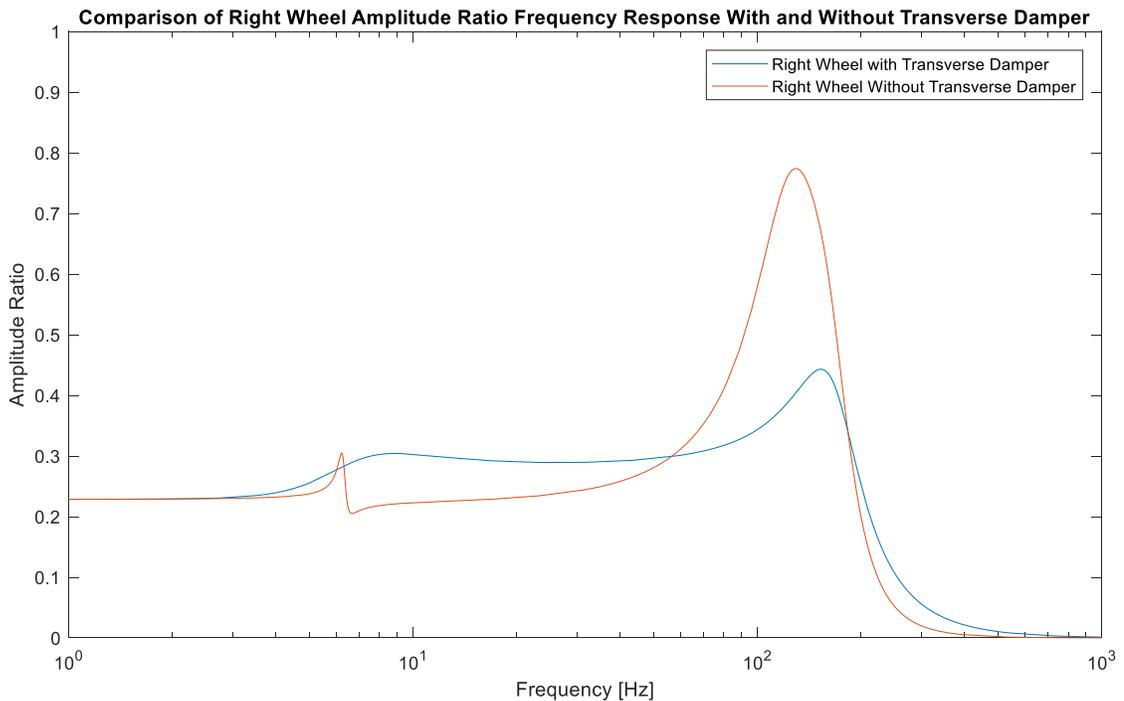


Figure 37: Comparison of right wheel response with and without transverse damper.

Through a grueling process, we have arrived at a visual plot that represents the amplitude ratios of the right un-sprung masses of the two compared systems, and it is clear that there is a winner among the two. Both of these systems were given the same values for all the parameters, (spring stiffness, tire stiffness, damping coefficients, etc.) aside from the obvious additional third damper which was an added parameter to the transverse damper half car model. However, it is clear that the system with the added transversely mounted damper caused the right wheel (the wheel that rides along flat

ground) to have a significantly lesser amplitude ratio in relation to the input velocity acting on the left side of the vehicle. Ultimately, this means that the contact patch and contact forces between the right tire and the road are less affected, so power transmission and grip are more steady, permitting greater tangential force from the ground and a faster race car.

It is imperative to note that the values assigned to the parameters in this system by no means represent an optimally tuned vehicle. The values assigned were found from FSAE forums and online sources according to what teams have used in the past to get reasonably good results. This means that there is a great deal of tuning to be done to actually get the frequency response to something that will benefit the driving experience, rather than accidentally cause resonance between masses that can go out of control very quickly for a given input frequency. However, for a primitive approach to the suspension design of a vehicle, in considering whether or not to make an addition of a third, transversely mounted damper, it appears that it would provide a great benefit to the vehicle's power transmission to the ground by having a lesser effect on wheels traveling on flat ground that are pulled up by the anti-roll bar torsion.

### **Semi-Active System Control with Magnetorheological Dampers**

There exist three types of suspension system control: passive, semi-active, and active. In passive systems, the energy parameters of the system (spring stiffness and damping coefficient) are set to static values. This only allows the suspension setup to be effective over a specific frequency range. However, if the system and its operating conditions are changed, the optimum performance cannot be achieved. Active control

systems however impart a force or velocity on the system to directly counteract specific behaviors over a very wide range of frequencies, but generally require a large power supply or are heavy and are therefore not ideal for racing situations. Semi-active control changes the vehicle parameters to change the behavioral response of the system without imparting an additional force from an outside energy source. This type of control is highly appealing with its low-mass implementation and no need for additional energy inputs, with applications over a much wider frequency range than passive suspension.

With the rise in automotive electronics and sensor monitoring contributing to computational control of vehicle dynamics parameters, the opportunities for suspension control are increased greatly. Magnetorheological (MR) damper technology hit the commercial market in about 2004, developed by the General Motors Delphi Automotive Division. The opportunities of this semi-active control mechanism were limited at first due to a lack of understanding and controllability. However, further iterations and newer models coupled with more accurate sensors and faster reactive computing has led to an innovative means of suspension design. Koenigsegg included these MR dampers into their vehicles for a highly controllable system whose frequency response behavior can be semi-actively controlled through both reactive and predictive settings of the MR dampers to change the effective hydraulic fluid viscosity and alter the damping coefficient of the damper.

Magnetorheological fluid is a fluid with iron particles in it that responds to a magnetic stimulus, changing shape and/or viscosity. Engineers have implemented the use of magnetorheological fluid within dampers by adding iron particles to hydraulic fluid,

such that when a magnetic field is induced within the fluid, it causes these iron particles to align themselves into lines. This particle alignment increases the viscosity of the fluid by making it more resistant to deformation by shear stress. Additionally, the magnetic field is generated by a series of wires running perpendicular to the length of the cylinder that have electric current sent through them. Ampere's Law dictates that the magnitude of the magnetic field generated around a current carrying wire is proportional to the current in the wire serving as its source. In this way, as magnetorheological fluid viscosity is proportional to the strength of the magnetic field, so it is also proportional to the magnitude of the current sent through the wires.

Further, to demonstrate the superior adjustability of magnetorheological dampers, magnetorheological dampers have been designed with current carrying wires going in opposite directions and timed such that the residual eddy-currents from the magnetic field generated by the first set of wires is negated by the second set. For example, one set of wires all travel in parallel directions, and the current through them induce a magnetic field, aligning iron particles in the fluid. When the viscosity is desired to be lower however, the current is shut off, but it takes a fraction of a second for the magnetic field to dissipate due to residual eddy currents, meaning that those aligned iron particles are not yet released from their positions in the fluid so the viscosity remains higher even after current is cut off. These eddy currents are counteracted by adding an additional second set of wires to the damper so that as the first set are being shut off, the second set activates and produces a magnetic field in the opposite direction working to immediately disrupt the field that existed and scatter the aligned particles, only activating long enough

to disrupt these residual eddy currents and then shutting off before realigning the particles with the new magnetic field orientation.

However, the practical use of MR dampers is significantly hindered by its characteristic hysteretic and non-linear dynamics. Such behavior makes the effective damping coefficient and damping forces with a given velocity difficult to calculate and predict. Therefore, for a successful implementation of MR dampers into a system, it demands an accurate model of the damper, and an effective control strategy in how the system calculates the current to be produced through the damper based on input variables. There exist a few different models that are designed to help predict the dynamics of the MR damper. These models include NARX, viscoelastic-plastic, nonlinear blackbox, fuzzy, neural network developed by Chang and Roschke, phenomenological based on a Bouc-Wen hysteresis model, and polynomial among others.

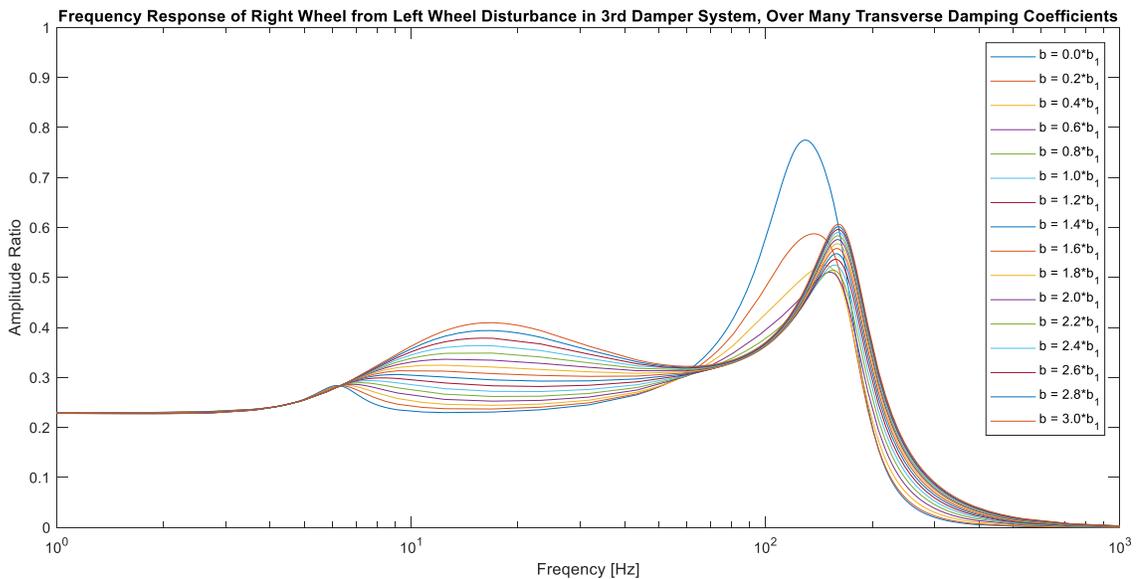


Figure 38: Right wheel frequency response plotted over a range of transverse damping coefficients. These coefficients are given as ratios of the damping coefficient at each individual wheel ( $b_1$ ) which has a damping ratio of 0.6, for reference with respect to other vehicle parameters.

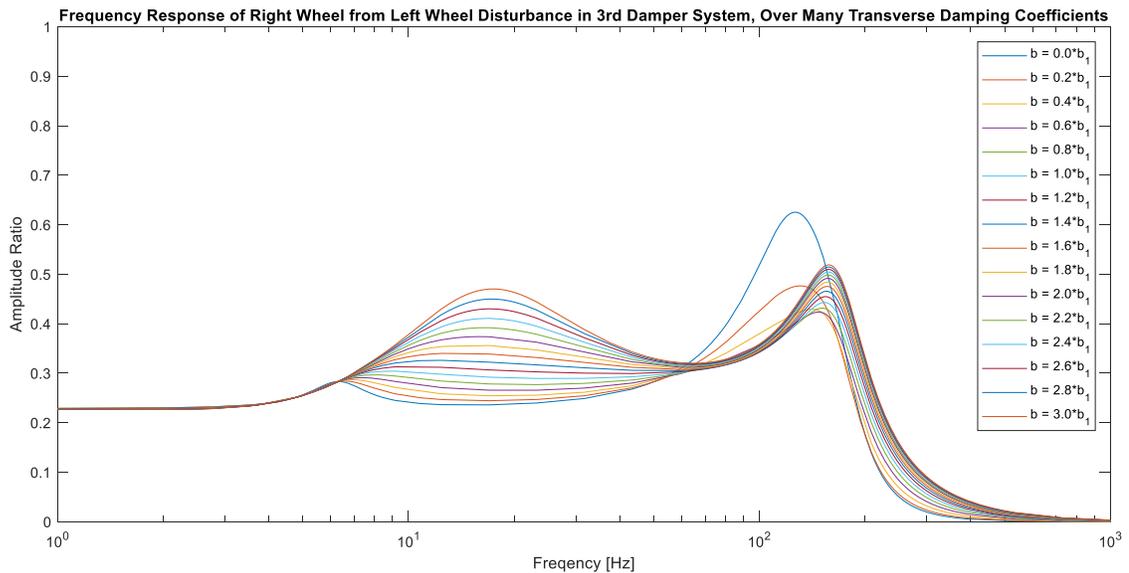


Figure 39: Right wheel frequency response plotted over a range of transverse damping coefficients. The  $b_1$  damping coefficient is given a damping ratio of 0.707 as opposed to 0.6 in Figure 38.

Figure 38 and 39 give great insight into the areas of control available with a variable damping coefficient. Using an MR damper the frequency response of a system is no longer limited to a single curve, but a whole range based on a variable damping coefficient. If a successful model of the MR damper behavior is created such that the effective damping coefficient is accurately related to an electrical input current, the system could be coded to input a specific current based on sensor data to immediately set a damping coefficient that will yield optimum control at the frequency observed. For example, as the vehicle races around the track, the sensors will constantly be feeding data into the onboard computer. The desired damping force can be calculated according to the force and velocity of the input, and if the damper motion velocity is also measured by a sensor, the input current can be calculated according to the model generated of the MR damper behavior to yield the desired damping force.

### **Conclusion**

Though a suspension system is simple in principle, the design process is no small feat – especially for racing applications where every fraction of a second counts. Further, the analysis of suspension systems is arguably more complicated to yield helpful results that may be interpreted allowing for improvement to the system as a whole. The dynamic systems bond graph approach to this comparison of suspension designs makes for an easy to follow process, but is sometimes not so easy to manipulate and interpret the results. In the process of calculating the necessary determinants, transfer functions, amplitude ratios, and phase angles necessary to plot frequency response of different aspects of the system, at times the mathematics seemed overwhelming. Eight state variables results in an 8x8 linear system when converted to the Laplacian domain, such that when transfer functions were calculated through MATLAB, the resulting equation put into a Word document was well over a page in length. This goes to show the overall complexity of the systems being analyzed, and the power of modern computation. Symbolic 8x8 determinants alone would take dozens of hours to calculate, but when the matrices of the system are entered correctly into MATLAB, they can be calculated in fractions of a second. There is much work to be done as far as tuning the values assigned to the system parameters to yield more desirable frequency response plots, which will be done in further work on the subject outside of this report.

Ultimately, this investigation and comparison was a challenging process that required the full application of Dynamic Systems Modeling to work through the analysis

from start to finish. I am glad for the challenge and look forward to continuing this work in the full design of the Liberty University FSAE suspension system.

### **Limitations and Future Work**

This work could be improved by mathematical plots with input values and parameters measured from a physical system for greater accuracy and meaningful results. In my limited knowledge of suspension design, though I have learned a great deal to reach the current design of this system, I have a great deal left to learn in optimizing each aspect of force and movement that the suspension is responsible for. With this being Liberty University's first FSAE vehicle, I have no previous model to test and learn from, only my textbooks and online resources.

Further investigation is planned to be done by empirical testing on the manufactured Formula SAE vehicle with and without the transverse third damper connected to the system to gather squat and dive curves, and best fit an effective quarter car damping coefficient to simulate the anti-squat and anti-dive characteristics of a third damper addition. Additionally, single-sided disturbances will be investigated in an effort to gather vertical displacement data of the opposite wheel to validate the mathematical model created by the dynamic systems analysis and plotted in MATLAB. Measured coefficients from the vehicle such as mass, spring rate, damping coefficient etc. can be gathered and input into the MATLAB program to establish the mathematical baseline against which to compare empirical data plots.

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## Appendix A

**Mass Spring Damper System: 2<sup>nd</sup> Order Differential Equation Development:**

Some input of displacement is put into the system by the road that causes the spring and damper to compress, but they still apply a force to the sprung mass that they are suspending, which is oscillatory in nature (due to the spring) once the input disturbance is returned to a steady flat surface. However, if the stiffness value of the spring and the damping coefficient of the damper are established correctly, the system can be controlled to dampen out oscillatory effects very quickly. The second order differential equation that describes this behavior is given generally as follows.

$$m \frac{d^2 x_{out}}{dt^2} + b \frac{dx_{out}}{dt} + kx_{out} = F_{in} \quad (28)$$

Here  $m$  is the sprung mass,  $b$  is the damping coefficient, and  $k$  is the spring stiffness. If the derivatives of  $x_{out}$  are replaced by Laplacian operators to form a characteristic equation, it looks like this for the homogenous case:

$$ms^2 + bs + k = 0 \quad (29)$$

Solved for  $s$ , this comes out to:

$$s_{1,2} = -\frac{b}{2m} \pm \left( \left( \frac{b}{2m} \right)^2 - \frac{k}{m} \right)^{0.5} \quad (30)$$

In which if there is no damping, in other words the damping coefficient  $b=0$ ,

$$s_{1,2} = \pm j \sqrt{\frac{k}{m}} \quad (31)$$

Signifying pure oscillatory motion, in which the natural frequency is given by:

$$\omega_n = \sqrt{\frac{k}{m}} \quad (32)$$

Following this point in the analysis, there are three possible cases in which the solved radicand is either equal to zero (critical damping), less than zero (underdamped), or greater than zero (overdamped). For the first case in the event in which the two solutions of  $s$  are equal and real, in which,

$$s_1 = s_2 = -\frac{b}{2m} \quad (33)$$

the resulting homogenous solution to the differential equation is therefore:

$$x_{out_h}(t) = c_1 e^{-\frac{b}{2m}t} + c_2 t e^{-\frac{b}{2m}t} \quad (34)$$

It should therefore be noted that if,

$$\frac{b}{2m} = \sqrt{\frac{k}{m}} \quad (35)$$

then the damping coefficient  $b$  in this case is the critical damping coefficient,  $b_c$ , because the exponent in the homogeneous solution would be equal to the previously determined resonant frequency. The critical damping coefficient may be algebraically solved for and is determined to be,

$$b_c = 2m \sqrt{\frac{k}{m}} \quad (36)$$

Now a new term zeta may be introduced to simplify the relationship, being equal to the ratio of the damping coefficient to the critical damping coefficient.

$$\zeta = \frac{b}{b_c} = \frac{b}{2m\omega_n} \quad (37)$$

Now equation 4 may be re-written with the zeta variable as follows:

$$s_{1,2} = \zeta\omega_n \pm \omega_n\sqrt{\zeta^2 - 1} \quad (38)$$

The second case of the solution is one in which the radicand is less than zero, and there exist two complex distinct roots as the solution to the above solution in equation 12, given by

$$s_{1,2} = -\zeta\omega_n \pm \omega_n j\sqrt{1 - \zeta^2} \quad (39)$$

This way the homogenous solution is given by:

$$x_{out_h}(t) = e^{-\zeta\omega_n t} \left\{ c_1 \cos \left[ \left( \omega_n \sqrt{1 - \zeta^2} \right) t \right] + c_2 \sin \left[ \left( \omega_n \sqrt{1 - \zeta^2} \right) t \right] \right\} \quad (40)$$

And the final possibility is where the radicand is greater than zero, such that there exist 2 distinct real roots to the equation, given by

$$s_{1,2} = -\zeta\omega_n \pm \omega_n\sqrt{\zeta^2 - 1} \quad (41)$$

In which the resulting homogenous solution is written simply as

$$x_{out_h}(t) = c_1 e^{(-\zeta\omega_n + \omega_n\sqrt{\zeta^2 - 1})t} + c_2 e^{(-\zeta\omega_n - \omega_n\sqrt{\zeta^2 - 1})t} \quad (42)$$

The total solution then is to plot together the applicable case of the homogenous solution with the particular solution, in which the particular solution is given as follows:

$$x_{outp} = \frac{F_{in}}{k} \quad (43)$$

$$x_{outTotal} = x_{outh} + x_{outp} \quad (44)$$

## Appendix B

## MATLAB Code Generated for Four-Bar Linkage Analysis

```

% Suspension Geometry 4-bar Linkage Analysis
clear vars

%Required Characteristics for FSAE Design:
% Minimum wheel travel = 2.00 inches
% Minimum Jounce = 1.00 inch
% Minimum Rebound = 1.00 inch

%Set lengths of suspension members [all dimensions in inches] for 25% RC/CG height:
a = 13.71; % a = upper wishbone length
b = 8.08; % b = upright connection distance of wishbones
c = 14.84; % c = lower wishbone length
d = 6.0; % d = fixed distance between frame members mounting points
e = 0; % e = horizontal displacement between top and bottom frame members
%f = 300/25.4; %total vertical height from absolute bottom frame member to top mounting
frame member
if c+b <= c+a
    disp('Grashof Law Verified')
else
    disp('Grashof Law Violated')
end

%Calculation of Position for the System:
theta_1 = 0; %by the way the diagram is drawn:

whl_trvl = 2.25;
% x = [0 e e+a*cosd(theta_4-theta_7-theta_2) c*cosd(theta_7) 0];
% y = [0 d*cosd(theta_5) d*cosd(theta_5)-a*sind(theta_4-theta_7-theta_2) c*sind(theta_7)
0];
% labels = {'Bottom Frame','Top Frame','Top of Upright','Bottom of Upright','Bottom
Frame'};

theta_2 = (80:105);
for theta_2 = (90:100)
    K1 = d/a;
    K2 = d/c;
    K3 = (a^2 -b^2 +c^2 + d^2)/(2*a*c);
    K4 = d/b;
    K5 = (c^2-d^2-a^2-b^2)/(2*a*b);

    A = cosd(theta_2)-K1-K2*cosd(theta_2)+K3;
    B = -2*sind(theta_2);
    C = K1-(K2+1)*cosd(theta_2) + K3;
    D = cosd(theta_2)-K1 + K4*cosd(theta_2) +K5;
    E = -2*sind(theta_2);
    F = K1+(K4-1)*cosd(theta_2)+K5;
    theta_3 = 2*atand((-E-(E^2-4*D*F)^0.5)/(2*D));
    w = [0 a*sind(90-theta_2) a*sind(90-theta_2)+b*cosd(theta_3) d 0];
    z = [0 a*cosd(90-theta_2) a*cosd(90-theta_2)+b*sind(theta_3) 0 0];
    pause(0.02);
    hold on

%Plot of all 4 links of suspension geometry through wheel travel
% subplot(2,2,1);
figure(1)
plot(w,z,'.-');
title('Suspension Geometry: Modeled as 4-bar Linkage');
xlabel('Vertical Distance [inches]');
ylabel('Horizontal Distance [inches]');
% text(x,y,labels,'VerticalAlignment','bottom','HorizontalAlignment','left');
axis([-10 18 -10 18]);
hold on

```

```

%Plotting the axis of equilibrium ride height for desired ground clearance
grnd_clrnce = 1.5; %distance from ground surface to center of lowest frame member.
%desired actual ground clearance is 1.0 inches, to permit wheel travel of
%1in and retain clearance without vehicle contact to ground.

%Front Suspension Geometry
F_upper = 11.13; %height in inches from ground
F_lower = 4.63; %height in inches from ground
%Rear Suspension Geometry
R_upper = 10.74; %height in inches from ground
R_lower = 4.74; %height in inches from ground

d_tire = 16;
r_tire = d_tire/2;
h = r_tire - grnd_clrnce; %height of wheel center above bottom frame member
line([R_upper-r_tire,R_upper-r_tire], get(gca, 'ylim'));
text(R_upper-r_tire,-2,'\leftarrow Axis of Equilibrium Ride Height','FontSize',8);

%Plot of only Upright through wheel travel
upright_w = [a*sind(90-theta_2) a*sind(90-theta_2)+b*cosd(theta_3)];
upright_z = [a*cosd(90-theta_2) a*cosd(90-theta_2)+b*sind(theta_3)];
% subplot(2,2,2);
figure(2)
plot(upright_w,upright_z,'.-');
title('Close-Up of Exaggerated Upright Camber Angle Variation Throughout Wheel Travel')
xlabel('Vertical Distance [inches]');
ylabel('Horizontal Distance [inches]');
axis([-3 10 13 15.5]);
hold on

%Plotting the axis of equilibrium ride height for desired ground clearance
line([R_upper-r_tire,R_upper-r_tire], get(gca, 'ylim'));
text(R_upper-r_tire,5,'\leftarrow Axis of Equilibrium Ride Height','FontSize',8);
line([R_upper-r_tire+whl_trvl/2,R_upper-r_tire+whl_trvl/2], get(gca, 'ylim'));
text(R_upper-r_tire+whl_trvl/2,6,'\leftarrow Maximum Rebound Distance','FontSize',8);
line([R_upper-r_tire-whl_trvl/2,R_upper-r_tire-whl_trvl/2], get(gca, 'ylim'));
text(R_upper-r_tire-whl_trvl/2,4,'\leftarrow Maximum Jounce Distance','FontSize',8);

wheel_center = [(a*sind(90-theta_2)+a*sind(90-theta_2)+b*cosd(theta_3))/2 (a*cosd(90-
theta_2)+a*cosd(90-theta_2)+b*sind(theta_3))/2];
disp('wheel center coordinates:');
disp(wheel_center)

wheel_center_h = (a*sind(90-theta_2)+a*sind(90-theta_2)+b*cosd(theta_3))/2;
wheel_center_v = (a*cosd(90-theta_2)+a*cosd(90-theta_2)+b*sind(theta_3))/2;
% subplot(2,2,3);
figure(3)
plot(wheel_center_h,wheel_center_v,'.-');
title('Wheel Center Point through Wheel Travel');
xlabel('Vertical Position');
ylabel('Horizontal Position');
hold on
line([R_upper-r_tire,R_upper-r_tire], get(gca, 'ylim'));
line([R_upper-r_tire+whl_trvl/2,R_upper-r_tire+whl_trvl/2], get(gca, 'ylim'));
text(R_upper-r_tire+whl_trvl/2,10,'\leftarrow Maximum Rebound Distance','FontSize',8);
line([R_upper-r_tire-whl_trvl/2,R_upper-r_tire-whl_trvl/2], get(gca, 'ylim'));
text(R_upper-r_tire-whl_trvl/2,10.5,'\leftarrow Maximum Jounce Distance','FontSize',8);
axis([1 4.5 13.5 15]);

Initial_slope = 0;

%A plot of camber angle versus input theta_2
y2 = a*cosd(90-theta_2)+b*sind(theta_3);
y1 = a*cosd(90-theta_2);
x2 = a*sind(90-theta_2)+b*cosd(theta_3);
x1 = a*sind(90-theta_2);
hold on

```

```

slope = (y2-y1)/(x2-x1)+Initial_slope;
camber = atand(slope)-8.3;
% subplot(2,2,4);
figure(4)
plot(theta_2, camber,'.-');
hold on
title('Camber Angle versus Theta 2');
xlabel('Theta 2 (degrees)');
ylabel('Camber Angle (degrees)');
%axis([70 95 2 8]);
end

%Putting labels on the final points to give reference to the plots
% subplot(2,2,1);
labels1 = {'Top Frame', 'Top of Upright', 'Bottom of Upright', 'Bottom Wishbone Mount', 'Top
Frame'};
text(w,z,labels1, 'VerticalAlignment', 'bottom', 'HorizontalAlignment', 'left');
%subplot(2,2,2);
labels2 = {'Top of Upright', 'Bottom of Upright'};
text(upright_w, upright_z, labels2, 'VerticalAlignment', 'bottom', 'HorizontalAlignment', 'left
');
%Listing output variables
theta_4 = 2*atand((-B-(B^2-4*A*C)^0.5)/(2*A));
%(These angles are for the proper orientation x-y attempt...)
% theta_5 = asind(e/a);
% theta_6 = theta_4-theta_2;
% theta_7 = 180-theta_4;
Angular_Position_Variables = {'theta_1'; 'theta_2'; 'theta_3'; 'theta_4'};
angles = [theta_1; theta_2; theta_3; theta_4];% theta_5; theta_6; theta_7];
T = table(Angular_Position_Variables, angles);

%calculating range of theta_2 that yields desired total wheel travel:

%calculate wheel center distance between first theta 2 input and last theta
%2 input

%Camber off-set to consider from inclination angle = 8.3 degrees

```

## Appendix C

## MATLAB Code Generated for Dynamic Systems Vehicle Response Analysis

```

%Honors Petition Analysis

%Quarter Car Model:
syms S m_us m_s k b k_t omega
A_1 = [0 0 -1/m_us 0; 0 0 1/m_us -1/m_s; k_t -k -b/m_us b/m_s; 0 k b/m_us -b/m_s];
B_1 = [1 0 0 0];
sIA = S*eye(4)-A_1;
Num_1 = [sIA(:,1:3) B_1(:,1)];

%Transfer Function Vm/Vin
Vm_Vin = 1/m_s*det(Num_1)/det(sIA);
collect(Vm_Vin);

%Replace S with j*omega for harmonic input
TF1 = subs(Vm_Vin, S, omega*1i);

%Setting Values
m_us = 8; %kg
k = 4000; %N/m - this is the lowest available spring rate from KAZ tech, equivalent to
175lbf/in
k_t = 123464; %N/m
m_s = 182/4; %kg - for an estimated 400lb vehicle
b = 0.6*2*m_s*sqrt(k/m_s); %For damping ratio = 0.6
assume(omega, 'real');

TF1_num = -(k_t*(k + b*omega*1i));
TF1_denom = (k*m_s*omega^2 - k*k_t + k*m_us*omega^2 + k_t*m_s*omega^2 - m_s*m_us*omega^4
- b*k_t*omega*1i + b*m_s*omega^3*1i + b*m_us*omega^3*1i);
TF1_actual = TF1_num/TF1_denom;

%Amplitude Ratio and Phase Angle Calculations
amp_ratio_TF1 =
sqrt(real(TF1_num)^2+imag(TF1_num)^2)/sqrt(real(TF1_denom)^2+imag(TF1_denom)^2);
phi_1 = atan2d(imag(TF1_num),real(TF1_num))-atan2d(imag(TF1_denom),real(TF1_denom));
phi_1 =subs(phi_1,omega,[0:1000]);
phi_1 = double(phi_1);
phi_1 = phi_1.*(phi_1 >= 0) + (phi_1 + 360) .* (phi_1 < 0);

figure(1);
subplot(2,1,1)
fplot(amp_ratio_TF1, [0 1000]);
title('Quarter Car Model Amplitude Ratio vs Frequency');
xlabel('Frequency [Hz]');
ylabel('Amplitude Ratio Tf V_m_s / V_i_n');
ylim([-0.1 2.0]);

subplot(2,1,2);
plot(0:1000,phi_1);
title('Quarter Car Model Phase Angle Frequency Response');
xlabel('Frequency [Hz]');
ylabel('Phase Angle [deg]');

%%
%Conventional Transverse Half Car Model

syms S m_us1 m_us2 m_s b_b_1 b_2 k_k_1 k_2 k_t1 k_t2 V_in g omega
A_2 = [0 0 0 0 0 -1/m_us1 0 0; 0 0 0 0 0 1/m_us1 -1/m_s -1/m_s;
0 0 0 0 0 -1/m_s 1/m_us2; 0 0 0 0 0 0 1/m_us2;
0 0 0 0 1/m_us1 0 -1/m_us2; k_t1 -k_1 0 0 -k -b_1/m_us1 b_1/m_s 0;
0 k_1 k_2 0 0 b_1/m_us1 -(b_1-b_2)/m_s b_2/m_us2; 0 0 -k_2 -k_t2 k 0 b_2/m_s -
b_2/m_us2];
B_3 = [1 0 0 0; 0 0 0 0; 0 0 0 0; 0 0 0 0; 0 0 0 0; 0 1 0 0; 0 0 0 -1; 0 0 1 0];

```

```

sIA = S*eye(8)-A_2;

%Transfer Function of Vmus1/Vin
Num_2 = [sIA(:,1:5),B_3(:,1),sIA(:,7:8)]; %Cramer's Rule: replacing the p_4 column with
V_in
TF2 = 1/m_us1*det(Num_2)/det(sIA);
collect(TF2);

%Transfer Function of Vmus2/Vin
Num_3 = [sIA(:,1:7),B_3(:,1)];
TF3 = 1/m_us2*det(Num_3)/det(sIA);

%Transfer Function of Vms/Vin
Num_4 = [sIA(:,1:6),B_3(:,1),sIA(:,8)];
TF4 = 1/m_s*det(Num_4)/det(sIA);
collect(TF4);

%Determining the Amplitude Ratio and Phase Angle Frequency Response of the
%Conventional Half Car Model Transfer Functions:
TF2 = subs(TF2, S, omega*1i);
TF3 = subs(TF3, S, omega*1i);
TF4 = subs(TF4, S, omega*1i);

%Assign values to symbolic parameters to give a plot for comparison:
%These values are made considering a typical FSAE racecar
m_s = 200; %kg
m_us1 = 8; %kg
m_us2 = m_us1; %kg
k_1 = 4000; %N/m
k_2 = k_1; %N/m
k_t1 = 123464; %N/m
k_t2 = k_t1; %N/m
k = 50000;
b_1 = 0.6*2*m_s/4*sqrt(k_1/(m_s/4)); %Should be between 0.5 and 0.7 damping ratio.
b_2 = b_1;

%Plotting Transfer Function 2, left unsprung mass velocity to V_in
TF2_num = -(k_t1*(2*b_2^2*m_s*omega^4 - b_2^2*m_s^2*omega^5*1i - k*m_s^2*omega^4 -
k_2*m_s^2*omega^4 - k_t2*m_s^2*omega^4 + m_s^2*m_us2*omega^6 - b_1*b_2*m_s*omega^4 +
b_1*k*m_s*omega^3*1i - b_2*k*m_s*omega^3*1i + b_1*k_t2*m_s*omega^3*1i +
b_2*k_1*m_s*omega^3*1i - b_2*k_2*m_s*omega^3*2i + b_2*k_1*m_us2*omega^3*1i +
b_1*k_t2*m_s*omega^3*1i - b_2*k_t2*m_s*omega^3*1i - b_1*m_s*m_us2*omega^5*1i +
b_2*m_s*m_us2*omega^5*1i + k*k_1*m_s*omega^2 + k*k_2*m_s*omega^2 + k_1*k_2*m_s*omega^2 +
k_1*k_2*m_us2*omega^2 + k_1*k_t2*m_s*omega^2 + k_2*k_t2*m_s*omega^2 -
k_1*m_s*m_us2*omega^4 - k_2*m_s*m_us2*omega^4));
TF2_denom = (2*b_2^2*m_s*m_us1*omega^6 - b_1*b_2^2*m_s^2*omega^6 + b_1*k*m_s^2*omega^5*1i -
2*b_2^2*k*m_s*omega^4 + b_2*k*k_1*m_s^2*omega^5*1i - 2*b_2^2*k_1*m_s*omega^4 +
b_1*k_2*m_s^2*omega^5*1i + b_2*k_1*m_s^2*omega^5*1i - 2*b_2^2*k_t1*m_s*omega^4 +
b_1*k_t2*m_s^2*omega^5*1i + b_2*k_t1*m_s^2*omega^5*1i - b_1*b_2^2*m_s*omega^5*2i -
b_1*m_s^2*m_us2*omega^7*1i - b_2*m_s^2*m_us1*omega^7*1i + k*k_1*m_s^2*omega^4 +
k*k_2*m_s^2*omega^4 + k*k_t1*m_s^2*omega^4 + k*k_t2*m_s^2*omega^4 + k_1*k_2*m_s^2*omega^4
+ k_1*k_t2*m_s^2*omega^4 + k_2*k_t1*m_s^2*omega^4 + k_t1*k_t2*m_s^2*omega^4 -
k*m_s^2*m_us1*omega^6 - k*m_s^2*m_us2*omega^6 - k_1*m_s^2*m_us2*omega^6 -
k_2*m_s^2*m_us1*omega^6 - k_t1*m_s^2*m_us2*omega^6 - k_t2*m_s^2*m_us1*omega^6 +
m_s^2*m_us1*m_us2*omega^8 - 2*b_1*b_2*k*k_1*m_s*omega^4 - 2*b_1*b_2*k_2*m_s*omega^4 +
b_1*b_2*k_t1*m_s*omega^4 - b_1*b_2*k_t2*m_s*omega^4 - b_1*b_2*m_s*m_us1*omega^6 +
b_1*b_2*m_s*m_us2*omega^6 + b_2*k*k_1*m_s*omega^3*2i + b_2*k*k_2*m_s*omega^3*2i -
b_2*k*k_1*m_us2*omega^3*2i - b_1*k*k_t1*m_s*omega^3*1i - b_1*k*k_t2*m_s*omega^3*1i +
b_2*k*k_t1*m_s*omega^3*1i + b_2*k*k_t2*m_s*omega^3*1i + b_2*k_1*k_2*m_s*omega^3*2i -
b_1*k_2*k_t1*m_s*omega^3*1i - b_2*k_1*k_t1*m_s*omega^3*1i - b_1*k_2*k_t2*m_s*omega^3*1i +
b_2*k_1*k_t2*m_s*omega^3*1i + b_2*k_2*k_t1*m_s*omega^3*2i - b_2*k_1*k_t1*m_us2*omega^3*1i
- b_1*k_t1*k_t2*m_s*omega^3*1i + b_2*k_t1*k_t2*m_s*omega^3*1i +
b_1*k*k_1*m_s*m_us1*omega^5*1i + b_1*k*k_1*m_s*m_us2*omega^5*1i - b_2*k*k_1*m_s*m_us1*omega^5*1i -
b_2*k*k_1*m_s*m_us2*omega^5*1i + b_1*k_2*m_s*m_us1*omega^5*1i + b_2*k_1*m_s*m_us1*omega^5*1i
+ b_1*k_2*m_s*m_us2*omega^5*1i - b_2*k_1*m_s*m_us2*omega^5*1i -
b_2*k_2*m_s*m_us1*omega^5*2i + b_2*k_1*m_s*m_us1*omega^5*1i +
b_1*k_t1*m_s*m_us2*omega^5*1i + b_1*k_t2*m_s*m_us1*omega^5*1i -
b_2*k_t1*m_s*m_us2*omega^5*1i - b_2*k_t2*m_s*m_us1*omega^5*1i -

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b_1*m_s*m_us1*m_us2*omega^7*1i + b_2*m_s*m_us1*m_us2*omega^7*1i - k*k_1*k_t1*m_s*omega^2
- k*k_1*k_t2*m_s*omega^2 - k*k_2*k_t1*m_s*omega^2 - k*k_2*k_t2*m_s*omega^2 -
k_1*k_2*k_t1*m_s*omega^2 - k_1*k_2*k_t2*m_s*omega^2 - k_1*k_2*k_t1*m_us2*omega^2 -
k_1*k_t1*k_t2*m_s*omega^2 - k_2*k_t1*k_t2*m_s*omega^2 + k*k_1*m_s*m_us1*omega^4 +
k*k_2*m_s*m_us1*omega^4 + k*k_2*m_s*m_us2*omega^4 + k_1*k_2*m_s*m_us1*omega^4 +
k_1*k_2*m_s*m_us2*omega^4 + k_1*k_2*m_us1*m_us2*omega^4 + k_1*k_t1*m_s*m_us2*omega^4 +
k_1*k_t2*m_s*m_us1*omega^4 + k_2*k_t1*m_s*m_us2*omega^4 + k_2*k_t2*m_s*m_us1*omega^4 -
k_1*m_s*m_us1*m_us2*omega^6 - k_2*m_s*m_us1*m_us2*omega^6);
TF2_actual = TF2_num/TF2_denom;
% H2 = tf([-56937720257911902017945600000000 -3819498389416736554363503247360000i
3506737252964536091714189532186536543 151045883307874263858068069490434806976i -
74997029763368346898221501022877676397339 -2012966146597960745377582809939607156765866i
561096677042127519547757619417240724567297966
37209735545605006698306619193361073923567616i -
44238226098683991777087695400023451811545964544
1722010398720942306000393237584310330982400000i
877489276075031452779253757246774968320000000000], [3689348814741910323200000000 0 -
291797976455567166397935396635342 0 9050216191234562096011300626214101788 0 -
127139589893816555139911723637762535244094 0
735097324995800450811324296348669012956987392 0 -
57550527568803929465446408905987755443662880768 0
113742093056565281007156335129061752832000000000]);
% figure(3);
% bode(H2);
% title('Conventional Half Car Model Bode Diagram of V_m_u_s_1/V_i_n');
amp_ratio_TF2 =
sqrt((real(TF2_num))^2+(imag(TF2_num))^2)/sqrt((real(TF2_denom))^2+(imag(TF2_denom))^2);

figure(4);
subplot(2,1,1);
fplot(amp_ratio_TF2,[0,1000]);
title('Conventional Half Car Model Amplitude Ratio of V_m_u_s_1 / V_i_n');
xlabel('Frequency [Hz]');
ylabel('Amplitude Ratio of V_m_u_s_1 / V_i_n');
ylim([0 1.75]);
phi_2 = atan2d(imag(TF2_num),real(TF2_num))-atan2d(imag(TF2_denom),real(TF2_denom));
phi_2 =subs(phi_2,omega,[0:1000]);
phi_2 = double(phi_2);
phi_2 = phi_2.* (phi_2 >= 0) + (phi_2 + 360) .* (phi_2 < 0);
subplot(2,1,2);
plot(0:1000,phi_2);
title('Conventional Half Car Model Phase Angle of V_m_u_s_1 / V_i_n');
xlabel('Frequency [Hz]');
ylabel('Phase Angle [deg]');

%Plotting Transfer Function 3, right unsprung mass Velocity to V_in
TF3 num = -(k_t1*m_s*omega*(k*m_s*omega^3*1i - k*k_1*omega*1i - k*k_2*omega*1i -
k_1*k_2*omega*1i + b_1*b_2*omega^3*1i + b_1*k*omega^2 - b_2*k*omega^2 + b_1*k_2*omega^2 +
b_2*k_1*omega^2)*1i);
TF3_denom = (2*b_2^2*m_s*m_us1*omega^6 - b_1*b_2*m_s^2*omega^6 + b_1*k*m_s^2*omega^5*1i -
2*b_2^2*k*m_s*omega^4 + b_2*k*m_s^2*omega^5*1i - 2*b_2^2*k_1*m_s*omega^4 +
b_1*k_2*m_s^2*omega^5*1i + b_2*k_1*m_s^2*omega^5*1i - 2*b_2^2*k_t1*m_s*omega^4 +
b_1*k_t2*m_s^2*omega^5*1i + b_2*k_t1*m_s^2*omega^5*1i - b_1*b_2^2*m_s*omega^5*2i -
b_1*m_s^2*m_us2*omega^7*1i - b_2*m_s^2*m_us1*omega^7*1i + k*k_1*m_s^2*omega^4 +
k*k_2*m_s^2*omega^4 + k*k_t1*m_s^2*omega^4 + k*k_t2*m_s^2*omega^4 + k_1*k_2*m_s^2*omega^4
+ k_1*k_t2*m_s^2*omega^4 + k_2*k_t1*m_s^2*omega^4 + k_t1*k_t2*m_s^2*omega^4 -
k*m_s^2*m_us1*omega^6 - k*m_s^2*m_us2*omega^6 - k_1*m_s^2*m_us2*omega^6 -
k_2*m_s^2*m_us1*omega^6 - k_t1*m_s^2*m_us2*omega^6 - k_t2*m_s^2*m_us1*omega^6 +
m_s^2*m_us1*m_us2*omega^8 - 2*b_1*b_2*k*m_s*omega^4 - 2*b_1*b_2*k_2*m_s*omega^4 +
b_1*b_2*k_t1*m_s*omega^4 - b_1*b_2*k_t2*m_s*omega^4 - b_1*b_2*m_s*m_us1*omega^6 +
b_1*b_2*m_s*m_us2*omega^6 + b_2*k*k_1*m_s*omega^3*2i + b_2*k*k_2*m_s*omega^3*2i -
b_2*k*k_1*m_s*omega^3*2i - b_1*k*k_t1*m_s*omega^3*1i - b_1*k*k_t2*m_s*omega^3*1i +
b_2*k*k_t1*m_s*omega^3*1i + b_2*k*k_t2*m_s*omega^3*1i + b_2*k_1*k_2*m_s*omega^3*2i -
b_1*k_2*k_t1*m_s*omega^3*1i - b_2*k_1*k_t1*m_s*omega^3*1i - b_1*k_2*k_t2*m_s*omega^3*1i +
b_2*k_1*k_t2*m_s*omega^3*1i + b_2*k_2*k_t1*m_s*omega^3*2i - b_2*k_1*k_t1*m_us2*omega^3*1i
- b_1*k_t1*k_t2*m_s*omega^3*1i + b_2*k_t1*k_t2*m_s*omega^3*1i +
b_1*k*m_s*m_us1*omega^5*1i + b_1*k*m_s*m_us2*omega^5*1i - b_2*k*m_s*m_us1*omega^5*1i -
b_2*k*m_s*m_us2*omega^5*1i + b_1*k_2*m_s*m_us1*omega^5*1i + b_2*k_1*m_s*m_us1*omega^5*1i

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+ b_1*k_2*m_s*m_us2*omega^5*1i - b_2*k_1*m_s*m_us2*omega^5*1i -
b_2*k_2*m_s*m_us1*omega^5*2i + b_2*k_1*m_us1*m_us2*omega^5*1i +
b_1*k_t1*m_s*m_us2*omega^5*1i + b_1*k_t2*m_s*m_us1*omega^5*1i -
b_2*k_t1*m_s*m_us2*omega^5*1i - b_2*k_t2*m_s*m_us1*omega^5*1i -
b_1*m_s*m_us1*m_us2*omega^7*1i + b_2*m_s*m_us1*m_us2*omega^7*1i - k*k_1*k_t1*m_s*omega^2
- k*k_1*k_t2*m_s*omega^2 - k*k_2*k_t1*m_s*omega^2 - k*k_2*k_t2*m_s*omega^2 -
k_1*k_2*k_t1*m_s*omega^2 - k_1*k_2*k_t2*m_s*omega^2 - k_1*k_2*k_t1*m_us2*omega^2 -
k_1*k_t1*k_t2*m_s*omega^2 - k_2*k_t1*k_t2*m_s*omega^2 + k*k_1*m_s*m_us1*omega^4 +
k*k_2*m_s*m_us1*omega^4 + k*k_2*m_s*m_us2*omega^4 + k_1*k_2*m_s*m_us1*omega^4 +
k_1*k_2*m_s*m_us2*omega^4 + k_1*k_2*m_us1*m_us2*omega^4 + k_1*k_t1*m_s*m_us2*omega^4 +
k_1*k_t2*m_s*m_us1*omega^4 + k_2*k_t1*m_s*m_us2*omega^4 + k_2*k_t2*m_s*m_us1*omega^4 -
k_1*m_s*m_us1*m_us2*omega^6 - k_2*m_s*m_us1*m_us2*omega^6);
TF3_actual = TF3_num/TF3_denom;
amp_ratio_TF3 =
sqrt((real(TF3_num))^2+(imag(TF3_num))^2)/sqrt((real(TF3_denom))^2+(imag(TF3_denom))^2);

% figure(5);
% H3 = tf([1464438165033494136472618401792000000,
195863877409290252677973733408329303525*1i, -71069840071612843540548377317557308533450, -
4301395089965414876537350139078609409597775*1i,
651628432086768101407606450390502562224777200, -
145563079048854427592038881035805667652096000*1i, -
52039244978822833689677944100460585884232908800,
12707123512235269799761546958005306654720000000*1i +
1039726617962485429169238376175370240000000000000], [14757395258967641292800000000, 0, -
1167191905822268665591741586541368, 0, 36200864764938248384045202504856407152, 0, -
508558359575266220559646894551050140976376, 0,
2940389299983201803245297185394676051827949568, 0, -
230202110275215717861785635623951021774651523072, 0,
4549683722262611240286253405162470113280000000000]);
% bode(H3);
% title('Conventional Half Car Model Bode Diagram V_m_u_s_2/V_i_n')

% figure(6);
% bode(H2);
% hold on
% bode(H3);
% p1=bode(H2);
% p2=bode(H3);
% title('Conventional Half Car Model Left and Right Wheel Frequency Response');
% legend([p1 p2], 'V_m_u_s_1/V_i_n', 'V_m_u_s_2/V_i_n');

figure(7);
subplot(2,1,1);
fplot(amp_ratio_TF3, [0,1000]);
title('Conventional Half Car Model Amplitude Ratio of V_m_u_s_2 / V_i_n');
xlabel('Frequency [Hz]');
ylabel('Amplitude Ratio of V_m_u_s_2 / V_i_n');
ylim([0 1]);
phi_3 = atan2d((imag(TF3_num)), real(TF3_num))-atan2d(imag(TF3_denom), real(TF3_denom));
phi_3 =subs(phi_3,omega, [0:1000]);
phi_3 = double(phi_3);
phi_3 = phi_3.*(phi_3 >= 0) + (phi_3 + 360) .* (phi_3 < 0);
subplot(2,1,2);
plot(0:1000,phi_3);
title('Conventional Half Car Model Phase Angle of V_m_u_s_2 / V_i_n');
xlabel('Frequency [Hz]');
ylabel('Phase Angle [deg]');

%Plotting Transfer Function 4, sprung mass velocity to V_in
TF4_num = -(k_t1*(k*k_1*m_s*omega^2 + b_1*k*m_s*omega^3*1i + b_2*k*m_s*omega^3*1i +
b_1*k_2*m_s*omega^3*1i + b_2*k_1*m_s*omega^3*1i + b_1*k_t2*m_s*omega^3*1i -
b_1*m_s*m_us2*omega^5*1i - b_1*b_2*m_s*omega^4 + k*k_2*m_s*omega^2 - k*k_1*m_us2*omega^2
+ k_1*k_2*m_s*omega^2 + k_1*k_t2*m_s*omega^2 - k_1*m_s*m_us2*omega^4));
TF4_denom = (2*b_2^2*k*m_s*omega^4 + b_2*k*m_s^2*omega^5*1i - 2*b_2^2*k_1*m_s*omega^4 +
b_1*k_2*m_s^2*omega^5*1i + b_2*k_1*m_s^2*omega^5*1i - 2*b_2^2*k_t1*m_s*omega^4 +
b_1*k_t2*m_s^2*omega^5*1i + b_2*k_t1*m_s^2*omega^5*1i - b_1*b_2^2*m_s*omega^5*2i -

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b_1*m_s^2*m_us2*omega^7*1i - b_2*m_s^2*m_us1*omega^7*1i + k*k_1*m_s^2*omega^4 +
k*k_2*m_s^2*omega^4 + k*k_t1*m_s^2*omega^4 + k*k_t2*m_s^2*omega^4 + k_1*k_2*m_s^2*omega^4
+ k_1*k_t2*m_s^2*omega^4 + k_2*k_t1*m_s^2*omega^4 + k_t1*k_t2*m_s^2*omega^4 -
k*m_s^2*m_us1*omega^6 - k*m_s^2*m_us2*omega^6 - k_1*m_s^2*m_us2*omega^6 -
k_2*m_s^2*m_us1*omega^6 - k_t1*m_s^2*m_us2*omega^6 - k_t2*m_s^2*m_us1*omega^6 +
m_s^2*m_us1*m_us2*omega^8 - 2*b_1*b_2*k*m_s*omega^4 - 2*b_1*b_2*k_2*m_s*omega^4 +
b_1*b_2*k_t1*m_s*omega^4 - b_1*b_2*k_t2*m_s*omega^4 - b_1*b_2*m_s*m_us1*omega^6 +
b_1*b_2*m_s*m_us2*omega^6 + b_2*k*k_1*m_s*omega^3*2i + b_2*k*k_2*m_s*omega^3*2i -
b_2*k*k_1*m_us2*omega^3*2i - b_1*k*k_t1*m_s*omega^3*1i - b_1*k*k_t2*m_s*omega^3*1i +
b_2*k*k_t1*m_s*omega^3*1i + b_2*k*k_t2*m_s*omega^3*1i + b_2*k_1*k_2*m_s*omega^3*2i -
b_1*k_2*k_t1*m_s*omega^3*1i - b_2*k_1*k_t1*m_s*omega^3*1i - b_1*k_2*k_t2*m_s*omega^3*1i +
b_2*k_1*k_t2*m_s*omega^3*1i + b_2*k_1*k_t1*m_us2*omega^3*2i - b_2*k_1*k_t1*m_us2*omega^3*1i
- b_1*k_t1*k_t2*m_s*omega^3*1i + b_2*k_t1*k_t2*m_s*omega^3*1i +
b_1*k*m_s*m_us1*omega^5*1i + b_1*k*m_s*m_us2*omega^5*1i - b_2*k*m_s*m_us1*omega^5*1i -
b_2*k*m_s*m_us2*omega^5*1i + b_1*k_2*m_s*m_us1*omega^5*1i + b_2*k_1*m_s*m_us1*omega^5*1i
+ b_1*k_2*m_s*m_us2*omega^5*1i - b_2*k_1*m_s*m_us2*omega^5*1i -
b_2*k_2*m_s*m_us1*omega^5*2i + b_2*k_1*m_us1*m_us2*omega^5*1i +
b_1*k_t1*m_s*m_us2*omega^5*1i + b_1*k_t2*m_s*m_us1*omega^5*1i -
b_2*k_t1*m_s*m_us2*omega^5*1i - b_2*k_t2*m_s*m_us1*omega^5*1i -
b_1*m_s*m_us1*m_us2*omega^7*1i + b_2*m_s*m_us1*m_us2*omega^7*1i - k*k_1*k_t1*m_s*omega^2
- k*k_1*k_t2*m_s*omega^2 - k*k_2*k_t1*m_s*omega^2 - k*k_2*k_t2*m_s*omega^2 -
k_1*k_2*k_t1*m_s*omega^2 - k_1*k_2*k_t2*m_s*omega^2 - k_1*k_2*k_t1*m_us2*omega^2 -
k_1*k_t1*k_t2*m_s*omega^2 - k_2*k_t1*k_t2*m_s*omega^2 + k*k_1*m_s*m_us1*omega^4 +
k*k_2*m_s*m_us1*omega^4 + k*k_2*m_s*m_us2*omega^4 + k_1*k_2*m_s*m_us1*omega^4 +
k_1*k_2*m_s*m_us2*omega^4 + k_1*k_2*m_s_us1*omega^4 + k_1*k_t1*m_s*m_us2*omega^4 +
k_1*k_t2*m_s*m_us1*omega^4 + k_2*k_t1*m_s*m_us2*omega^4 + k_2*k_t2*m_s*m_us1*omega^4 -
k_1*m_s*m_us1*m_us2*omega^6 - k_2*m_s*m_us1*m_us2*omega^6);
amp_ratio_TF4 =
sqrt((real(TF4_num))^2+(imag(TF4_num))^2)/sqrt((real(TF4_denom))^2+(imag(TF4_denom))^2);

```

```

figure(8);
subplot(2,1,1)
fplot(amp_ratio_TF4, [0, 1000]);
title('Conventional Half Car Model Amplitude Ratio of V_m_s/V_i_n');
xlabel('Frequency [Hz]');
ylabel('Amplitude Ratio of V_m_s/V_i_n');
ylim([0 6]);
phi_4 = atan2d((imag(TF4_num)),real(TF4_num))-atan2d(imag(TF4_denom),real(TF4_denom));
phi_4 =subs(phi_4,omega,[0:1000]);
phi_4 = double(phi_4);
phi_4 = phi_4.*(phi_4 >= 0) + (phi_4 + 360) .* (phi_4 < 0);
subplot(2,1,2);
plot(0:1000,phi_4);
title('Conventional Half Car Model Phase Angle of V_m_s/V_i_n');
xlabel('Frequency [Hz]');
ylabel('Phase Angle [deg]');

```

```

figure(9)
fplot(amp_ratio_TF2,[0,1000]);
hold on
fplot(amp_ratio_TF3,[0,1000]);
ylim([0 1.5]);
xlabel('Frequency [Hz]');
ylabel('Amplitude Ratio');
legend('Left Wheel (experiences bump)','Right Wheel (changes velocity from sway bar
connection)');
title('Conventional Half Car Model Comparison of Amplitude Ratio Frequency Response of
Left and Right Wheels');

```

```

%%
%Koenigsegg Transverse 3rd Damper Addition System
syms S m_us1 m_us2 m_s b_b1 b_b2 k_k1 k_k2 k_t1 k_t2 V_in g omega
A_3 = [0 0 0 0 0 -1/m_us1 0 0; 0 0 0 0 0 1/m_us1 -1/m_s -1/m_s;
0 0 0 0 0 0 -1/m_s 1/m_us2; 0 0 0 0 0 0 0 1/m_us2;
0 0 0 0 0 1/m_us1 0 -1/m_us2; k_t1 -k_1 0 0 -k (-b_1-b)/m_us1 (b_1)/m_s -b/m_us2;
0 k_1 k_2 0 0 (b_1)/m_us1 (-b_1-b_2)/m_s (b_2)/m_us2; 0 0 -k_2 -k_t2 k -b/m_us1
(b_2)/m_s (-b-b_2)/m_us2];
B_3 = [1 0 0 0; 0 0 0 0; 0 0 0 0; 0 0 0 0; 0 0 0 0; 0 1 0 0; 0 0 0 -1; 0 0 1 0];

```

```

sIA = S*eye(8)-A_3;
det(sIA);

%Transfer Function of Vmus1/Vin
Num_5 = [sIA(:,1:5),B_3(:,1),sIA(:,7:8)]; %Cramer's Rule: replacing the p_4 column with
V_in
TF5 = 1/m_us1*det(Num_5)/det(sIA);
collect(TF5);

%Transfer Function of Vmus2/Vin
Num_6 = [sIA(:,1:7),B_3(:,1)];
TF6 = 1/m_us2*det(Num_6)/det(sIA);

%Transfer Function of Vms/Vin
Num_7 = [sIA(:,1:6),B_3(:,1),sIA(:,8)];
TF7 = 1/m_s*det(Num_7)/det(sIA);
collect(TF4);

%Determining the Amplitude Ratio and Phase Angle Frequency Response of the
%Conventional Half Car Model Transfer Functions:
TF5 = subs(TF5, S, omega*1i);
TF6 = subs(TF6, S, omega*1i);
TF7 = subs(TF7, S, omega*1i);

%Assign values to symbolic parameters to give a plot for comparison:
%These values are made considering a typical FSAE racecar
m_s = 200; %kg
m_us1 = 8; %kg
m_us2 = m_us1; %kg
k_1 = 4000; %N/m
k_2 = k_1; %N/m
k_t1 = 123464; %N/m
k_t2 = k_t1; %N/m
k = 50000;
b_1 = 0.707*2*m_s/4*sqrt(k_1/(m_s/4)); %Damping ratio should be between 0.5 and 0.7
b_2 = b_1;
b = 1*b_1;

%Plotting Transfer Function 5, left unsprung mass velocity to V_in
TF5_num = (k_t1*(m_s^2*m_us2*omega^6 - b_2*m_s^2*omega^5*1i - k*m_s^2*omega^4 -
k_2*m_s^2*omega^4 - k_t2*m_s^2*omega^4 - b*m_s^2*omega^5*1i + b*k*m_s*omega^3*1i +
b*k_1*m_s*omega^3*1i + b_2*k*m_s*omega^3*1i + b*k_1*m_us2*omega^3*1i +
b*k_t2*m_s*omega^3*1i + b_2*k_1*m_s*omega^3*1i + b_2*k_1*m_us2*omega^3*1i +
b_2*k_t2*m_s*omega^3*1i - b*m_s*m_us2*omega^5*1i - b_2*m_s*m_us2*omega^5*1i +
k*k_1*m_s*omega^2 + k*k_2*m_s*omega^2 + k_1*k_2*m_s*omega^2 + k_1*k_2*m_us2*omega^2 +
k_1*k_t2*m_s*omega^2 + k_2*k_t2*m_s*omega^2 - k_1*m_s*m_us2*omega^4 -
k_2*m_s*m_us2*omega^4));
TF5_denom = (b*b_1^2*m_s*omega^5*1i + b*b_1*m_s^2*omega^6 + b*b_2*m_s^2*omega^6 +
b_1^2*b_2*m_s*omega^5*1i + b_1*b_2*m_s^2*omega^6 - b*k*m_s^2*omega^5*4i +
b_1^2*k*m_s*omega^4 - b*k_1*m_s^2*omega^5*1i - b_1*k*m_s^2*omega^5*1i -
b*k_2*m_s^2*omega^5*1i - b_2*k*m_s^2*omega^5*1i - b*k_t1*m_s^2*omega^5*1i -
b*k_t2*m_s^2*omega^5*1i + b_1^2*k_2*m_s*omega^4 - b_1*k_2*m_s^2*omega^5*1i -
b_2*k_1*m_s^2*omega^5*1i + b_1^2*k_t2*m_s*omega^4 - b_1*k_t2*m_s^2*omega^5*1i -
b_2*k_t1*m_s^2*omega^5*1i + b*m_s^2*m_us1*omega^7*1i + b*m_s^2*m_us2*omega^7*1i -
b_1^2*m_s*m_us2*omega^6 + b_1*m_s^2*m_us2*omega^7*1i + b_2*m_s^2*m_us1*omega^7*1i -
k*k_1*m_s^2*omega^4 - k*k_2*m_s^2*omega^4 - k*k_t1*m_s^2*omega^4 - k*k_t2*m_s^2*omega^4 -
k_1*k_2*m_s^2*omega^4 - k_1*k_t2*m_s^2*omega^4 - k_2*k_t1*m_s^2*omega^4 -
k_t1*k_t2*m_s^2*omega^4 + k*m_s^2*m_us1*omega^6 + k*m_s^2*m_us2*omega^6 +
k_1*m_s^2*m_us2*omega^6 + k_2*m_s^2*m_us1*omega^6 + k_t1*m_s^2*m_us2*omega^6 +
k_t2*m_s^2*m_us1*omega^6 - m_s^2*m_us1*m_us2*omega^8 + 3*b*b_1*k*m_s*omega^4 -
b*b_2*k*m_s*omega^4 + b*b_1*k_1*m_s*omega^4 - b*b_2*k_1*m_s*omega^4 +
b_1*b_2*k*m_s*omega^4 + b*b_1*k_1*m_us2*omega^4 - b*b_2*k_1*m_us2*omega^4 +
b*b_1*k_t2*m_s*omega^4 - b*b_2*k_t2*m_s*omega^4 + b_1*b_2*k_1*m_s*omega^4 -
b_1*b_2*k_t2*m_s*omega^4 - b*b_1*m_s*m_us2*omega^6 + b*b_2*m_s*m_us2*omega^6 +
b_1*b_2*m_s*m_us2*omega^6 + b*k*k_1*m_s*omega^3*1i + b*k*k_2*m_s*omega^3*1i +
b*k*k_1*m_us2*omega^3*1i + b*k*k_t1*m_s*omega^3*1i + b*k*k_t2*m_s*omega^3*1i -
b_1*k*k_1*m_s*omega^3*1i + b*k_1*k_2*m_s*omega^3*1i - b_1*k*k_2*m_s*omega^3*1i +
b_1*k*k_1*m_us2*omega^3*1i + b*k_1*k_2*m_us2*omega^3*1i + b*k_1*k_t1*m_s*omega^3*1i +

```

```

b_2*k*k_t1*m_s*omega^3*1i + b*k_2*k_t2*m_s*omega^3*1i + b_2*k*k_t2*m_s*omega^3*1i +
b*k_1*k_t1*m_us2*omega^3*1i + b*k_t1*k_t2*m_s*omega^3*1i - b_1*k_1*k_2*m_s*omega^3*1i -
b_1*k_1*k_t2*m_s*omega^3*1i + b_2*k_1*k_t1*m_s*omega^3*1i + b_1*k_2*k_t2*m_s*omega^3*1i +
b_2*k_1*k_t2*m_s*omega^3*1i - b*k*m_s*m_us1*omega^5*1i - b*k*m_s*m_us2*omega^5*1i -
b*k_1*m_s*m_us1*omega^5*1i - b*k_1*m_s*m_us2*omega^5*1i - b_2*k*m_s*m_us1*omega^5*1i -
b*k_2*m_s*m_us2*omega^5*1i - b_2*k*m_s*m_us2*omega^5*1i - b*k_1*m_us1*m_us2*omega^5*1i -
b*k_t1*m_s*m_us2*omega^5*1i - b*k_t2*m_s*m_us1*omega^5*1i + b_1*k_1*m_s*m_us2*omega^5*1i
- b_2*k_1*m_s*m_us1*omega^5*1i - b_1*k_2*m_s*m_us2*omega^5*1i -
b_2*k_1*m_s*m_us2*omega^5*1i - b_2*k_1*m_us1*m_us2*omega^5*1i -
b_2*k_t1*m_s*m_us2*omega^5*1i - b_2*k_t2*m_s*m_us1*omega^5*1i +
b*m_s*m_us1*m_us2*omega^7*1i + b_2*m_s*m_us1*m_us2*omega^7*1i + k*k_1*k_t1*m_s*omega^2 +
k*k_1*k_t2*m_s*omega^2 + k*k_2*k_t1*m_s*omega^2 + k*k_2*k_t2*m_s*omega^2 +
k_1*k_2*k_t1*m_s*omega^2 + k_1*k_2*k_t2*m_s*omega^2 + k_1*k_2*k_t1*m_us2*omega^2 +
k_1*k_t1*k_t2*m_s*omega^2 + k_2*k_t1*k_t2*m_s*omega^2 - k*k_1*m_s*m_us1*omega^4 -
k*k_2*m_s*m_us1*omega^4 - k*k_2*m_s*m_us2*omega^4 - k_1*k_2*m_s*m_us1*omega^4 -
k_1*k_2*m_s*m_us2*omega^4 - k_1*k_2*m_us1*m_us2*omega^4 - k_1*k_t1*m_s*m_us2*omega^4 -
k_1*k_t2*m_s*m_us1*omega^4 - k_2*k_t1*m_s*m_us2*omega^4 - k_2*k_t2*m_s*m_us1*omega^4 +
k_1*m_s*m_us1*m_us2*omega^6 + k_2*m_s*m_us1*m_us2*omega^6);
amp_ratio_TF5 =
sqrt((real(TF5_num))^2+(imag(TF5_num))^2)/sqrt((real(TF5_denom))^2+(imag(TF5_denom))^2);

```

```

figure(10);
subplot(2,1,1);
fplot(amp_ratio_TF5,[0,1000]);
title('Transverse Damper Half Car Model Amplitude Ratio of V_m_u_s_1 / V_i_n');
xlabel('Frequency [Hz]');
ylabel('Amplitude Ratio of V_m_u_s_1 / V_i_n');
ylim([0 1.75]);
phi_5 = atan2d((imag(TF5_num)),real(TF5_num))-atan2d(imag(TF5_denom),real(TF5_denom));
phi_5 =subs(phi_5,omega,[0:1000]);
phi_5 = double(phi_5);
phi_5 = phi_5.*(phi_5 >= 0) + (phi_5 + 360) .* (phi_5 < 0);
subplot(2,1,2);
plot(0:1000,phi_5);
title('Transverse Damper Half Car Model Phase Angle of V_m_u_s_1 / V_i_n');
xlabel('Frequency [Hz]');
ylabel('Phase Angle [deg]');

```

```

%Plotting Transfer Function 7, right unsprung mass Velocity to V_in
TF6_num = (k_t1*m_s*omega*(k*m_s*omega^3*1i - k*k_1*omega*1i - k*k_2*omega*1i -
k_1*k_2*omega*1i + b*b_1*omega^3*1i + b_1*b_2*omega^3*1i + b*k*omega^2 + b_2*k*omega^2 +
b_1*k_2*omega^2 + b_2*k_1*omega^2 + b*m_s*omega^4)*1i);
TF6_denom = (b*b_1^2*m_s*omega^5*1i + b*b_1*m_s^2*omega^6 + b*b_2*m_s^2*omega^6 +
b_1^2*b_2*m_s*omega^5*1i + b_1*b_2*m_s^2*omega^6 - b*k*m_s^2*omega^5*4i +
b_1^2*k*m_s*omega^4 - b*k_1*m_s^2*omega^5*1i - b_1*k*m_s^2*omega^5*1i -
b*k_2*m_s^2*omega^5*1i - b_2*k*m_s^2*omega^5*1i - b*k_t1*m_s^2*omega^5*1i -
b*k_t2*m_s^2*omega^5*1i + b_1^2*k_2*m_s*omega^4 - b_1*k_2*m_s^2*omega^5*1i -
b_2*k_1*m_s^2*omega^5*1i + b_1^2*k_t2*m_s*omega^4 - b_1*k_t2*m_s^2*omega^5*1i -
b_2*k_t1*m_s^2*omega^5*1i + b*m_s^2*m_us1*omega^7*1i + b*m_s^2*m_us2*omega^7*1i -
b_1^2*m_s*m_us2*omega^6 + b_1*m_s^2*m_us2*omega^7*1i + b_2*m_s^2*m_us1*omega^7*1i -
k*k_1*m_s^2*omega^4 - k*k_2*m_s^2*omega^4 - k*k_t1*m_s^2*omega^4 - k*k_t2*m_s^2*omega^4 -
k_1*k_2*m_s^2*omega^4 - k_1*k_t2*m_s^2*omega^4 - k_2*k_t1*m_s^2*omega^4 -
k_t1*k_t2*m_s^2*omega^4 + k*m_s^2*m_us1*omega^6 + k*m_s^2*m_us2*omega^6 +
k_1*m_s^2*m_us2*omega^6 + k_2*m_s^2*m_us1*omega^6 + k_t1*m_s^2*m_us2*omega^6 +
k_t2*m_s^2*m_us1*omega^6 - m_s^2*m_us1*m_us2*omega^8 + 3*b*b_1*k*m_s*omega^4 -
b*b_2*k*m_s*omega^4 + b*b_1*k_1*m_s*omega^4 - b*b_2*k_1*m_s*omega^4 +
b_1*b_2*k*m_s*omega^4 + b*b_1*k_1*m_us2*omega^4 - b*b_2*k_1*m_us2*omega^4 +
b*b_1*k_t2*m_s*omega^4 - b*b_2*k_t2*m_s*omega^4 + b_1*b_2*k_1*m_s*omega^4 -
b_1*b_2*k_t2*m_s*omega^4 - b*b_1*m_s*m_us2*omega^6 + b*b_2*m_s*m_us2*omega^6 +
b_1*b_2*k_t1*m_s*omega^6 + b*k*k_1*m_s*omega^3*1i + b*k*k_2*m_s*omega^3*1i +
b*k*k_1*m_us2*omega^3*1i + b*k*k_2*t1*m_s*omega^3*1i + b*k*k_2*t2*m_s*omega^3*1i -
b_1*k*k_1*m_s*omega^3*1i + b*k_1*k_2*m_s*omega^3*1i - b_1*k*k_2*m_s*omega^3*1i +
b_1*k*k_2*m_s*omega^3*1i + b*k_1*k_2*m_us2*omega^3*1i + b*k_1*k_t1*m_s*omega^3*1i +
b_2*k*k_2*t1*m_s*omega^3*1i + b*k_2*k_t2*m_s*omega^3*1i + b_2*k*k_2*t2*m_s*omega^3*1i +
b*k_1*k_t1*m_us2*omega^3*1i + b*k_t1*k_t2*m_s*omega^3*1i - b_1*k_1*k_2*m_s*omega^3*1i -
b_1*k_1*k_t2*m_s*omega^3*1i + b_2*k_1*k_t1*m_s*omega^3*1i + b_1*k_2*k_t2*m_s*omega^3*1i +
b_2*k_1*k_t2*m_s*omega^3*1i + b_2*k_1*k_t1*m_us2*omega^3*1i +

```

```

b_2*k_t1*k_t2*m_s*omega^3*1i - b*k*m_s*m_us1*omega^5*1i - b*k*m_s*m_us2*omega^5*1i -
b*k_1*m_s*m_us1*omega^5*1i - b*k_1*m_s*m_us2*omega^5*1i - b_2*k*m_s*m_us1*omega^5*1i -
b*k_2*m_s*m_us2*omega^5*1i - b_2*k*m_s*m_us2*omega^5*1i - b*k_1*m_us1*m_us2*omega^5*1i -
b*k_t1*m_s*m_us2*omega^5*1i - b*k_t2*m_s*m_us1*omega^5*1i + b_1*k_1*m_s*m_us2*omega^5*1i
- b_2*k_1*m_s*m_us1*omega^5*1i - b_1*k_2*m_s*m_us2*omega^5*1i -
b_2*k_1*m_s*m_us2*omega^5*1i - b_2*k_1*m_us1*m_us2*omega^5*1i -
b*m_s*m_us1*m_us2*omega^7*1i + b_2*m_s*m_us1*m_us2*omega^7*1i + k*k_1*k_t1*m_s*omega^2 +
k*k_1*k_t2*m_s*omega^2 + k*k_2*k_t1*m_s*omega^2 + k*k_2*k_t2*m_s*omega^2 +
k_1*k_2*k_t1*m_s*omega^2 + k_1*k_2*k_t2*m_s*omega^2 + k_1*k_2*k_t1*m_us2*omega^2 +
k_1*k_t1*k_t2*m_s*omega^2 + k_2*k_t1*k_t2*m_s*omega^2 - k*k_1*m_s*m_us1*omega^4 -
k*k_2*m_s*m_us1*omega^4 - k*k_2*m_s*m_us2*omega^4 - k_1*k_2*m_s*m_us1*omega^4 -
k_1*k_2*m_s*m_us2*omega^4 - k_1*k_2*m_us1*m_us2*omega^4 - k_1*k_t1*m_s*m_us2*omega^4 -
k_1*k_t2*m_s*m_us1*omega^4 - k_2*k_t1*m_s*m_us2*omega^4 - k_2*k_t2*m_s*m_us1*omega^4 +
k_1*m_s*m_us1*m_us2*omega^6 + k_2*m_s*m_us1*m_us2*omega^6);
amp_ratio_TF6 =
sqrt((real(TF6_num))^2+(imag(TF6_num))^2)/sqrt((real(TF6_denom))^2+(imag(TF6_denom))^2);

```

```

figure(11);
subplot(2,1,1);
fplot(amp_ratio_TF6,[0,1000]);
title('Transverse Damper Half Car Model Amplitude Ratio of V_m_u_s_2 / V_i_n');
xlabel('Frequency [Hz]');
ylabel('Amplitude Ratio of V_m_u_s_2 / V_i_n');
ylim([0 1]);
phi_6 = atan2d((imag(TF6_num)),real(TF6_num))-atan2d(imag(TF6_denom),real(TF6_denom));
phi_6 =subs(phi_6,omega,[0:1000]);
phi_6 = double(phi_6);
phi_6 = phi_6.*(phi_6 >= 0) + (phi_6 + 360) .* (phi_6 < 0);
subplot(2,1,2);
plot(0:1000,phi_6);
title('Transverse Damper Half Car Model Phase Angle of V_m_u_s_2 / V_i_n');
xlabel('Frequency [Hz]');
ylabel('Phase Angle [deg]');

```

```

%Plotting Transfer Function 8, sprung mass velocity to V_in
TF7 num = (k_t1*(k*k_1*m_s*omega^2 - b_1*b_2*m_s*omega^4 + b*k*m_s*omega^3*2i +
b*k_1*m_s*omega^3*1i + b_1*k*m_s*omega^3*1i + b_2*k*m_s*omega^3*1i +
b*k_1*m_us2*omega^3*1i + b*k_t2*m_s*omega^3*1i + b_1*k_2*m_s*omega^3*1i +
b_2*k_1*m_s*omega^3*1i + b_1*k_t2*m_s*omega^3*1i - b*m_s*m_us2*omega^5*1i -
b_1*m_s*m_us2*omega^5*1i - b*b_1*m_s*omega^4 + k*k_2*m_s*omega^2 - k*k_1*m_us2*omega^2 +
k_1*k_2*m_s*omega^2 + k_1*k_t2*m_s*omega^2 - k_1*m_s*m_us2*omega^4));
TF7 denom = (b*b_1^2*m_s*omega^5*1i + b*b_1*m_s^2*omega^6 + b*b_2*m_s^2*omega^6 +
b_1^2*b_2*m_s*omega^5*1i + b_1*b_2*m_s^2*omega^6 - b*k*m_s^2*omega^5*4i +
b_1^2*k*k_m_s*omega^4 - b*k_1*m_s^2*omega^5*1i - b_1*k*m_s^2*omega^5*1i -
b*k_2*m_s^2*omega^5*1i - b_2*k*m_s^2*omega^5*1i - b*k_t1*m_s^2*omega^5*1i -
b*k_t2*m_s^2*omega^5*1i + b_1^2*k_2*m_s*omega^4 - b_1*k_2*m_s^2*omega^5*1i -
b_2*k_1*m_s^2*omega^5*1i + b_1^2*k_t2*m_s*omega^4 - b_1*k_t2*m_s^2*omega^5*1i -
b_2*k_t1*m_s^2*omega^5*1i + b*m_s^2*m_us1*omega^7*1i + b*m_s^2*m_us2*omega^7*1i -
b_1^2*m_s*m_us2*omega^6 + b_1*m_s^2*m_us2*omega^7*1i + b_2*m_s^2*m_us1*omega^7*1i -
k*k_1*m_s^2*omega^4 - k*k_2*m_s^2*omega^4 - k*k_t1*m_s^2*omega^4 - k*k_t2*m_s^2*omega^4 -
k_1*k_2*m_s^2*omega^4 - k_1*k_t2*m_s^2*omega^4 - k_2*k_t1*m_s^2*omega^4 -
k_t1*k_t2*m_s^2*omega^4 + k*m_s^2*m_us1*omega^6 + k*m_s^2*m_us2*omega^6 +
k_1*m_s^2*m_us2*omega^6 + k_2*m_s^2*m_us1*omega^6 + k_t1*m_s^2*m_us2*omega^6 +
k_t2*m_s^2*m_us1*omega^6 - m_s^2*m_us1*m_us2*omega^8 + 3*b*b_1*k*m_s*omega^4 -
b*b_2*k*m_s*omega^4 + b*b_1*k_1*m_s*omega^4 - b*b_2*k_1*m_s*omega^4 +
b_1*b_2*k*k_m_s*omega^4 + b*b_1*k_1*m_s*omega^4 + b*b_1*k_2*m_s*omega^4 - b*b_2*k_1*m_us2*omega^4 +
b*b_1*k_t2*m_s*omega^4 - b*b_2*k_t2*m_s*omega^4 + b_1*b_2*k_1*m_s*omega^4 -
b_1*b_2*k_t2*m_s*omega^4 - b*b_1*m_s*m_us2*omega^6 + b*b_2*m_s*m_us2*omega^6 +
b_1*b_2*m_s*m_us2*omega^6 + b*k*k_1*m_s*omega^3*1i + b*k*k_2*m_s*omega^3*1i +
b*k*k_1*m_us2*omega^3*1i + b*k*k_t1*m_s*omega^3*1i + b*k*k_t2*m_s*omega^3*1i -
b_1*k*k_1*m_s*omega^3*1i + b*k_1*k_2*m_s*omega^3*1i - b_1*k*k_2*m_s*omega^3*1i +
b_1*k*k_k_1*m_us2*omega^3*1i + b*k_1*k_2*m_us2*omega^3*1i + b*k_1*k_t1*m_s*omega^3*1i +
b_2*k*k_k_1*m_s*omega^3*1i + b*k_2*k_t2*m_s*omega^3*1i + b_2*k*k_k_2*m_s*omega^3*1i +
b*k_1*k_t1*m_us2*omega^3*1i + b*k_t1*k_t2*m_s*omega^3*1i - b_1*k_1*k_2*m_s*omega^3*1i -
b_1*k_1*k_t2*m_s*omega^3*1i + b_2*k_1*k_t1*m_s*omega^3*1i + b_1*k_2*k_t2*m_s*omega^3*1i +
b_2*k_1*k_t2*m_s*omega^3*1i + b_2*k_1*k_t1*m_us2*omega^3*1i +
b_2*k_t1*k_t2*m_s*omega^3*1i - b*k*m_s*m_us1*omega^5*1i - b*k*m_s*m_us2*omega^5*1i -

```

```

b*k_1*m_s*m_us1*omega^5*1i - b*k_1*m_s*m_us2*omega^5*1i - b_2*k*m_s*m_us1*omega^5*1i -
b*k_2*m_s*m_us2*omega^5*1i - b_2*k*m_s*m_us2*omega^5*1i - b*k_1*m_us1*m_us2*omega^5*1i -
b*k_t1*m_s*m_us2*omega^5*1i - b*k_t2*m_s*m_us1*omega^5*1i + b_1*k_1*m_s*m_us2*omega^5*1i -
b_2*k_1*m_s*m_us1*omega^5*1i - b_1*k_2*m_s*m_us2*omega^5*1i -
b_2*k_1*m_s*m_us2*omega^5*1i - b_2*k_2*m_s*m_us1*omega^5*1i +
b*m_s*m_us1*m_us2*omega^7*1i + b_2*m_s*m_us1*m_us2*omega^7*1i + k*k_1*k_t1*m_s*omega^2 +
k*k_1*k_t2*m_s*omega^2 + k*k_2*k_t1*m_s*omega^2 + k*k_2*k_t2*m_s*omega^2 +
k_1*k_2*k_t1*m_s*omega^2 + k_1*k_2*k_t2*m_s*omega^2 + k_1*k_2*k_t1*m_us2*omega^2 +
k_1*k_t1*k_t2*m_s*omega^2 + k_2*k_t1*k_t2*m_s*omega^2 - k*k_1*m_s*m_us1*omega^4 -
k*k_2*m_s*m_us1*omega^4 - k*k_2*m_s*m_us2*omega^4 - k_1*k_2*m_s*m_us1*omega^4 -
k_1*k_2*m_s*m_us2*omega^4 - k_1*k_2*m_us1*m_us2*omega^4 - k_1*k_t1*m_s*m_us2*omega^4 -
k_1*k_t2*m_s*m_us1*omega^4 - k_2*k_t1*m_s*m_us2*omega^4 - k_2*k_t2*m_s*m_us1*omega^4 +
k_1*m_s*m_us1*m_us2*omega^6 + k_2*m_s*m_us1*m_us2*omega^6);
amp_ratio_TF7 =
sqrt((real(TF7_num))^2+(imag(TF7_num))^2)/sqrt((real(TF7_denom))^2+(imag(TF7_denom))^2);

```

```

figure(12);
subplot(2,1,1)
fplot(amp_ratio_TF7,[0 1000]);
title('Transverse Damper Half Car Model Amplitude Ratio of V_m_s / V_i_n');
xlabel('Frequency [Hz]');
ylabel('Amplitude Ratio of V_m_s / V_i_n');
ylim([0 2]);
phi_7 = atan2d((imag(TF7_num)),real(TF7_num))-atan2d(imag(TF7_denom),real(TF7_denom));
phi_7 =subs(phi_7,omega,[0:1000]);
phi_7 = double(phi_7);
phi_7 = phi_7.*(phi_7 >= 0) + (phi_7 + 360) .* (phi_7 < 0);
subplot(2,1,2);
plot(0:1000,phi_7);
title('Transverse Damper Half Car Model Phase Angle of V_m_s / V_i_n');
xlabel('Frequency [Hz]');
ylabel('Phase Angle [deg]');

```

```

figure(13)
fplot(amp_ratio_TF5,[0,1000]);
hold on
fplot(amp_ratio_TF6,[0,1000]);
ylim([0 1.5]);
xlabel('Frequency [Hz]');
ylabel('Amplitude Ratio');
legend('Left Wheel (experiences bump)','Right Wheel (changes velocity from sway bar
connection damped by transverse damper)');
title('Transverse Damper Half Car Model Comparison of Amplitude Ratio Frequency Response
of Left and Right Wheels');

```

```

figure (14)
fplot(amp_ratio_TF6,[0,1000]);
hold on
fplot(amp_ratio_TF3,[0,1000]);
xlabel('Frequency [Hz]');
ylabel('Amplitude Ratio');
legend('Right Wheel with Transverse Damper','Right Wheel Without Transverse Damper');
title('Comparison of Right Wheel Amplitude Ratio Frequency Response With and Without
Transverse Damper');

```

*%Plotting Frequency Response of Transverse Damper System with Various  
%Damping Coefficients*

```

figure (15);
for b = 0:b_1:0.2*b_1:3*b_1
TF6_num = (k_t1*m_s*omega*(k*m_s*omega^3*1i - k*k_1*omega*1i - k*k_2*omega*1i -
k_1*k_2*omega*1i + b*b_1*omega^3*1i + b_1*b_2*omega^3*1i + b*k*omega^2 + b_2*k*omega^2 +
b_1*k_2*omega^2 + b_2*k_1*omega^2 + b*m_s*omega^4)*1i);
TF6_denom = (b*b_1^2*m_s*omega^5*1i + b*b_1*m_s^2*omega^6 + b*b_2*m_s^2*omega^6 +
b_1^2*b_2*m_s*omega^5*1i + b_1*b_2*m_s^2*omega^6 - b*k*m_s^2*omega^5*4i +
b_1^2*k*m_s*omega^4 - b*k_1*m_s^2*omega^5*1i - b_1*k*m_s^2*omega^5*1i -
b*k_2*m_s^2*omega^5*1i - b_2*k*m_s^2*omega^5*1i - b*k_t1*m_s^2*omega^5*1i -
b*k_t2*m_s^2*omega^5*1i + b_1^2*k_2*m_s*omega^4 - b_1*k_2*m_s^2*omega^5*1i -

```

```

b_2*k_1*m_s^2*omega^5*1i + b_1^2*k_t2*m_s*omega^4 - b_1*k_t2*m_s^2*omega^5*1i -
b_2*k_t1*m_s^2*omega^5*1i + b*m_s^2*m_us1*omega^7*1i + b*m_s^2*m_us2*omega^7*1i -
b_1^2*m_s*m_us2*omega^6 + b_1*m_s^2*m_us2*omega^7*1i + b_2*m_s^2*m_us1*omega^7*1i -
k*k_1*m_s^2*omega^4 - k*k_2*m_s^2*omega^4 - k*k_t1*m_s^2*omega^4 - k*k_t2*m_s^2*omega^4 -
k_1*k_2*m_s^2*omega^4 - k_1*k_t2*m_s^2*omega^4 - k_2*k_t1*m_s^2*omega^4 -
k_t1*k_t2*m_s^2*omega^4 + k*m_s^2*m_us1*omega^6 + k*m_s^2*m_us2*omega^6 +
k_1*m_s^2*m_us2*omega^6 + k_2*m_s^2*m_us1*omega^6 + k_t1*m_s^2*m_us2*omega^6 +
k_t2*m_s^2*m_us1*omega^6 - m_s^2*m_us1*m_us2*omega^8 + 3*b*b_1*k*m_s*omega^4 -
b*b_2*k*m_s*omega^4 + b*b_1*k_1*m_s*omega^4 - b*b_2*k_1*m_s*omega^4 +
b_1*b_2*k*m_s*omega^4 + b*b_1*k_1*m_us2*omega^4 - b*b_2*k_1*m_us2*omega^4 +
b*b_1*k_t2*m_s*omega^4 - b*b_2*k_t2*m_s*omega^4 + b_1*b_2*k_1*m_s*omega^4 -
b_1*k*k_1*m_us2*omega^6 + b*b_1*m_s*m_us1*omega^6 + b*b_2*m_s*m_us2*omega^6 +
b_1*b_2*m_s*m_us2*omega^6 + b*k*k_1*m_s*omega^3*1i + b*k*k_2*m_s*omega^3*1i +
b*k*k_1*m_us2*omega^3*1i + b*k*k_t1*m_s*omega^3*1i + b*k*k_t2*m_s*omega^3*1i -
b_1*k*k_1*m_s*omega^3*1i + b*k_1*k_2*m_s*omega^3*1i - b_1*k*k_2*m_s*omega^3*1i +
b_1*k*k_1*m_us2*omega^3*1i + b*k_1*k_2*m_us2*omega^3*1i + b*k_1*k_t1*m_s*omega^3*1i +
b_2*k*k_t1*m_s*omega^3*1i + b*k_2*k_t2*m_s*omega^3*1i + b_2*k*k_t2*m_s*omega^3*1i +
b*k_1*k_t1*m_us2*omega^3*1i + b*k_t1*k_t2*m_s*omega^3*1i - b_1*k_1*k_2*m_s*omega^3*1i -
b_1*k_1*k_t2*m_s*omega^3*1i + b_2*k_1*k_t1*m_s*omega^3*1i + b_1*k_2*k_t2*m_s*omega^3*1i +
b_2*k_1*k_t2*m_s*omega^3*1i + b_2*k_1*k_t1*m_us2*omega^3*1i +
b_2*k_t1*k_t2*m_s*omega^3*1i - b*k*m_s*m_us1*omega^5*1i - b*k*m_s*m_us2*omega^5*1i -
b*k_1*m_s*m_us1*omega^5*1i - b*k_1*m_s*m_us2*omega^5*1i - b_2*k*m_s*m_us1*omega^5*1i -
b_2*k*m_s*m_us2*omega^5*1i - b_2*k*m_s*m_us2*omega^5*1i - b*k_1*m_us1*m_us2*omega^5*1i -
b*k_t1*m_s*m_us2*omega^5*1i - b*k_t2*m_s*m_us1*omega^5*1i + b_1*k_1*m_s*m_us2*omega^5*1i
- b_2*k_1*m_s*m_us1*omega^5*1i - b_1*k_2*m_s*m_us2*omega^5*1i -
b_2*k_1*m_s*m_us2*omega^5*1i - b_2*k_1*m_us1*m_us2*omega^5*1i -
b_2*k_t1*m_s*m_us2*omega^5*1i - b_2*k_t2*m_s*m_us1*omega^5*1i +
b*m_s*m_us1*m_us2*omega^7*1i + b_2*m_s*m_us1*m_us2*omega^7*1i + k*k_1*k_t1*m_s*omega^2 +
k*k_1*k_t2*m_s*omega^2 + k*k_2*k_t1*m_s*omega^2 + k*k_2*k_t2*m_s*omega^2 +
k_1*k_2*k_t1*m_s*omega^2 + k_1*k_2*k_t2*m_s*omega^2 + k_1*k_2*k_t1*m_us2*omega^2 +
k_1*k_t1*k_t2*m_s*omega^2 + k_2*k_t1*k_t2*m_s*omega^2 - k*k_1*m_s*m_us1*omega^4 -
k*k_2*m_s*m_us1*omega^4 - k*k_2*m_s*m_us2*omega^4 - k_1*k_2*m_s*m_us1*omega^4 -
k_1*k_2*m_s*m_us2*omega^4 - k_1*k_t1*m_s*m_us1*omega^4 - k_1*k_t1*m_s*m_us2*omega^4 -
k_1*k_t2*m_s*m_us1*omega^4 - k_2*k_t1*m_s*m_us2*omega^4 - k_2*k_t2*m_s*m_us1*omega^4 +
k_1*m_s*m_us1*m_us2*omega^6 + k_2*m_s*m_us1*m_us2*omega^6);
amp_ratio_TF6MR =
sqrt((real(TF6_num))^2+(imag(TF6_num))^2)/sqrt((real(TF6_denom))^2+(imag(TF6_denom))^2);
fplot(amp_ratio_TF6MR, [0 600]);
legend('Right Wheel with Transverse Damper')
hold on
end
Legend = cell(16,1);
Legend{1} = 'b = 0.0*b_1';
Legend{2} = 'b = 0.2*b_1';
Legend{3} = 'b = 0.4*b_1';
Legend{4} = 'b = 0.6*b_1';
Legend{5} = 'b = 0.8*b_1';
Legend{6} = 'b = 1.0*b_1';
Legend{7} = 'b = 1.2*b_1';
Legend{8} = 'b = 1.4*b_1';
Legend{9} = 'b = 1.6*b_1';
Legend{10} = 'b = 1.8*b_1';
Legend{11} = 'b = 2.0*b_1';
Legend{12} = 'b = 2.2*b_1';
Legend{13} = 'b = 2.4*b_1';
Legend{14} = 'b = 2.6*b_1';
Legend{15} = 'b = 2.8*b_1';
Legend{16} = 'b = 3.0*b_1';

legend(Legend);
title('Frequency Response of Right Wheel from Left Wheel Disturbance in 3rd Damper
System, Over Many Transverse Damping Coefficients');
xlabel('Frequency [Hz]');
ylabel('Amplitude Ratio');

```

